Detecting the Cosmic Neutrino Background

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What we will cover

• Lecture 1: Introduction to the $C\nu B$

Lecture 2: Direct detection proposals

 Lecture 3: Indirect detection, constraints and future prospects

Contents

- Recap 🧲
- Indirect detection proposals:
 Cosmic ray attentuation
 - Exclusion principle
- Constraints and sensitivities
- What's next?



• PTOLEMY:

$$\nu_e + {}^3\mathrm{H} \to e^- + {}^3\mathrm{He}^+$$



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• Event rate:

$$\Gamma \sim 4 \,\mathrm{y}^{-1} \qquad \Delta \leq 2m_{\nu}$$



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• Uncertainty principle causes issues



$$\nu + X \to \nu + X$$



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• Acceleration:

$$a_{\nu} \simeq 10^{-28} \left(\frac{m_{\nu}}{0.1 \,\mathrm{eV}}\right)^2 \mathrm{cm} \,\mathrm{s}^{-2}$$



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• Future techniques could significantly improve this







• Energy shift giving rise to magnetic field (larger helicity term):

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$$B_{\perp} \simeq 10^{-25} \delta_{\nu} \,\mathrm{T}$$

 $B_{\mathrm{ref}} \simeq 10^{-16} \,\mathrm{T}$

• Asymmetry constrained by BBN:

 $B_{\perp} \lesssim 10^{-24} \,\mathrm{T}$

Recap

• Accelerator:



Recap

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• Need target with small gap

Recap

• Accelerator:



- Need target with small gap
- No suitable targets discovered (yet)

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Recap

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• Limits the proton flux above these energies (see Mauricio's lectures)

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 $\bar{\nu}_{\rm CR} + \nu_{\rm C\nu B} \to Z \to \dots$

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• No neutrinos observed at these energies

• Other processes? [1]

 $\bar{\nu}_{\rm CR} + \nu_{\rm C\nu B} \to X \to \dots$

[1] V. Brdar, P. S. B. Dev, R. Plestid, A. Soni, Phys. Lett. B 833, 137358 (2022)

• Other processes? [1]

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Meson resonances much lighter than EW bosons

[1] V. Brdar, P. S. B. Dev, R. Plestid, A. Soni, Phys. Lett. B 833, 137358 (2022)

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 $\sigma(\bar{\nu}\nu \to X) \propto \operatorname{Br}(X \to \bar{\nu}\nu)$

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• For EW bosons:

 $\operatorname{Br}(Z \to \bar{\nu}\nu) \simeq 0.07$

• Issues?

 $\sigma(\bar{\nu}\nu \to X) \propto \operatorname{Br}(X \to \bar{\nu}\nu)$

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 $\operatorname{Br}(X \to \bar{\nu}\nu) \lesssim 10^{-11}$

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$$\sigma(\bar{\nu}\nu \to X) \propto \operatorname{Br}(X \to \bar{\nu}\nu)$$

• For mesons:

$$\operatorname{Br}(X \to \bar{\nu}\nu) \lesssim 10^{-11}$$

• Need huge density of neutrinos

• What about the final states?

 $X \to \dots$

[2] IceCube, Nature **591** (2021) 7849, 220-224
Cosmic ray attentuation

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• Glashow resonance observed by IceCube [2]:

$$e^- + \bar{\nu}_e \to W^- \to e^- + \bar{\nu}_e$$

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Cosmic ray attentuation

• What about the final states?

 $X \to \dots$

• Glashow resonance observed by IceCube [2]:

$$e^- + \bar{\nu}_e \to W^- \to e^- + \bar{\nu}_e$$

• Perhaps we can do the same for $C\nu B$

[2] IceCube, Nature **591** (2021) 7849, 220-224

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• Neutrinos are fermions

• Neutrino emitting process:

$A \to B + \nu$

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• Exclusion principle:

$$\int d^3 p_\nu \to \int d^3 p_\nu \left(1 - f(p_\nu)\right)$$

Follow a Fermi-Dirac distribution:

$$f(p) = \frac{1}{\exp\left(\frac{p}{T}\right) + 1}$$

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• At very small momenta:

$$f(p) \to \frac{1}{2}$$

• Atomic and nuclear transitions [3]:

 $|e\rangle \rightarrow |g\rangle + \gamma + \nu + \bar{\nu}$

[3] M. Yoshimura, N. Sasao, M. Tanaka, Phys. Rev. D 91, 063516 (2015)

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$$|e\rangle \rightarrow |g\rangle + \gamma + \nu + \bar{\nu}$$

• Maximum photon energy → minimum neutrino energy

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• Atomic and nuclear transitions [3]:

$$|e\rangle \rightarrow |g\rangle + \gamma + \nu + \bar{\nu}$$

• Maximum photon energy → minimum neutrino energy

• Look for suppression at the tail end

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• Realistically?

