

EuCAPT Astroneutrino Theory Workshop 2024 Prague, Czech Republic, Sept. 2024

Large neutrino mass in cosmology: sterile neutrinos as the rescue



KIT – Die Forschungsuniversität in der Helmholtz-Gemeinschaft

HIDDe V Hunting Invisibles: Dark sectors, Dark matter and Neutrinos







Neutrino masses

Neutrino oscillations: • $|m_3^2 - m_1^2| \approx (2.5 \pm 0.03) \times 10^{-3} \,\mathrm{eV}^2$ • $m_2^2 - m_1^2 = (7.42 \pm 0.21) \times 10^{-5} \,\mathrm{eV}^2$

Absolute mass determinations:

- beta-decay spectrum(KATRIN)
- neutrinoless double-beta decay (assuming Majorana neutrinos)
- cosmology





$$m_{\beta} = \sqrt{\sum_{i} |U_{ei}|^2 m_i^2} < 0.45 \text{ eV}$$
$$m_{\beta\beta} = \left|\sum_{i} U_{ei}^2 m_i\right| \lesssim 0.07 \text{ eV}$$
$$\sum_{i} m_i \lesssim 0.1 \text{ eV}$$



Neutrino mass from cosmology

$$\Sigma \equiv \sum_{i=1}^{3} m_i = \begin{cases} m_0 + \sqrt{\Delta m_{21}^2 + m_0^2} + \sqrt{\Delta m_{31}^2 + m_0^2} \\ m_0 + \sqrt{|\Delta m_{32}^2| + m_0^2} + \sqrt{|\Delta m_{32}^2| - \Delta m_{21}^2 + m_0^2} \end{cases}$$











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• minimal values predicted from oscillation data for $m_0 = 0$:

$$\Sigma_{\rm min} = \begin{cases} 98.6 \pm 0.85 \,\mathrm{meV} & (\mathrm{IO}) - 58.5 \pm 0.48 \,\mathrm{meV} & (\mathrm{NO}) - 58.5 \,\mathrm{meV} & (\mathrm{NO})$$









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• Upper bounds from current data:

- $\Sigma m_{\nu} < 0.12 \,\mathrm{eV} \,(95 \,\% \,\mathrm{CL})$ Planck CMB+BAO 2018
- $\Sigma m_{\nu} < 0.072 \,\mathrm{eV} \,(95 \,\% \,\mathrm{CL})$ DESI + CMB 2024









Emerging tension between cosmology and terrestrial data



Gariazzo, Mena, Schwetz, 2302.14159



Th. Schwetz - DESY seminar, 6 June 2024

possible (near-term) future scenarios:





Gariazzo, Mena, Schwetz, 2302.14159



Complementarity between mass determinations from heaven and earth

link between neutrino mass observables in the standard scenario:



fig. by I. Esteban based on NuFit 5.0

neutrinoless double beta decay





- What if terrestrial experiments see a positive signal? How could this be consistent with cosmology?

What if cosmology does not see finite neutrino mass and upper bounds become tighter than the minimal value predicted by neutrino oscillation?



- What if terrestrial experiments see a positive signal? How could this be consistent with cosmology?

A seesaw model for "large" neutrino mass consistent with cosmology, including sterile neutrino dark matter

work with Miguel Escudero, Jorge Terol-Calvo, 2211.01729 Cristina Benso, Drona Vatsyayan, to appear

What if cosmology does not see finite neutrino mass and upper bounds become tighter than the minimal value predicted by neutrino oscillation?



Sterile neutrino at which mass scale?





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Cosmology bounds can be relaxed in non-standard scenarios

- neutrino decay into dark radiation Chacko et al. 1909.05275; 2002.08401; Escudero et al., 2007.04994; Barenboim et al.,2011.01502; Chacko et al. 2112.13862: $\sum m_{\nu} < 0.42 \, \text{eV}$
- time dependent neutrino mass Lorenz et al. 1811.01991; 2102.13618; Esteban, Salvado, 2101.05804; Sen, Smirnov, 2407.02462, 2306.15718; talk by A. Smirnov
- modified momentum distribution Cuoco et al., astro-ph/0502465; Barenboim et al., 1901.04352; Alvey, Sabti, Escudero, 2111.14870
- reduced neutrino density + dark radiation Beacom, Bell, Dodelson, 04; Farzan, Hannestad, 1510.02201; Renk, Stöcker et al., 2009.03286; Escudero, TS, Terol-Calvo, 2211.01729





