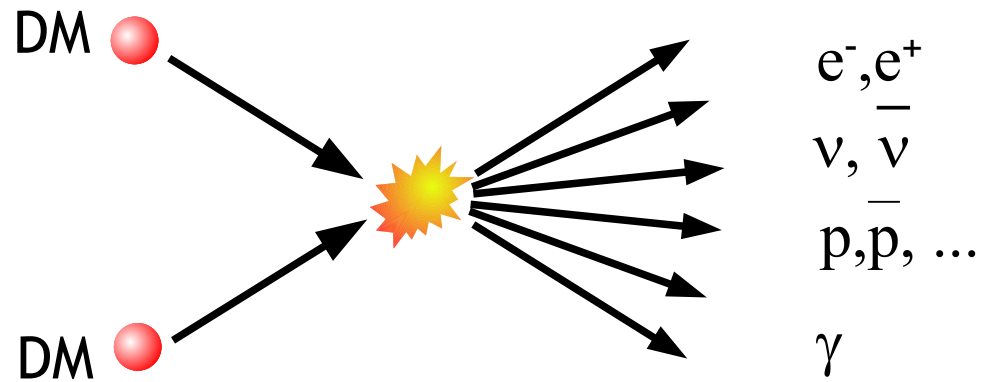


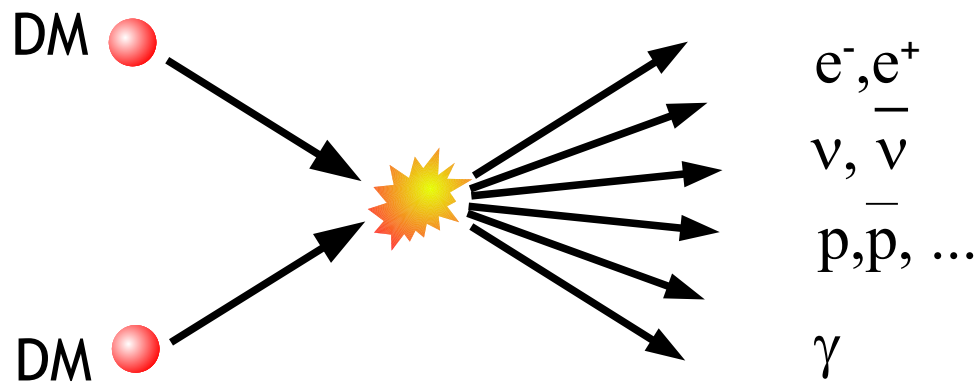
**Probing the DM annihilation  
cross-section using neutrinos**

# Indirect dark matter searches



Dark matter particles annihilate into ordinary particles, such as electrons and positrons, antiprotons, **neutrinos**, photons...

