Detecting the Cosmic Neutrino Background

Jack Shergold



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What we will cover

• Lecture 1: Introduction to the $C\nu B$

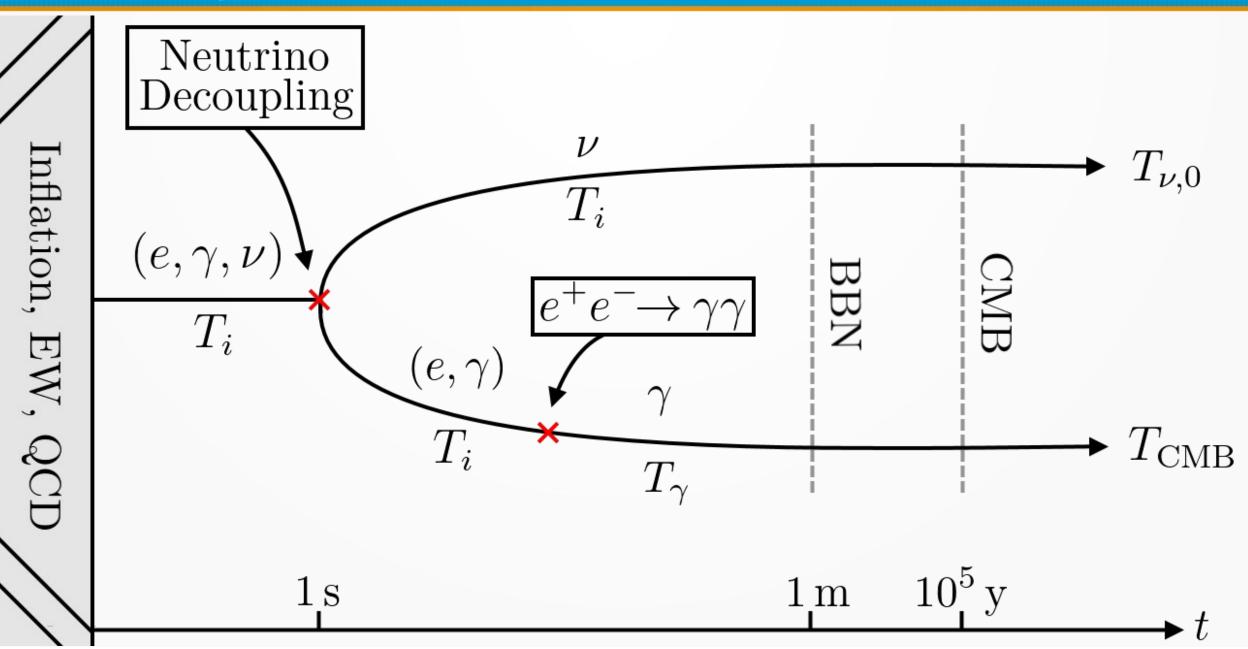
• Lecture 2: Direct detection proposals

 Lecture 3: Indirect detection, constraints and future prospects

Contents

- Recap
- Direct detection proposals:
 - PTOLEMY
 - Coherent scattering
 - Stodolsky effect
 - Accelerator





• Key properties: $T_{\nu,0} = \left(\frac{4}{11}\right)^{\frac{1}{3}} T_{\text{CMB}}$



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 $T_{\nu,0} = 0.168 \,\mathrm{meV}$



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• All left helicity (right helicity antineutrinos)

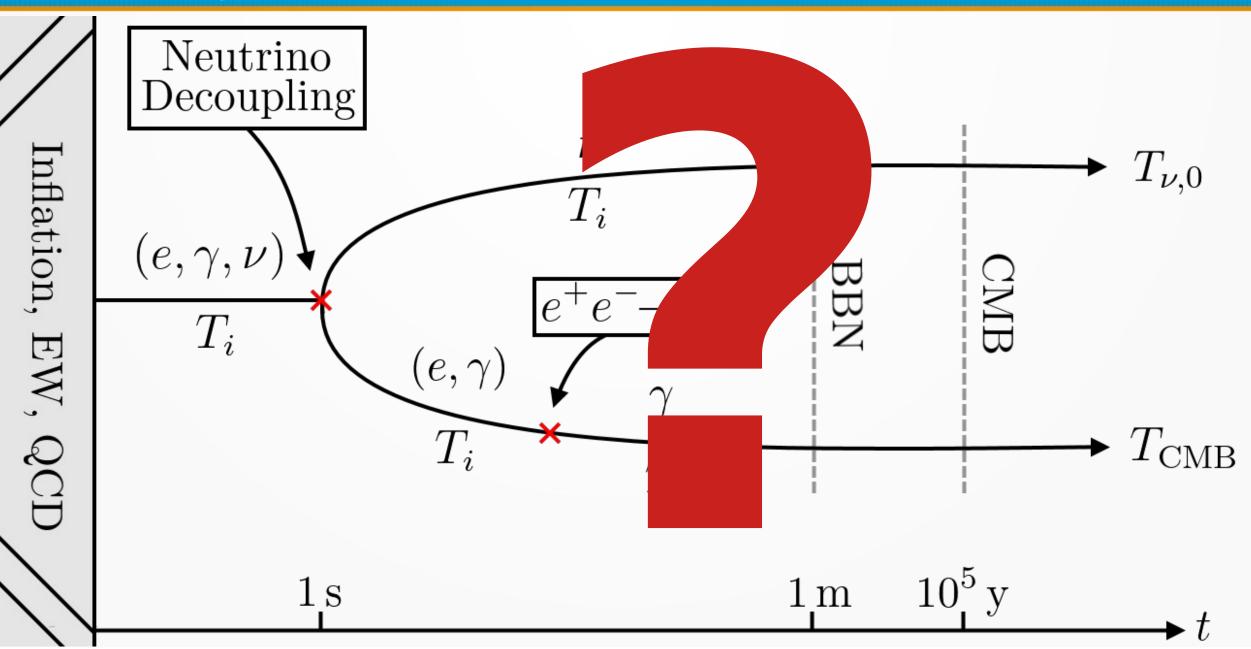
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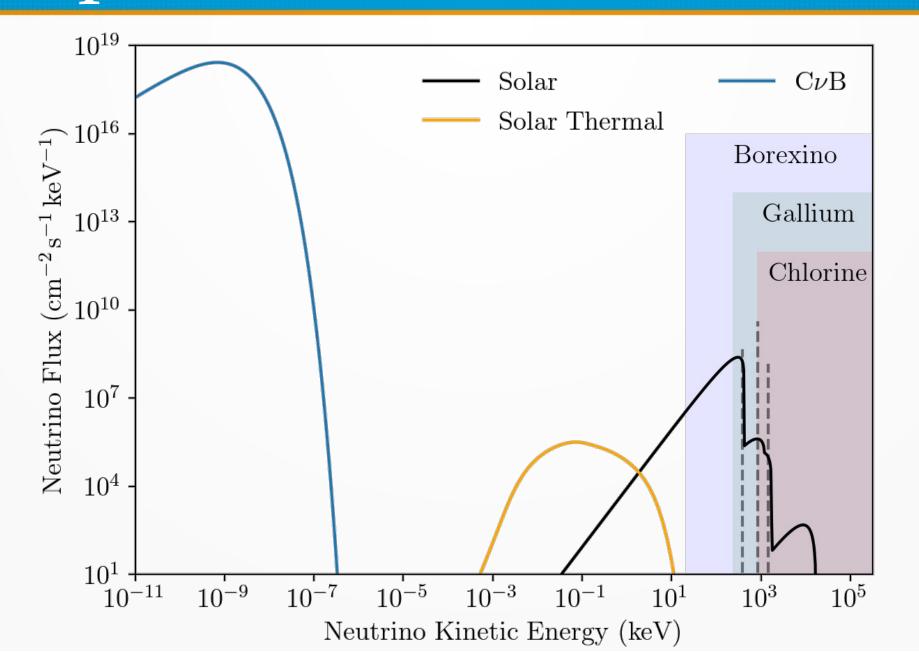
 $T_{\nu,0} = 0.168 \,\mathrm{meV}$ $|\nu_{\alpha}\rangle \to |U_{\alpha i}|^2 |\nu_i\rangle$