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Counting the number of neutrino flavours

$N_{\rm eff}$ affects

formation of light elements (BBN), T ~ MeV, t ~ 1 min

 $N_{\rm eff} = 2.78 \pm 0.28 \,(68\% \,{\rm CL})$

CMB decoupling, T ~ eV, t ~ 400 000 yr

$N_{\rm eff} = 2.99 \pm 0.17 \,(68\% \,{\rm CL})$





Relaxing the neutrino mass bound from cosmology

Cosmology is sensitive to:

energy density in non-relativistic neutrinos (late times)

 $\rho_{\nu}^{\text{non.rel.}} \approx n_{\nu} \sum m_{\nu} < 14 \,\text{eV}\,\text{cm}^{-3}$

energy density in relativistic neutrinos (early times, BBN, CMB)

 $N_{\rm eff}^{\rm relat.} = 2.99 \pm 0.17$ 'eff



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relax bound on m_{ν} by reducing neutrino number density

$$\sum m_{\nu} < 0.12 \,\mathrm{eV} \left(\frac{n_{\nu}^{\mathrm{SM}}}{n_{\nu}}\right)$$

introduce "dark radiation" to keep $N_{\rm eff}^{\rm relat.} \approx 3$

$$N_{\rm eff}^{\rm relat.} = N_{\rm eff}^{\nu} + N_{\rm eff}^{\rm DR} \approx 3$$





- after BBN but before CMB decoupling









Relaxed bound from cosmology

relaxing the present bound by converting neutrinos into N_{χ} generations of massless fermions with g_{χ} internal degrees of freedom:

$$\sum m_{\nu} < 0.12 \,\mathrm{eV} \,(1 + g_{\chi} N_{\chi} / 6)$$

need $\gtrsim 10$ massless species for $m_{\nu} \sim 1 \text{ eV}$



Farzan, Hannestad, 1510.02201 Escudero, TS, Terol-Calvo, 2211.01729









- 3 heavy right-handed neutrinos (seesaw)
- new abelian symmetry $U(1)_X$ local or global
- a scalar Φ charged under $U(1)_X$
- \bullet a set of $N_{\!\gamma}$ massless fermions charged under $U(1)_X$

Escudero, TS, Terol-Calvo, 2211.01729







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Yukawa sector

Escudero, TS, Terol-Calvo, 2211.01729

 $-\mathcal{L} = \overline{N_R} Y_{\nu} \ell_L \widetilde{H}^{\dagger} + \frac{1}{2} \overline{N_R} M_R N_R^c + \overline{N_R} Y_{\Phi} \chi_L \Phi + \text{h.c.}$





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Scalar potential $V = \mu_H^2 H^{\dagger} H + \lambda_H (H^{\dagger} H)^2 + \mu_{\Phi}^2 |\Phi|^2 + \lambda_{\Phi} |\Phi|^4 + \lambda_{H\Phi} |\Phi|^2 H^{\dagger} H$







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 $\mathscr{L}_{\text{int}} = g_X Z'_\mu \overline{\chi} \gamma^\mu \chi$ **Gauge interaction**

$$-\frac{1}{2}\overline{N_R}M_RN_R^c + \overline{N_R}Y_{\Phi}\chi_L\Phi + \text{h.c.}$$
$$H^{\dagger}H)^2 + \mu_{\Phi}^2|\Phi|^2 + \lambda_{\Phi}|\Phi|^4 + \lambda_{H\Phi}|\Phi|^2H^{\dagger}H$$







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$$\mathcal{M}_n = \begin{pmatrix} 0 & m_D & 0 \\ m_D^T & M_R & \Lambda \\ 0 & \Lambda^T & 0 \end{pmatrix}$$

Escudero, TS, Terol-Calvo, 2211.01729

 $m_D = \frac{v_{\rm EW}}{\sqrt{2}} Y_{\nu}, \quad \Lambda = \frac{v_{\Phi}}{\sqrt{2}} Y_{\Phi}$

 $\int_{\Phi} \chi_L \Phi + \text{h.c.}$

 $\Lambda \ll m_D \ll M_R$

 $m_{\rm heavy} \approx M_R$ $m_{\rm active} \approx m_D^2/M_R$ $m_{\chi} = 0, \quad \theta_{\nu\chi} \approx \Lambda/m_D$