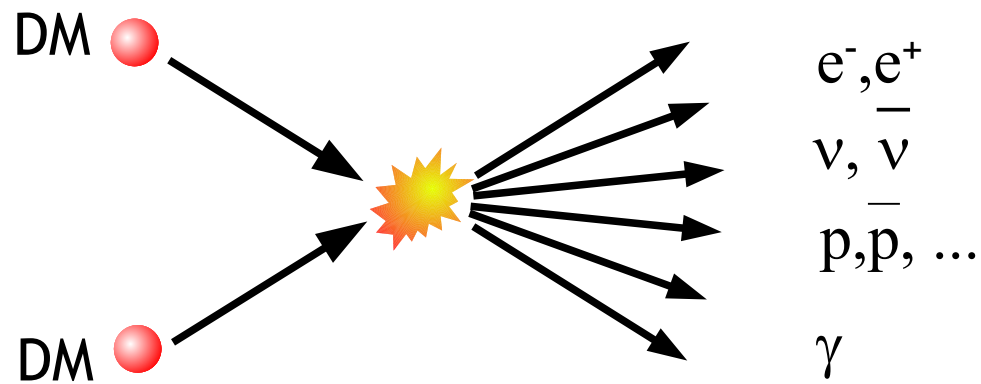
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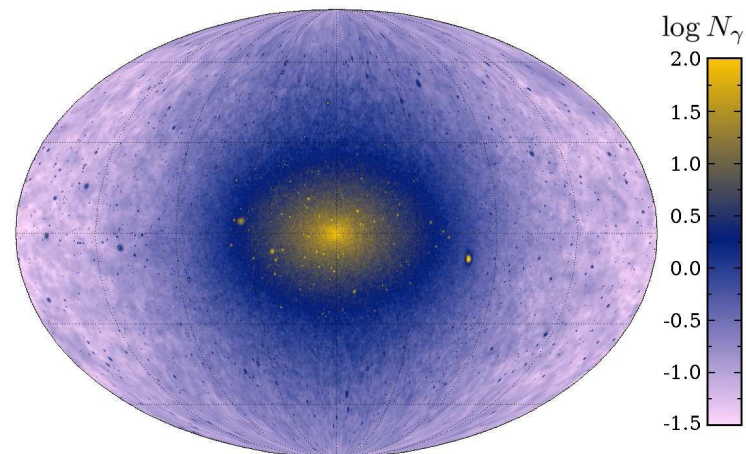
$$\frac{dJ_{\text{halo}}}{dE_\nu} = \frac{1}{4\pi} \underbrace{\left[ \frac{\langle \sigma_{\text{ann}} v \rangle}{2m_{\text{DM}}^2} \sum_f \frac{dN_\nu^f}{dE_\nu} B_f \right]}_{\text{Source term (particle physics)}} \times \underbrace{\int_{\text{l.o.s.}} \rho^2(\vec{l}) d\vec{l}}_{\text{Line-of-sight integral (astrophysics)}}$$

# Indirect dark matter searches



Dark matter particles annihilate into ordinary particles, such as electrons and positrons, antiprotons, **neutrinos**, photons...

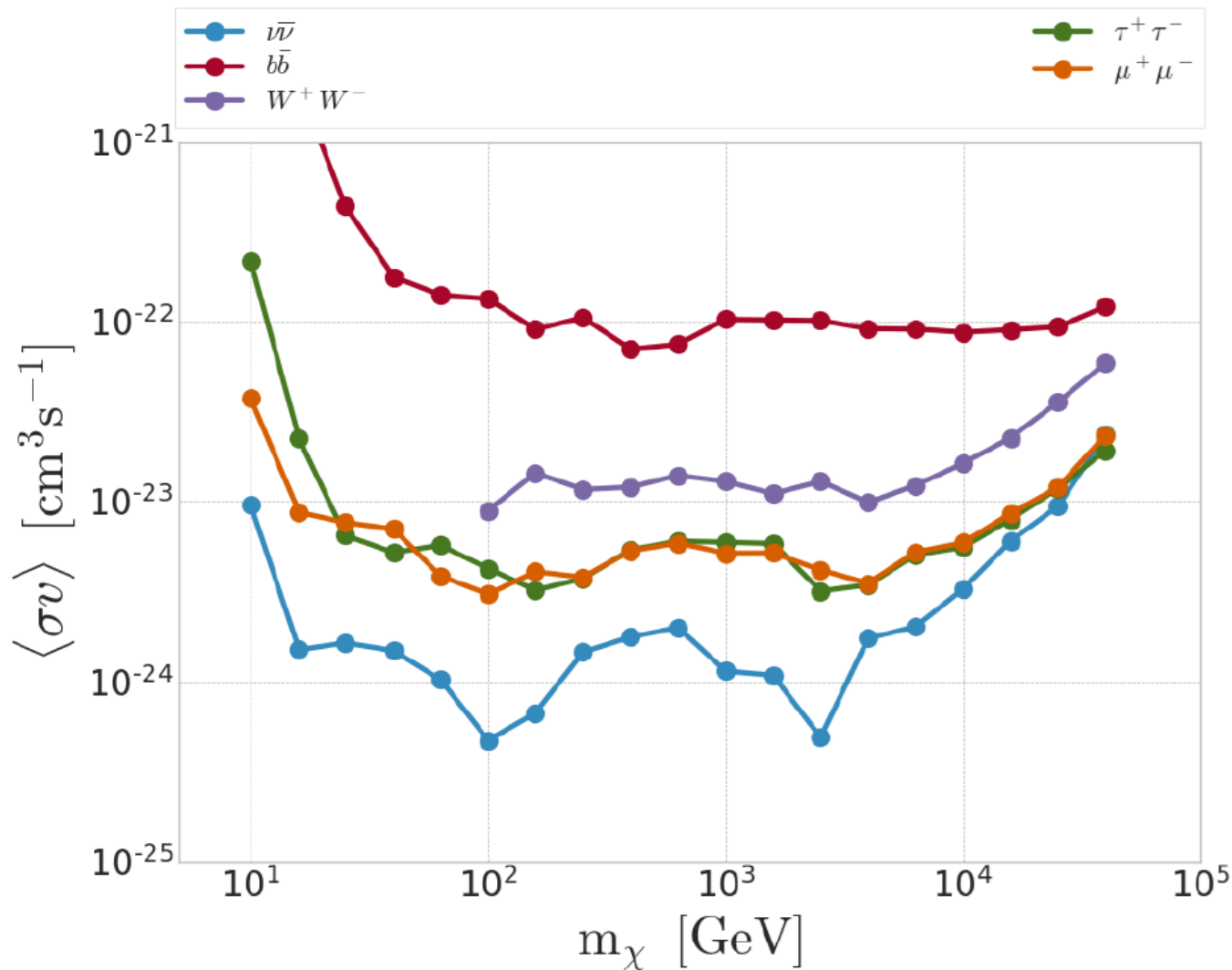
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# Limits on the annihilation cross-section