- All left helicity (right helicity antineutrinos)
- Follow massless Fermi-Dirac distribution:

$$f(p) = \frac{1}{\exp\left(\frac{p}{T}\right) + 1}$$







How to detect the $C\nu B$

Direct

- Unstable nuclei
- Coherent scattering
- Stodolsky effect

Indirect

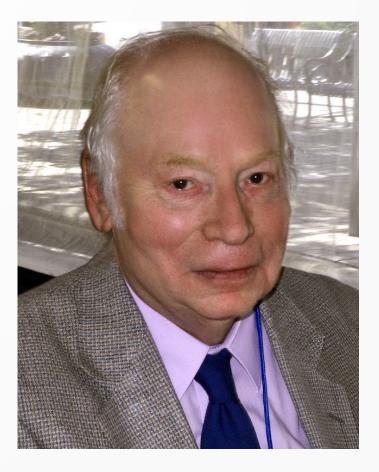
- Cosmic ray attenuation
- Pauli blocking
- Neutron star cooling

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• Proposed by Weinberg in 1962 [1]:

$$\nu_e + {}^3\mathrm{H} \to e^- + {}^3\mathrm{He}^+$$



[1] S. Weinberg, Phys. Rev. **128**, 1457 (1962)

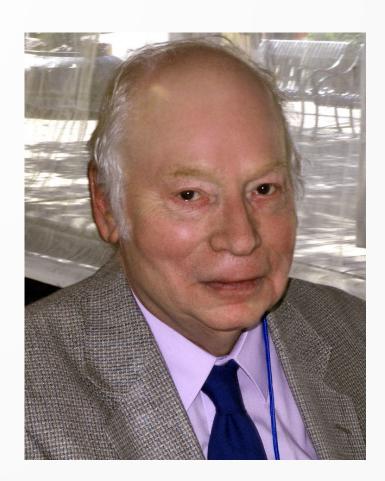
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• This process has no threshold:

 $^{3}\mathrm{H} \rightarrow ^{3}\mathrm{He}^{+} + e^{-} + \bar{\nu}_{e}$

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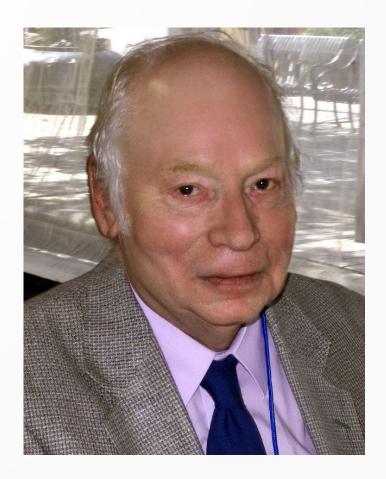
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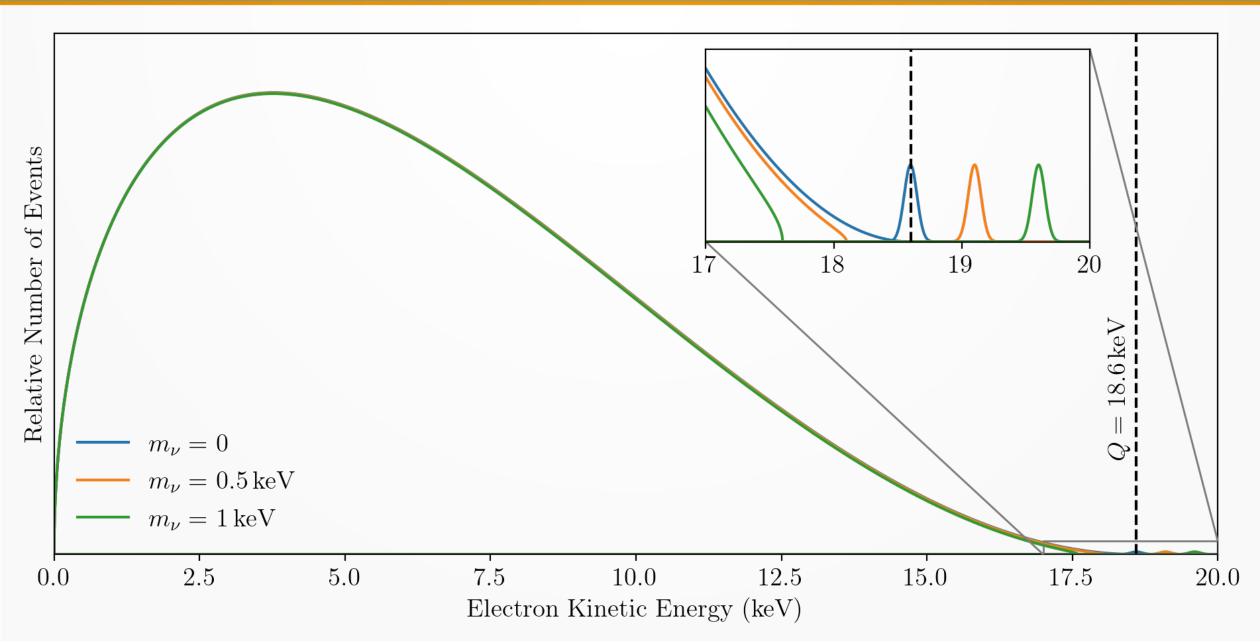
• This process has no threshold:

 $^{3}\mathrm{H} \rightarrow ^{3}\mathrm{He}^{+} + e^{-} + \bar{\nu}_{e}$

• Beta decay releases 18.6 keV



[1] S. Weinberg, Phys. Rev. **128**, 1457 (1962)



• Neutrino capture cross section [2]:

$$\langle \sigma \beta_{\nu} \rangle \propto G_F^2 E_e p_e \sim 10^{-45} \,\mathrm{cm}^2$$

[2] A. Long, C. Lunardini and E. Sabancilar, JCAP 08, 038 (2014)

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• Giving an event rate for 100g of tritium:

$$\Gamma = N_T n_\nu \langle \sigma \beta_\nu \rangle \sim 4 \, \mathrm{y}^{-1}$$

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• This event rate is doubled for Majorana neutrinos

[2] A. Long, C. Lunardini and E. Sabancilar, JCAP 08, 038 (2014)

• Extreme sensitivity required to detect signal:

 $\Delta \le 2m_{\nu}$

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• Obtaining and storing 100g of tritium

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- Obtaining and storing 100g of tritium
- cf. KATRIN, uses ~300µg of tritium [3]

[3] L. Kuckert *et al.*, Vacuum **158**, 195-205 (2018)

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 $\Delta\simeq 50\,\mathrm{meV}$

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 $m_{\nu_h} \gtrsim 50 \,\mathrm{meV}$