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- For a given photon energy, lots of combinations of momenta

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- Neutrinos are cold:

$$1 - f(p_{\nu}) \to 1$$

• So what now?

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- Need a restricted phase space, e.g. 2-body decay
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• Chemical potentials?

$$f(p_{\nu}) > \frac{1}{2}$$

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• KATRIN:

 $^{3}\mathrm{H} \rightarrow ^{3}\mathrm{He}^{+} + e^{-} + \bar{\nu}_{e}$

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• Analogous to PTOLEMY



• Current constraints:

$$\sqrt{\sum_{i} |U_{ei}|^2 m_{\nu_i}^2} < 0.45 \,\mathrm{eV}$$

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• Quasi-degenerate regime:

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• Quasi-degenerate regime:

$$m_{\nu_i} \gg \sqrt{\Delta m_{ij}^2} \implies m_{\nu_i} \simeq m_{\nu_i}$$

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 $m_{\nu} < 0.45 \,\mathrm{eV}$

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• Also weakly sensitive to density

 $\sum m_{\nu_i} < 0.07 \,\mathrm{eV}$

[4] DESI, arXiv:2404.03002

• DESI [4]:



• Heavily disfavours inverted hierarchy!

[4] DESI, arXiv:2404.03002

• What if neutrinos are unstable? [5]

 $\sum_{i} m_{\nu_i} < 1 \,\mathrm{eV}$

[5] M. Escudero, J. Lopez-Pavon, N. Rius, S. Sandner, JHEP 12 (2020) 119

• What if neutrinos are unstable? [5]

$$\sum_{i} m_{\nu_i} < 1 \,\mathrm{eV}$$

 Stronger KamLAND bound exists for Majorana neutrinos [6]:

$$\sum_{i} |U_{ei}|^2 m_{\nu_i} < 0.12 \,\mathrm{eV}$$

[5] M. Escudero, J. Lopez-Pavon, N. Rius, S. Sandner, JHEP **12** (2020) 119[6] KamLAND-ZEN, arXiv: 2406.11438

• Also heavily constrains the CvB

• Also heavily constrains the $C\nu B$

• $C\nu B$ has two components – clustered + free streaming:

$$n_{\nu} = n_{\nu,\rm C} + n_{\nu,\rm FS}$$

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• $C\nu B$ has two components – clustered + free streaming:

$$n_{\nu} = n_{\nu,\rm C} + n_{\nu,\rm FS}$$

• Ratio set by mass and temperature

• Clustered component:

$$n_{\nu,\mathrm{C}} \leq \frac{1}{(2\pi)^3} \underbrace{\left(\frac{4}{3}\pi p_f^3\right)}_{V_{\nu}} \underbrace{\left(\frac{4}{3}\pi p_f^3\right)}_{V_{\mathrm{tot}}}$$
• Clustered component:



• Fermi momentum set by escape velocity:

$$p_f \simeq \beta_{\rm esc} m_{\nu}$$

• Clustered component:

$$n_{\nu,\mathrm{C}} \le 12.8 \,\mathrm{cm}^{-3} \left(\frac{m_{\nu}}{0.1 \,\mathrm{eV}}\right)^3$$

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• Unclustered component:

$$n_{\nu,\text{FD}} = \frac{1}{2\pi^2} \int_{p_f}^{\infty} \frac{p^2}{\exp\left(\frac{p}{T} + 1\right)}$$







Constraints



Sensitivity



Sensitivity



Summary

- CvB detection is an almost impossible task:
 - Cold, almost non-interacting neutrinos
- Lots of reasons to look for the $C\nu B$:
 - New interactions, lepton asymmetries, $N_{\rm eff}$
- Lots of ideas to detect the $C\nu B!$
 - None really work...
 - But there are hints!

Thank you!





Děkuju!