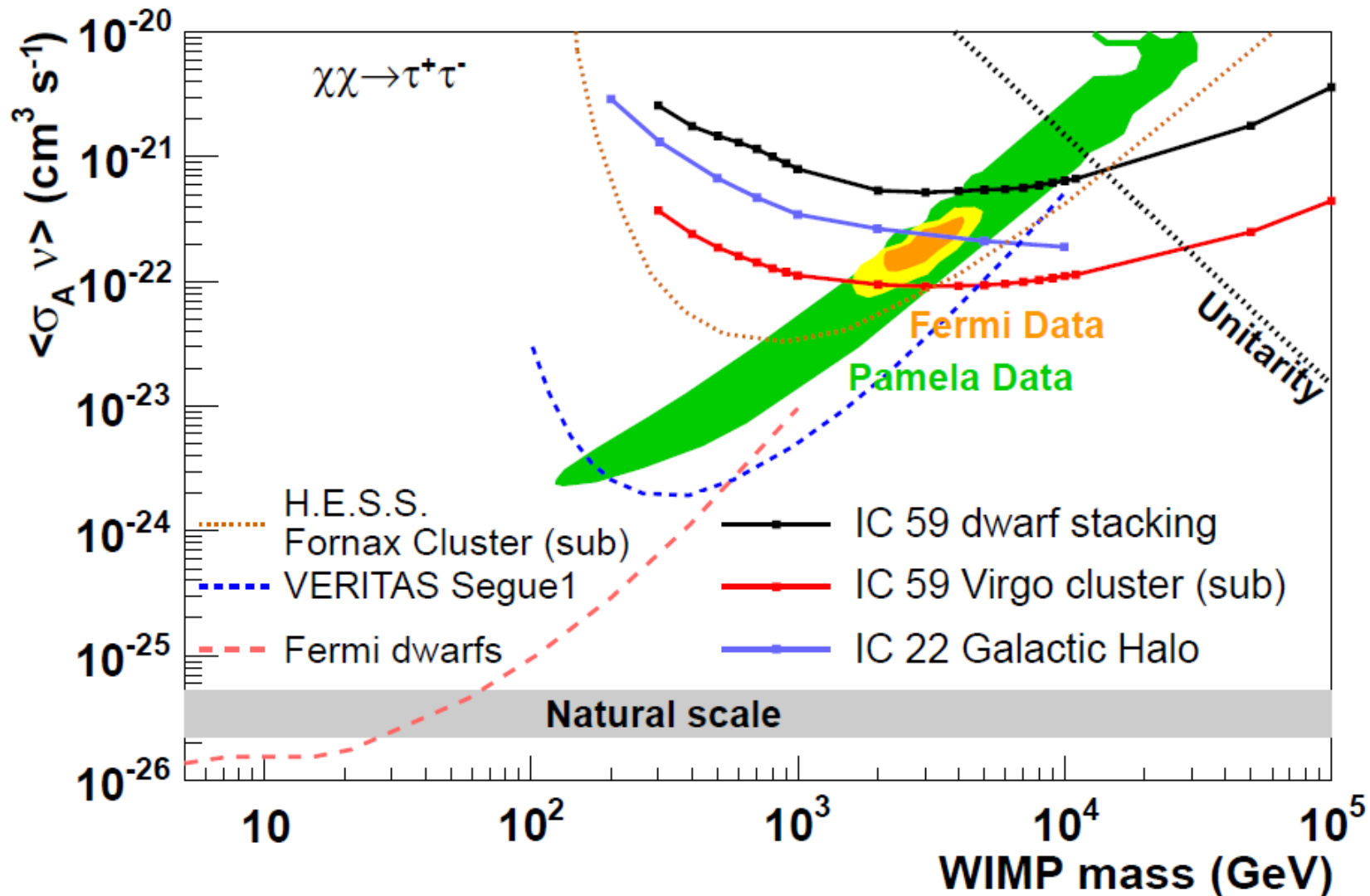
Neutrinos from dark matter annihilations in the Milky Way halo