• Extreme sensitivity required to detect signal:

$$\Delta < 2m_{\nu_h}$$

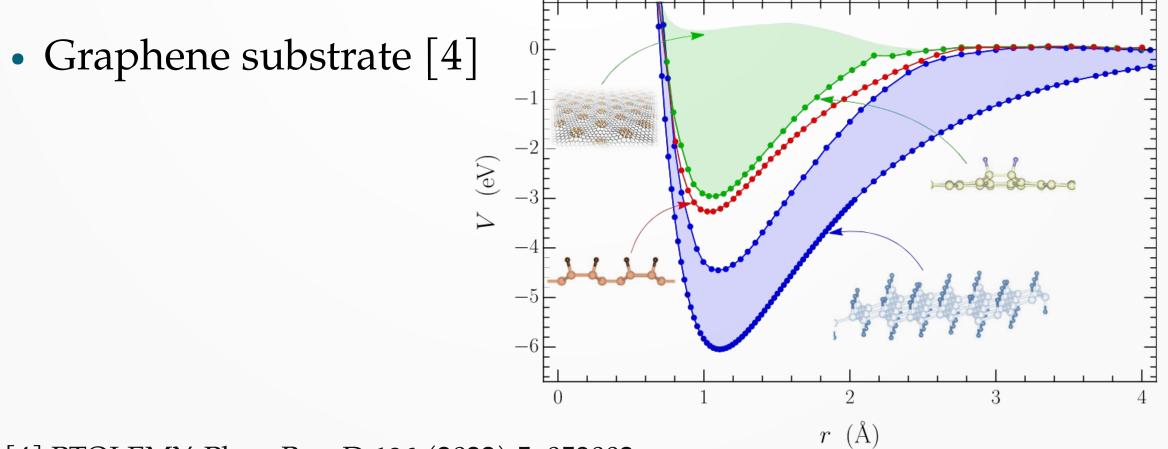
• Extreme sensitivity required to detect signal:

$$\Delta < 2m_{\nu_h}$$

• PTOLEMY should at least see something!

• Obtaining and storing 100g of tritium

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[4] PTOLEMY, Phys. Rev. D 106 (2022) 5, 053002

• Uncertainty principle:

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$$\Delta x \Delta p \simeq \frac{1}{2} \implies \Delta v \simeq \frac{1}{2m_T \Delta x}$$

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 $\Delta E_e \simeq p_e \Delta v$

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$$\Delta x \simeq 0.1 \,\text{\AA}$$

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 $\Delta E_e \gg 2m_{\nu_h}$

What's the catch?

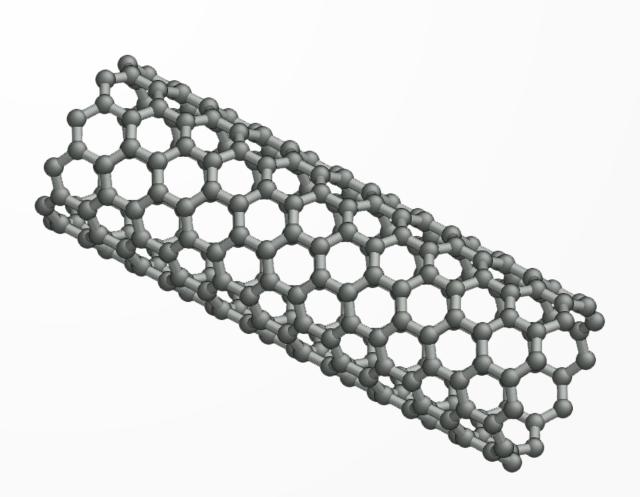
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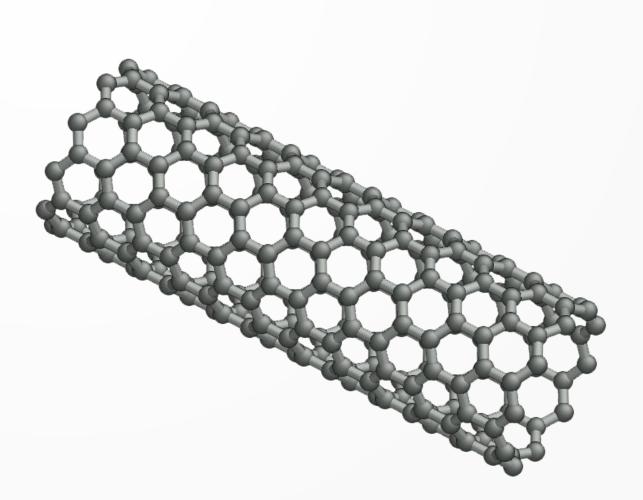
 $\Delta E_e \gg 2m_{\nu_h}$

• No longer guaranteed to see anything!

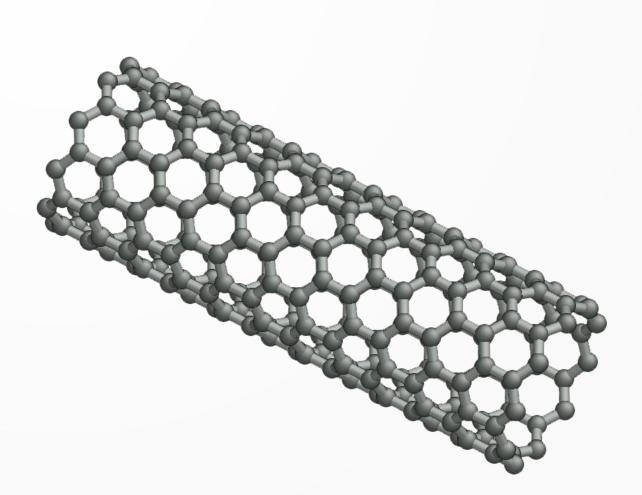
• Carbon nanotubes?



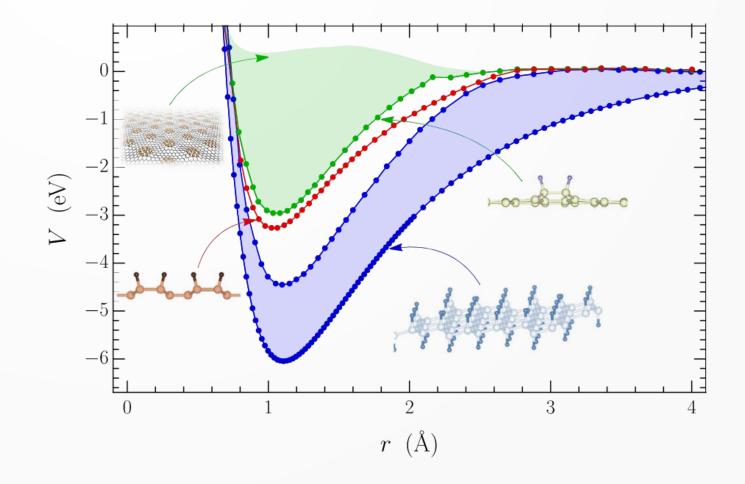
- Carbon nanotubes?
- Allow tritium to *slide*



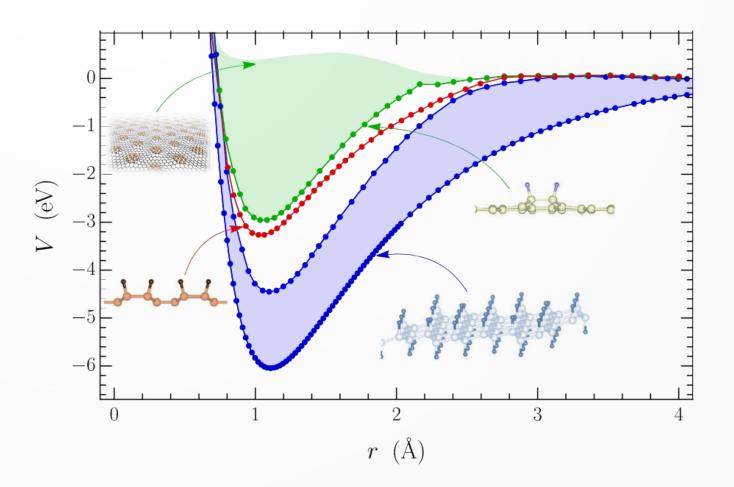
- Carbon nanotubes?
- Allow tritium to *slide*
- Increases $\Delta x!$



• Other ideas:



- Other ideas:
- Modified potentials?