- 3 heavy right-handed neutrinos (seesaw)
- new abelian symmetry $U(1)_X \rightarrow gauged$
- a scalar Φ charged under $U(1)_X$
- a set of N_{γ} massless fermions charged under $U(1)_X$ $\lambda_{\tau'}^{\chi\chi} = g_X$ $\int_{\Phi} \chi_L \Phi + \text{h.c.}$ $\lambda_{\tau'}^{\chi\nu} = g_X \theta_{\nu\chi}$ $m_{Z'}$ v_{Φ} $\lambda_{\tau'}^{\nu\nu} = g_X \theta_{\nu\gamma}^2$ couplings to neutrinos induced by mixing: $Z' \leftrightarrow \nu \nu l \nu \chi l \chi \chi$

$$-\mathcal{L} = \overline{N_R} Y_{\nu} \ell_L \widetilde{H}^{\dagger} + \frac{1}{2} \overline{N_R} M_R N_R^c + \overline{N_R} Y_{\lambda}^{\dagger}$$
$$\mathscr{L}_{\text{int}} = g_X Z'_{\mu} \overline{\chi} \gamma^{\mu} \chi \qquad g_X = -\frac{1}{2} g_X Z'_{\mu} \overline{\chi} \gamma^{\mu} \chi$$

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$$-\mathcal{L} = \overline{N_R} Y_{\nu} \ell_L \widetilde{H}^{\dagger} + \frac{1}{2} \overline{N_R} M_R N_R^c + \overline{N_R} Y_{\Phi} \chi_L$$
$$\mathscr{L}_{\text{int}} = g_X Z'_{\mu} \overline{\chi} \gamma^{\mu} \chi \qquad g_X = \frac{m_{Z'}}{v_{\Phi}}$$

Escudero, TS, Terol-Calvo, 2211.01729

indep. params for pheno:

- $\int_{\Phi} \chi_L \Phi + \text{h.c.}$

 $m_{\nu}, M_R, \theta_{\nu\chi}$ $v_{\Phi}, m_{Z'}$























• thermalization of the dark sector:

 $\Rightarrow \langle \Gamma(\nu\nu \to Z') \rangle \gtrsim H(T = m_{Z'}/3)$









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- avoid thermalization of the dark sector before BBN: $\langle \Gamma(\nu\nu \to Z') \rangle < H(T = 0.7 \,\text{MeV})$
- free-streaming of neutrinos & dark radiation before/around recombination $\langle \Gamma \rangle < H$ for $z < 10^5$ Taule, Escudero, Garny, 2207.04062







Neutrino mixing with massless states $\theta_{\nu\gamma}$

- avoid thermalization of χ prior neutrino decoupling due to oscillations
- take into account effective potential due to self-interactions











Neutrino mixing with massless states $\theta_{\nu\gamma}$



upper range potentially testable in oscillation experiments $10^{-4} \lesssim \theta_{\nu\chi} \lesssim 10^{-1}$ work in progress [Escudero, Maltoni, Ota, TS]





Constraints on heavy RH neutrinos

 $M_R \lesssim 10^{10} - 10^{14} \,\mathrm{GeV}$

- perturbativity of Yukawa $Y_{\Phi} \overline{N}_R \chi_I \Phi$
- loop-induced Higgs portal $\lambda_{\Phi H} | \Phi |^2 H^{\dagger} H$ remains small to avoid thermalization of Φ prior BBN







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- loop-induced Higgs portal $\lambda_{\Phi H} |\Phi|^2 H^{\dagger} H$ remains small to avoid thermalization of Φ prior BBN

Comment on leptogenesis:

- standard thermal LG works if $N \rightarrow HL$ dominates over $N \to \phi \chi$
- otherwise χ would thermalize and conflict with $N_{
 m eff}$ \Rightarrow require $T_{RH} < M_R$ (allows still for $T_{RH} \gg T_{EW}$)





Signatures in a super nova

• SN cooling arguments for SN1987A exclude

$$3 \times 10^{-7} \, \frac{\text{keV}}{m_{Z'}} \lesssim \lambda_{Z'}^{\nu\nu} \lesssim 10^{-4} \, \frac{\text{keV}}{m_{Z'}} \quad \begin{array}{c} \text{Fiorillo} \\ \text{Vitagli} \\ 2209. \end{array}$$

weaker than BBN constraint $\lambda_{Z'}^{\nu\nu} \lesssim 10^{-7} (\text{keV}/m_{Z'})$