IceCube 2023

# Limits on the annihilation cross-section

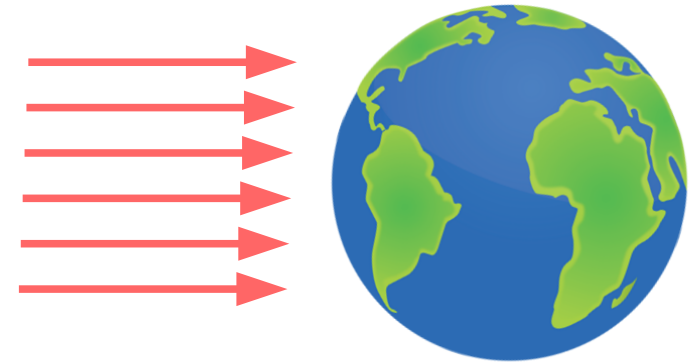
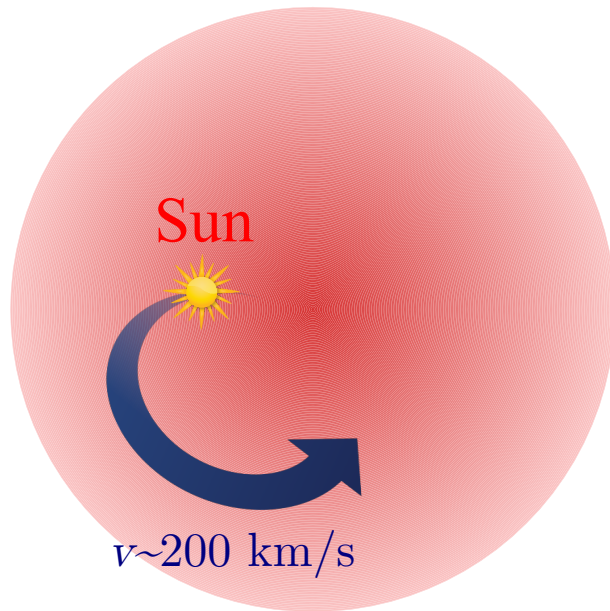
Neutrinos from dark matter annihilations in dwarf galaxies & galaxy clusters.