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$$\langle \vec{\beta}_{\nu} \rangle \neq \vec{0}$$

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• Variation of around 0.1% → distinguish with Fourier decomposition:

$$r_{\rm SN} \propto (t - t_0)$$

• Allows for worse energy resolution! [5] E. Akhmedov, JCAP **09** (2019) 031

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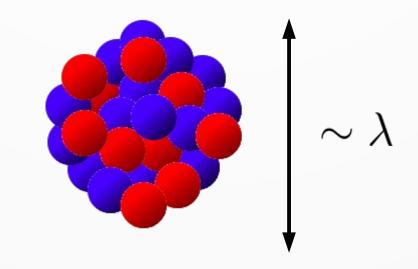
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 $\nu + X \to \nu + X$

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• Wavelength of order target size:



• Neutrino "sees" a larger charge:

$$\mathcal{M} \sim Q_T \qquad \qquad Q_T \sim N_T$$

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$$\mathcal{M} \sim Q_T \qquad \qquad Q_T \sim N_T$$

• Significantly enhanced event rates:

$$\Gamma \sim N_T \longrightarrow \Gamma_{\rm coh} \sim N_T^2$$

• More rigorously [6]:

 $\mathcal{M}_i \propto e^{-i(\vec{q}\cdot\vec{x}_i)}$

[6] M. Bauer, JS, JCAP **01** (2023), 003

• More rigorously [6]:

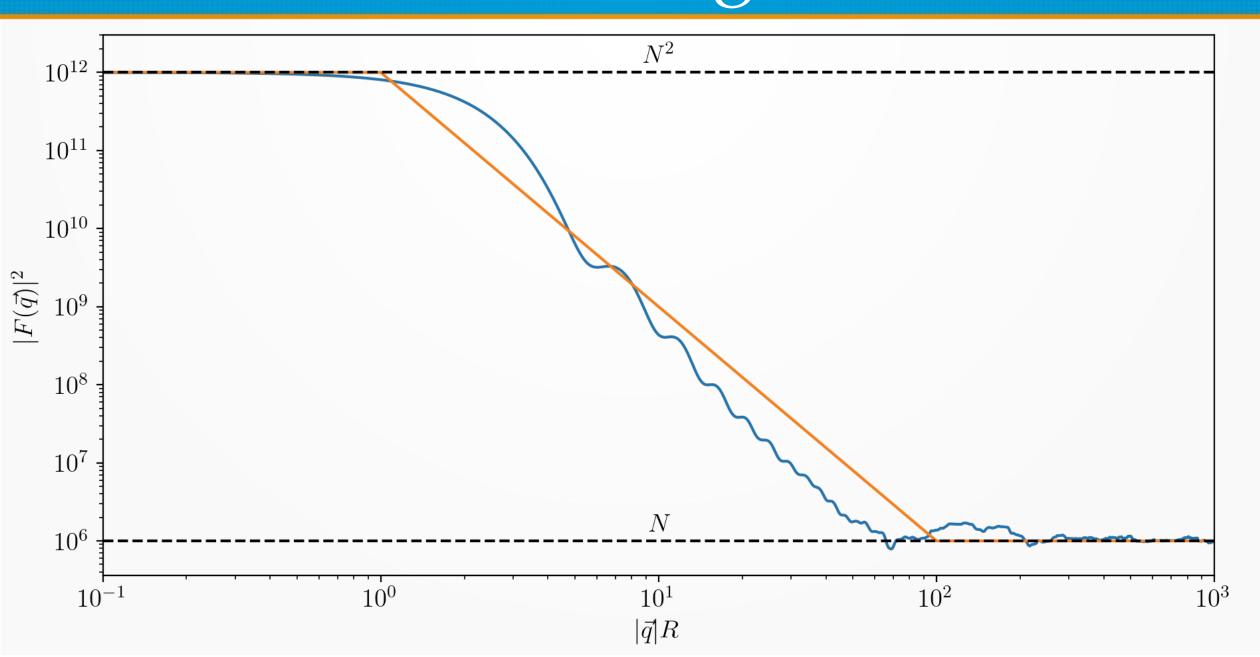
$$|\mathcal{M}|^2 \propto \left| \sum_i e^{-i(\vec{q} \cdot \vec{x}_i)} \right|^2$$

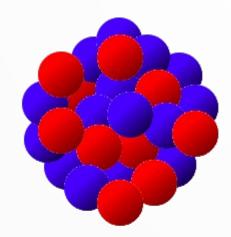
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- More rigorously [6]: $|\mathcal{M}|^2 \propto \left| \sum_i e^{-i(\vec{q} \cdot \vec{x}_i)} \right|^2$
- Small momentum transfer, $\langle q \rangle^{-1} \ll \langle x_i \rangle$:

$$\left|\sum_{i} e^{-i(\vec{q}\cdot\vec{x}_i)}\right|^2 \to N_T^2$$

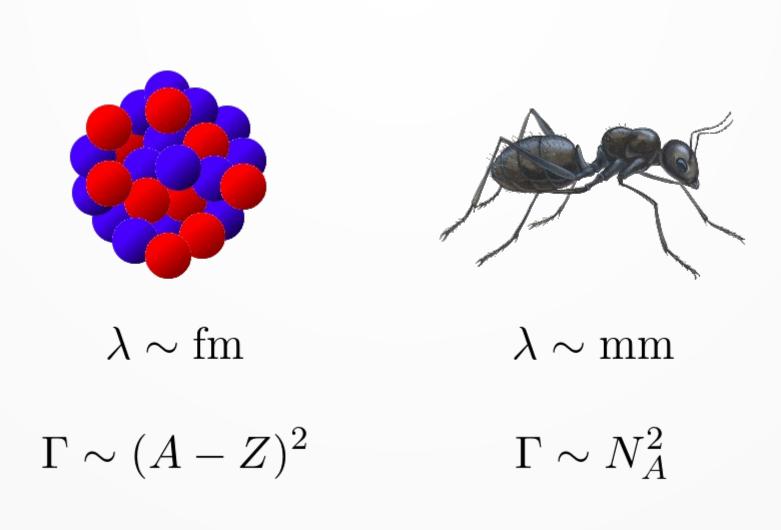
[6] M. Bauer, JS, JCAP **01** (2023), 003





 $\lambda \sim {\rm fm}$

 $\Gamma \sim (A - Z)^2$



• Recall from last lecture:

$$\frac{\Gamma}{N_T} \simeq 10^{-39} \,\mathrm{y}^{-1}$$

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• Per mole of target:

$$\Gamma_{\rm incoh} \simeq 10^{-16} \, {\rm y}^{-1} \qquad \Gamma_{\rm coh} \simeq 10^7 \, {\rm y}^{-1}$$

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• For a typical setup:

 $\Gamma_{\rm coh} \sim \mathcal{O}(\rm kHz)$

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• For a typical setup:

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• So why don't we see anything?