• Future galactic SN at 10 kpc: neutrino signal in HyperK from $Z' \rightarrow \nu \nu$: sensitivity down to

 $\lambda_{Z'}^{\nu\nu} \sim 10^{-9} ({\rm keV}/m_{Z'})$ Akita, Im, Masud, 2206.06852





C. Benso, TS, D. Vatsyayan, to appear

Original model: [Escudero, TS, Terol-Calvo]

- 3 heavy right-handed neutrinos (seesaw) N
- new abelian symmetry $U(1)_X$ local or global
- a scalar Φ charged under $U(1)_X$
- a set of N_{γ} massless fermions χ charged under $U(1)_X$
- add one more heavy RH neutrino N' \Rightarrow one of the χ will also pick up a seesaw induced mass $\rightarrow \psi$



C. Benso, TS, D. Vatsyayan, to appear

neutral fermion mass matrix in the basis $(\chi_L^c, \nu_L^c, \psi_L^c, N', N)$

$$\mathcal{M}_n = egin{pmatrix} 0 & 0 & 0 & \Lambda' & \Lambda \ 0 & 0 & 0 & m_D' & m_D \ 0 & 0 & \kappa' & \kappa \ \Lambda'^T & m_D^T' & \kappa'^T & M' & 0 \ \Lambda^T & m_D^T & \kappa^T & 0 & M \end{pmatrix}$$

assume hierarchies: $M \gg M' \gg m_D \gg \kappa' \gg \Lambda \gg m'_D, \kappa, \Lambda'$. $M' m_D^2 \ll M {\kappa'}^2$

$$m_{\chi} = 0$$

$$m_{\nu} \approx m_D M^{-1} m_D^T$$

$$m_{\psi} \approx \kappa' M'^{-1} \kappa'^T.$$

$$m_{N'} \approx M'$$

$$m_N \approx M.$$



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keV DM candidate

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$$\begin{split} m_{\chi} &= 0 \\ m_{\nu} \approx m_D M^{-1} m_D^T \\ m_{\psi} \approx \kappa' M'^{-1} \kappa'^T \, . \\ m_{N'} \approx M' \end{split}$$
 $m_N \approx M$.

mixing and interactions:

$$\theta_{\nu\psi} = \frac{m'_D}{\kappa'}, \quad \theta_{\chi\psi}$$

$$\mathscr{L}_{\rm int} = g_X Z'_\mu \overline{\psi} \gamma$$



DM production via dark freeze-out



- assume $m_{\psi} < m_{Z'}$
- ψ thermalizes with the dark fluid via $\psi\psi\leftrightarrow Z'$ • DM freeze-out for $T_{DS} \lesssim m_{\psi}$

$$\Omega_{\psi}h^{2} \simeq x_{f} \frac{10^{-10} \text{ GeV}^{-1}}{\langle \sigma v \rangle_{\psi\psi \to \chi\chi}}$$
$$\langle \sigma v \rangle_{\psi\psi \to \chi\chi} = N_{\chi} \frac{g^{4}}{4\pi} \frac{m_{\psi}^{2}}{(m_{Z'}^{2} - 4m_{\psi}^{2})^{2}}$$



Right DM abundance in the relevant parameter region

• DM mass $15 \,\mathrm{keV} \lesssim m_{\psi} \lesssim 100 \,\mathrm{keV}$

• DM stability and X-ray constraints: $\psi \rightarrow \nu \chi \chi, \psi \rightarrow \nu \gamma$ suppressed by $heta_{
u\psi}^2$ require $\theta_{\nu\psi} \lesssim 10^{-8}$





Summary & Outlook

- Relaxing cosmo bound on $\sum m_{\nu}$ requires exciting new physics
- Presented simple seesaw model:
 - large number of massless sterile neutrinos ($N_{\gamma} \gtrsim 10 30$)
 - dark U(1) symmetry with breaking scale between 10 MeV and 10 GeV
 - weakly coupled Z' with mass 1 100 keV with $\lambda_{\tau'}^{\nu\nu} \sim 10^{-9}$
- production due to DS interactions, mixing w active neutrinos can be very small
- possible signatures:
 - galactic SN observations
 - sterile neutrino searches at oscillation experiments [work in progress w M. Escudero, T. Ota, M. Maltoni]

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