**Probing the DM–nucleon  
scattering cross–section  
using neutrinos**

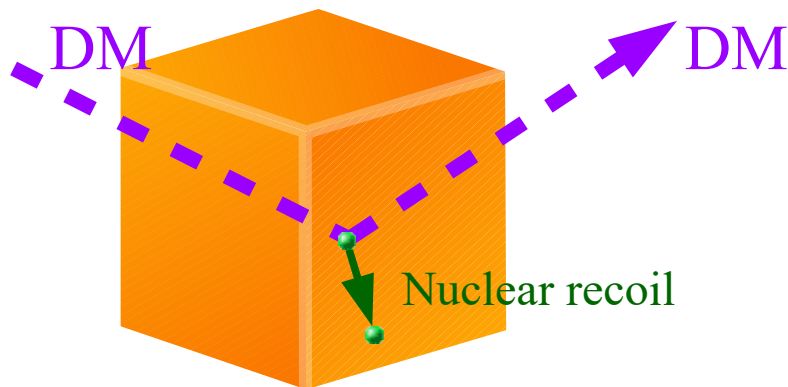
# Standard method: “direct detection” experiments

The Sun (and the Earth) might be moving through a “gas” of dark matter particles.

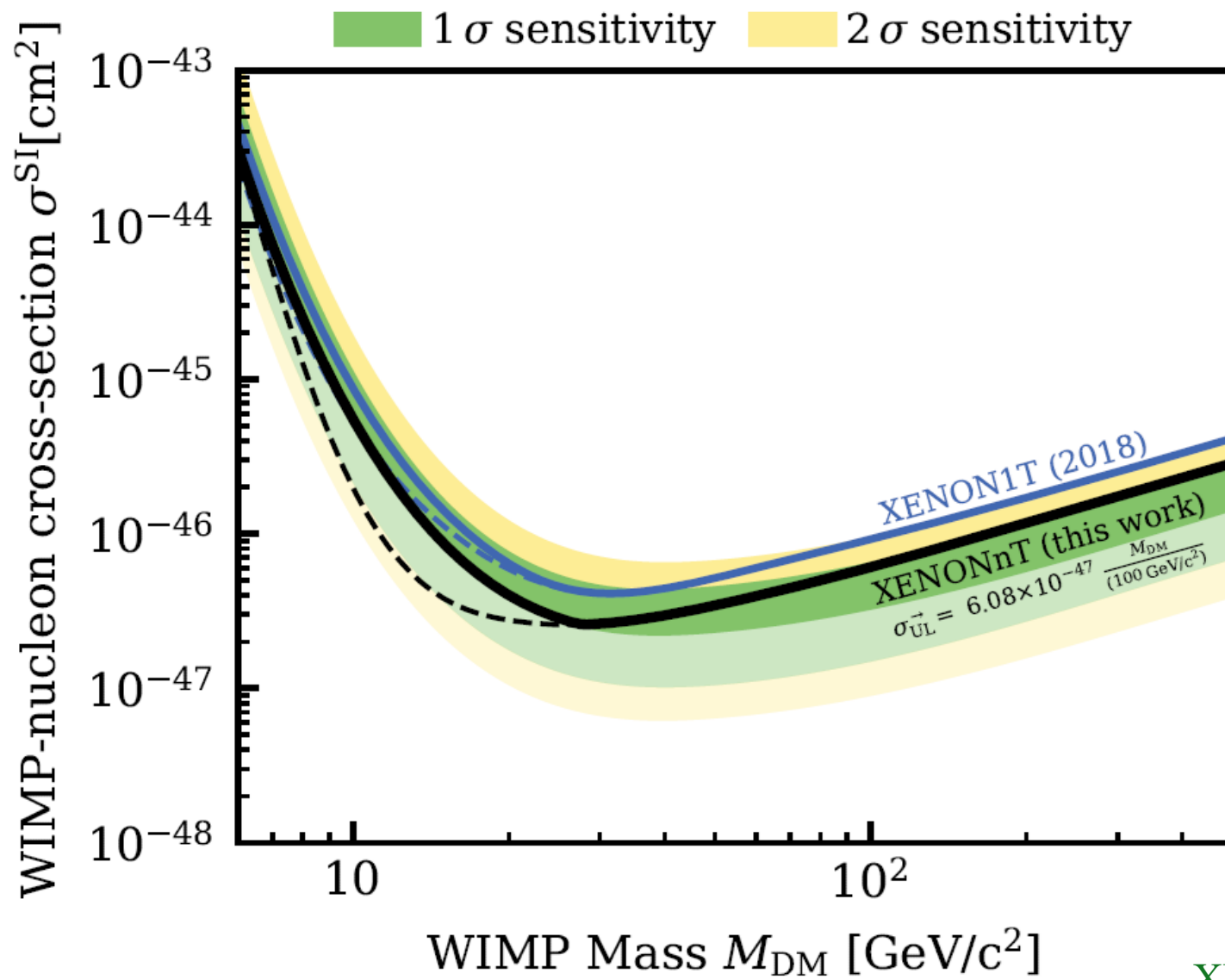