• Relic neutrinos are cold and light:

 $\langle p_{\nu} \rangle \simeq 0.5 \,\mathrm{meV}$

• Relic neutrinos are cold and light:

 $\Delta p_{\nu} \sim \beta_{\oplus} E_{\nu}$

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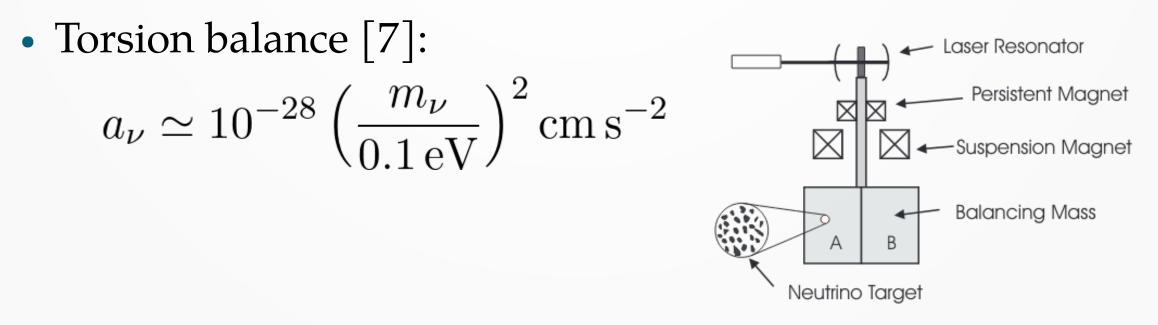


Fig: C. Hagmann, arXiv:astro-ph/9905258

[7] JS, JCAP 11 (2021), 052

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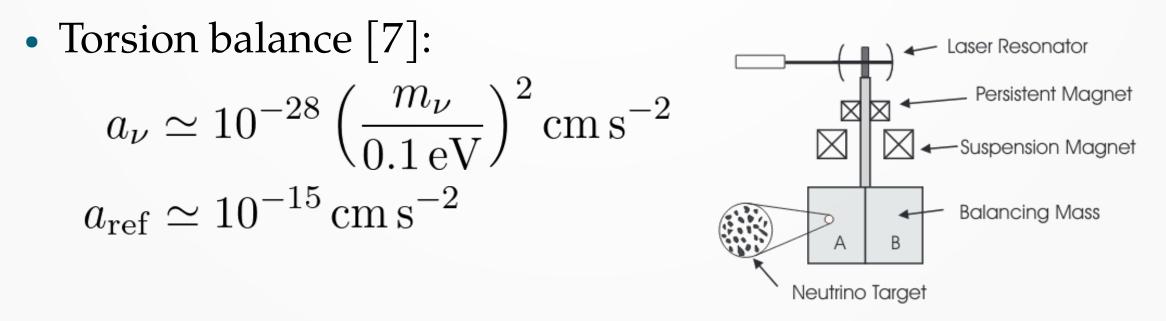


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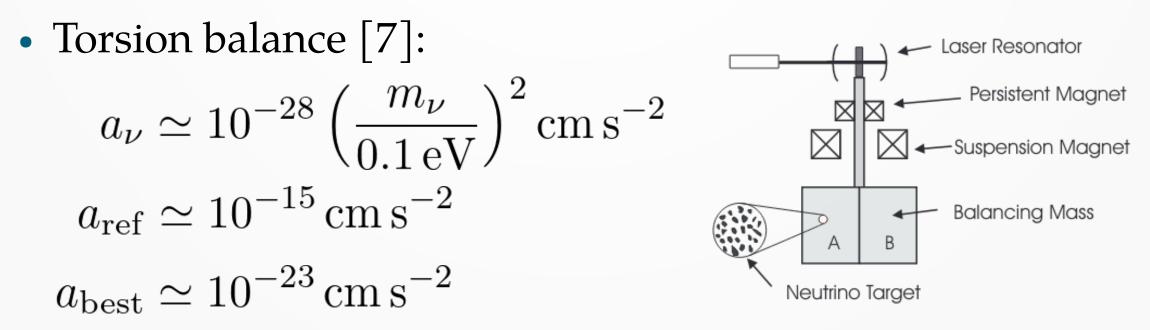


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• Majorana neutrinos:

$$\sigma_M \sim (2 - \beta_\nu^2) \sigma_D$$

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$$\sigma_M \sim (2 - \beta_{\nu}^2) \sigma_D$$

• Should be sensitive to Majorana neutrinos!

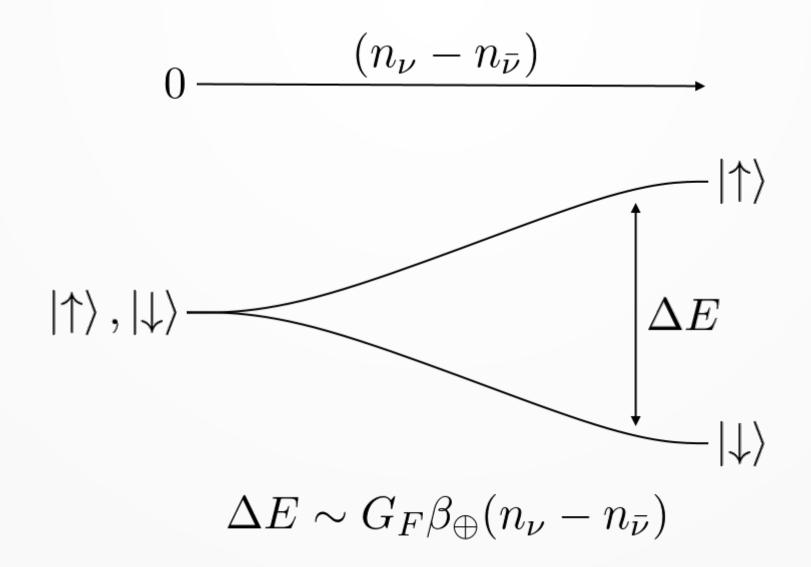
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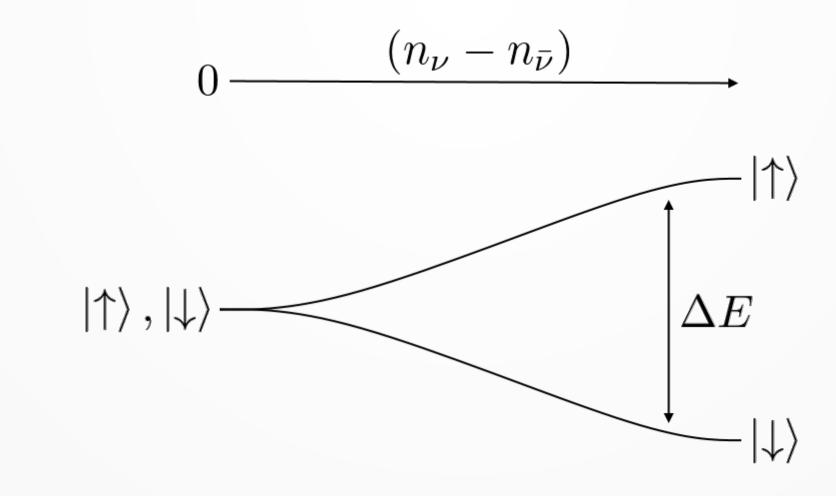
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• Stodolsky effect (Stodolsky '75, Duda '01, JS '23)

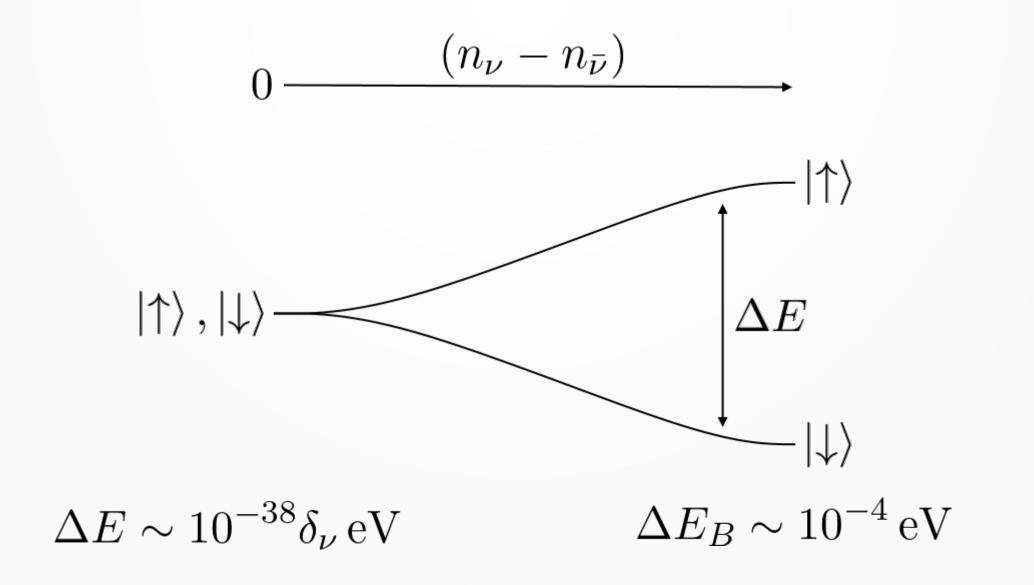
- Stodolsky effect (Stodolsky '75, Duda '01, JS '23)
- Analogous to the Zeeman effect

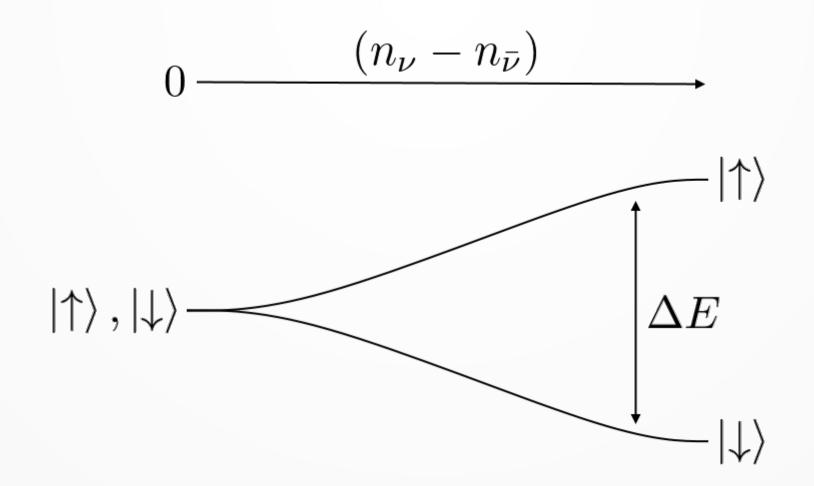
- Stodolsky effect (Stodolsky '75, Duda '01, JS '23)
- Analogous to the Zeeman effect
- Only* effect linear in G_F (Cabibbo '82, Langacker '83)





 $\Delta E \sim 10^{-38} \delta_{\nu} \,\mathrm{eV}$





• Also an effect that depends on helicity asymmetry

• So how do we see it?

- So how do we see it?
- Hamiltonian picks up spin-dependent terms:

 $H = H_0 + H_S$

- So how do we see it?
- Spin precession:

$$\frac{dS_{\perp}}{dt} = i[H, S_{\perp}] \neq 0$$

- Two observables:
- Torque [8]:

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$$a_{\mathrm{ref}} \simeq 10^{-15} \,\mathrm{cm}\,\mathrm{s}^{-2}$$
$$a_{\mathrm{best}} \simeq 10^{-23} \,\mathrm{cm}\,\mathrm{s}^{-2}$$

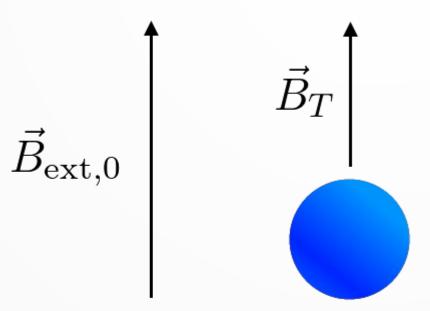
- Two observables:
- Magnetisation [8]:

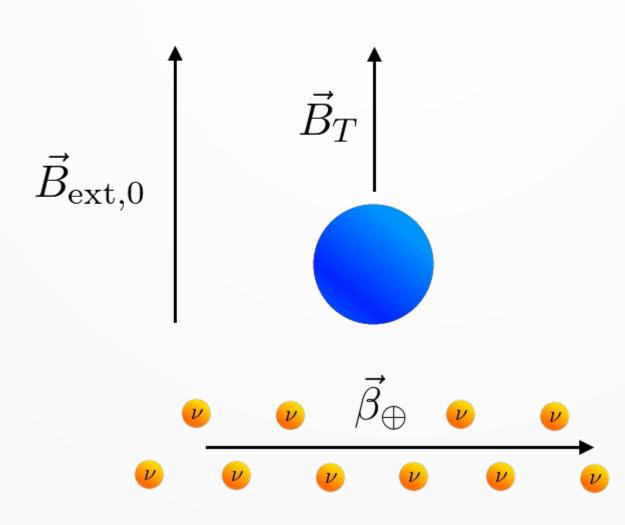
$$\frac{dS_{\perp}}{dt} \neq 0 \implies \frac{dB_{\perp}}{dt} \neq 0$$

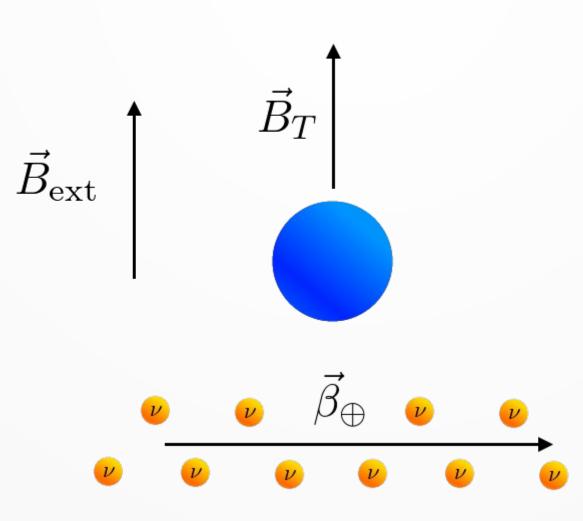


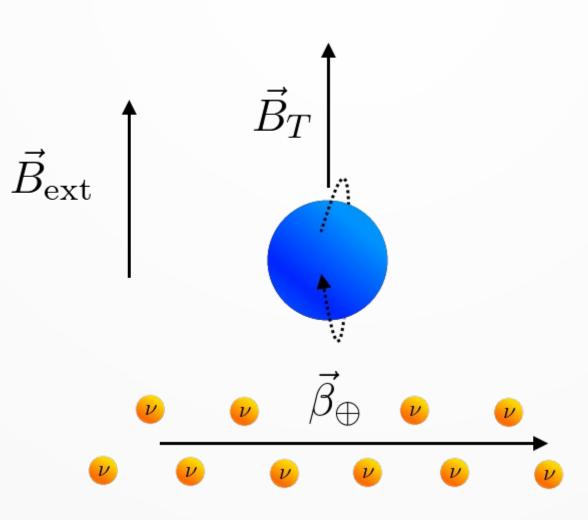
 $ec{B}_{\mathrm{ext},0}$

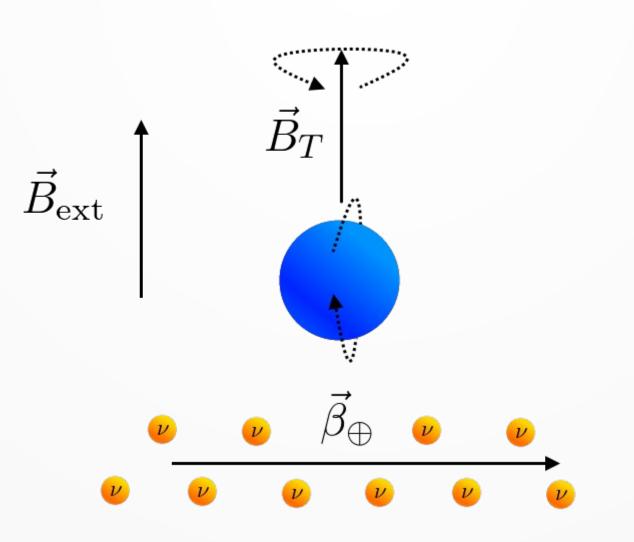


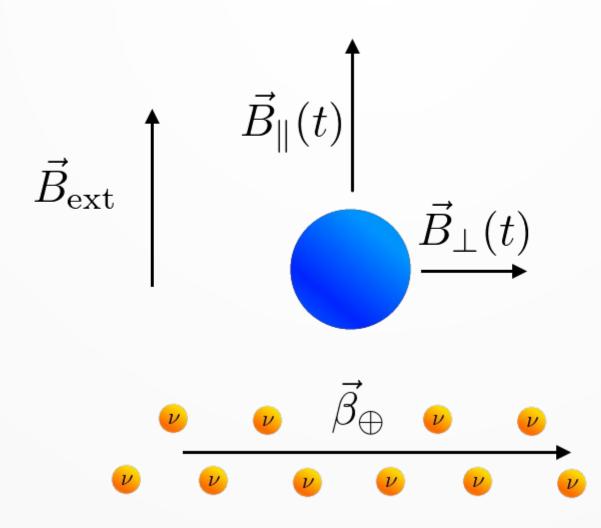


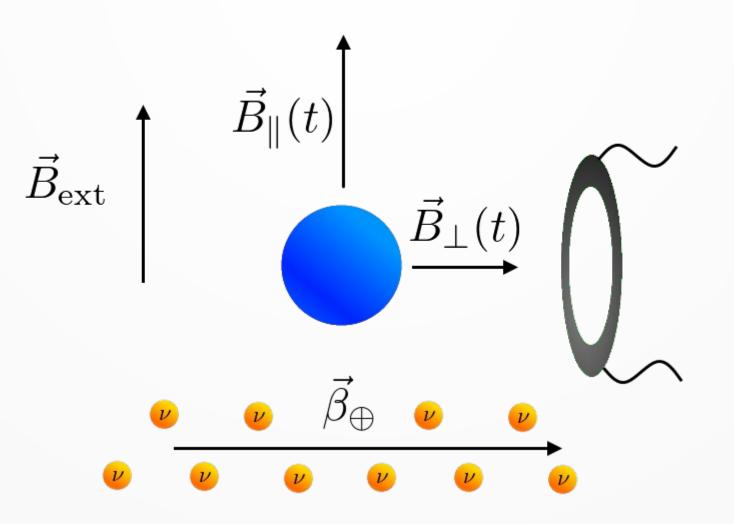


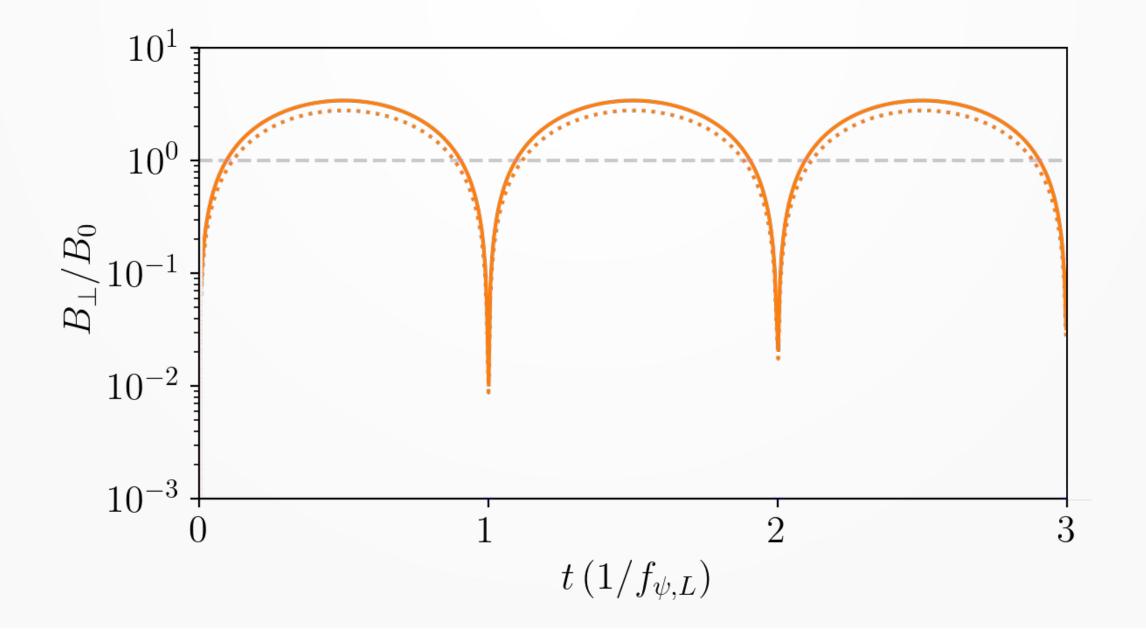












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• Helicity asymmetry term? 3 orders of magnitude larger!

• Asymmetry?

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- BBN constraint:

$$\delta_{\nu} = \delta_h \lesssim 10^{-2}$$

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• Still a way off, but better than coherent

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Stodolsky effect

- Majorana neutrinos?
- Only helicity dependent term survives
- This can be killed by frame transformations:

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• Only a problem for very slow neutrinos

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• PTOLEMY-on-a-beam?

$$\langle \sigma \beta_{\nu} \rangle \propto G_F^2 E_e p_e \sim 10^{-45} \,\mathrm{cm}^2$$

• Doesn't really work...

$$E_{\nu} = \frac{m_{\nu}}{M} E_{\text{beam}} \implies E_{\text{beam}} \sim \mathcal{O}(\text{PeV})$$

• Thresholded processes?

• Tiny cross sections \rightarrow use a resonance!

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- Tunable beam energy naturally invites resonances
- e.g. Z-resonance: $\nu + \bar{\nu}_{C\nu B} \rightarrow Z \rightarrow ?$
- Vastly larger cross section:

$$\sigma \propto \frac{1}{M_Z^2} \propto G_F$$

• Accelerator experiment [9]:

$${}^{A}_{Z}P + e^{-}(\text{bound}) + \bar{\nu}_{e} \rightarrow {}^{A}_{Z-1}D$$

 ${}^{A}_{Z}P + \nu_{e} \rightarrow {}^{A}_{Z+1}D + e^{-}(\text{bound})$

[9] M. Bauer, JS, Phys. Rev. D 104 (2021), 083039

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 ${}^{A}_{Z}P + \nu_{e} \rightarrow {}^{A}_{Z+1}D + e^{-}(\text{bound})$

• Huge cross section on resonance:

$$\sigma_{\rm peak} \propto \frac{1}{Q^2} \operatorname{Br}(D \to P)$$

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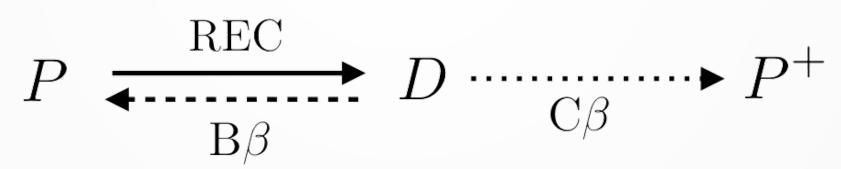
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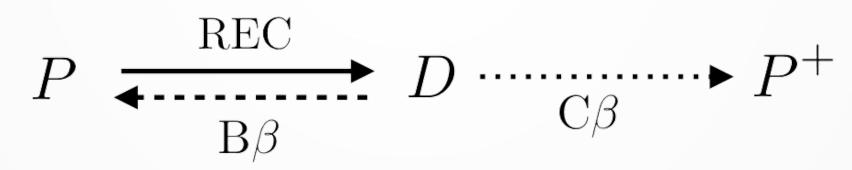
$$\sigma_{\rm peak} \simeq 10^{-15} \left(\frac{1\,{\rm keV}}{Q}\right)^2 {\rm Br}(D \to P)\,{\rm cm}^2$$

[9] M. Bauer, JS, Phys. Rev. D 104 (2021), 083039

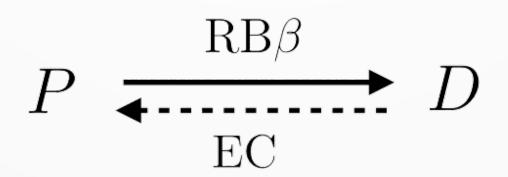
• Resonant electron capture:

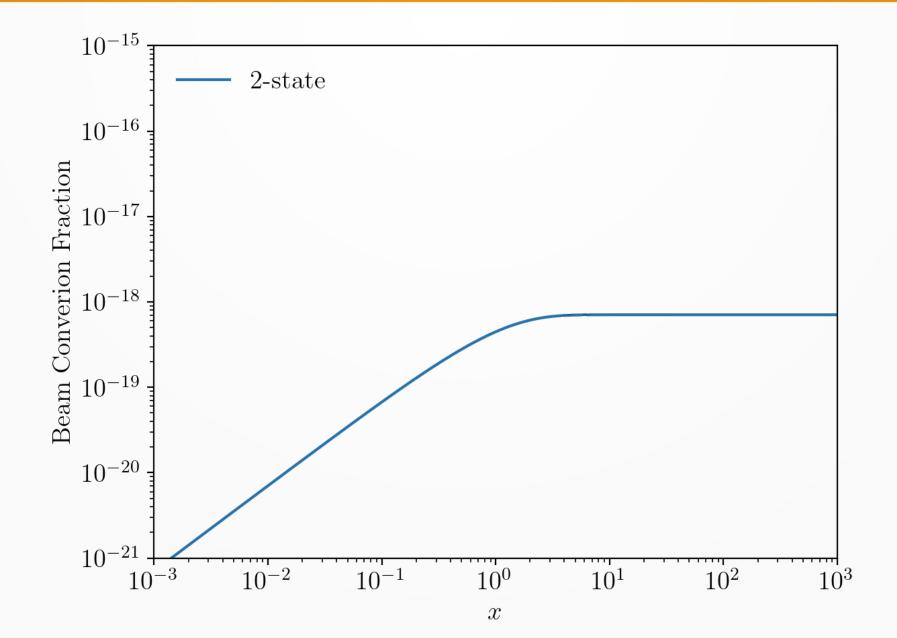


• Resonant electron capture:



• Resonant bound beta decay:

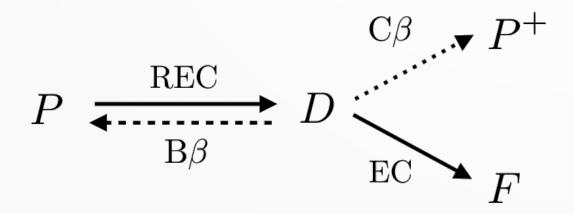




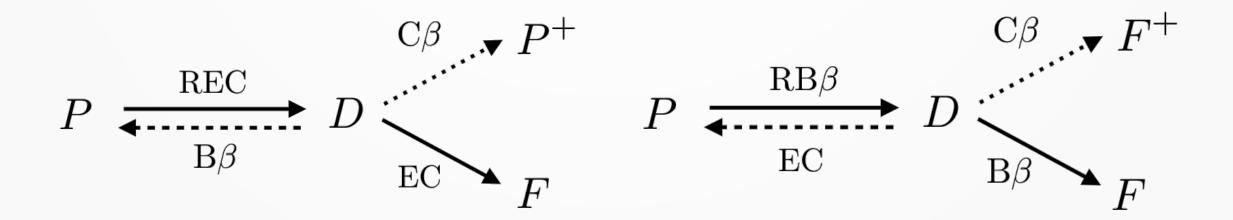
• Overcome this using 3-state systems!

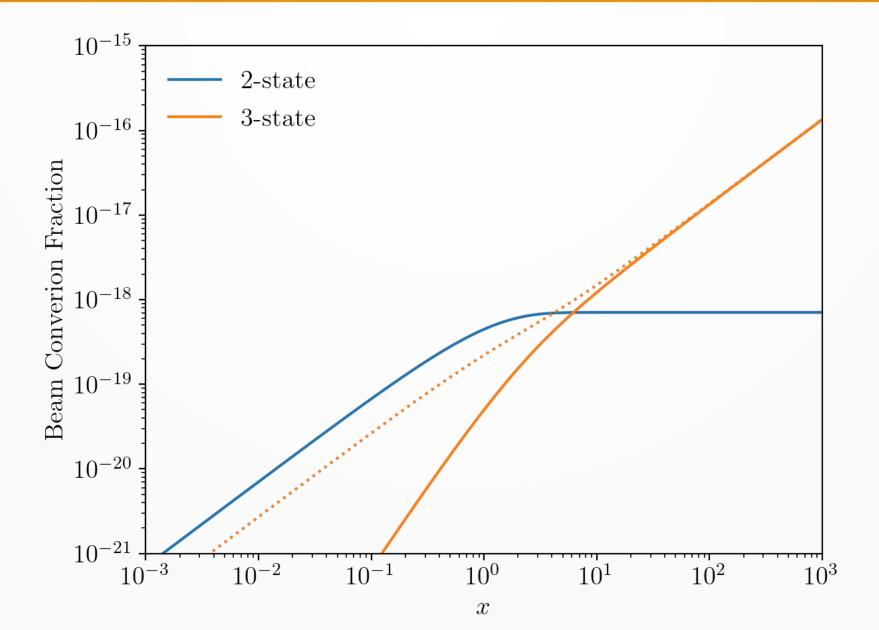
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 $Q \simeq$

$$157 \text{Gd} \xrightarrow{\text{RB}\beta} 157 \text{Tb}$$

$$10 \text{ keV} \qquad \frac{E}{A} \simeq 100 \text{ TeV} \qquad N_D \simeq 10^{-12}$$

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$$V_{D}^{157} \text{Gd} \xrightarrow{\text{RB}\beta} {}^{157} \text{Tb}$$

 $Q \simeq 10 \text{ keV}$ $\frac{E}{A} \simeq 100 \text{ TeV}$ $N_D \simeq 10^{-12}$

• In general:

$$N_D \sim \frac{1}{Q^7} \implies Q \sim 0.1 \,\mathrm{keV}$$

• Where next?

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$$Q \to Q - E^*$$

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• Why use nuclei at all?

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 - Effects far too small to observe currently
- Stodolsky effect:
 - Better than coherent, but requires asymmetry
- Accelerator:
 - Limited by targets

Thank you! Questions?