WIMPs  
 $v \sim 200 \text{ km/s}$

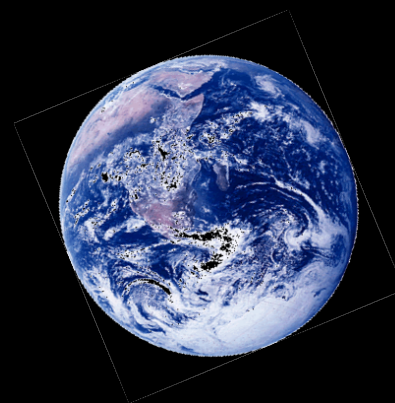
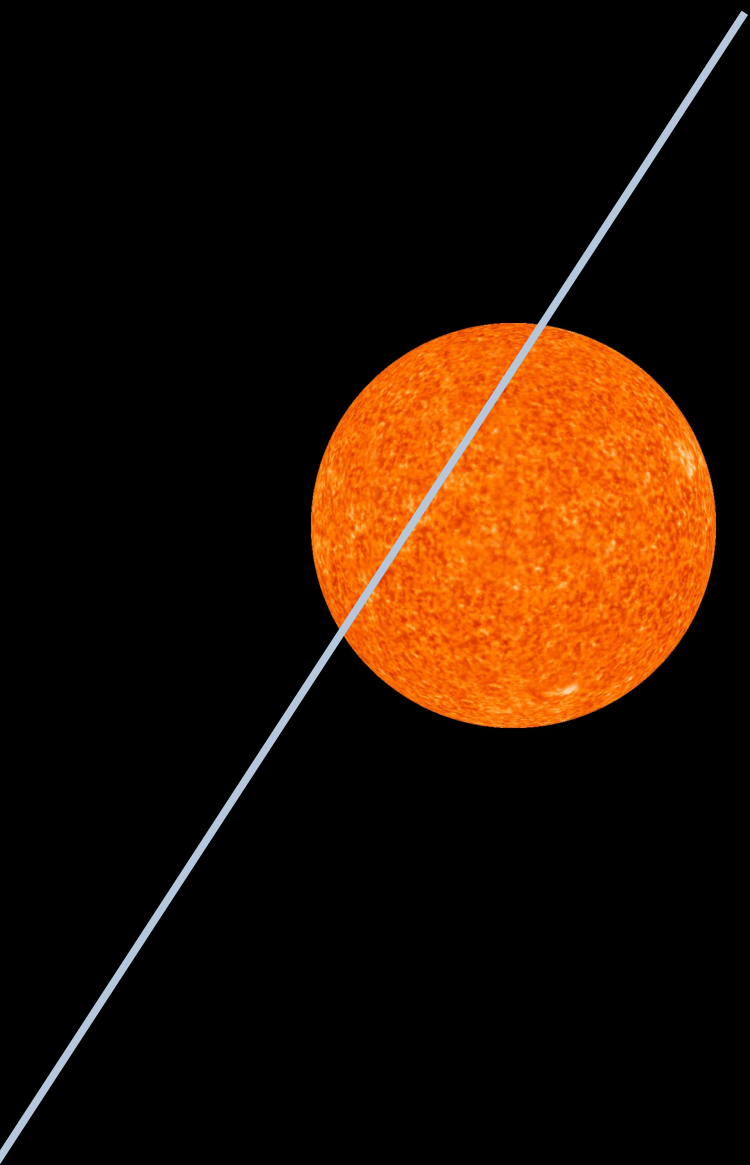
Once in a while a dark matter particle will interact with a nucleus. The nucleus then recoils, producing vibrations, ionizations or scintillation light in the detector.

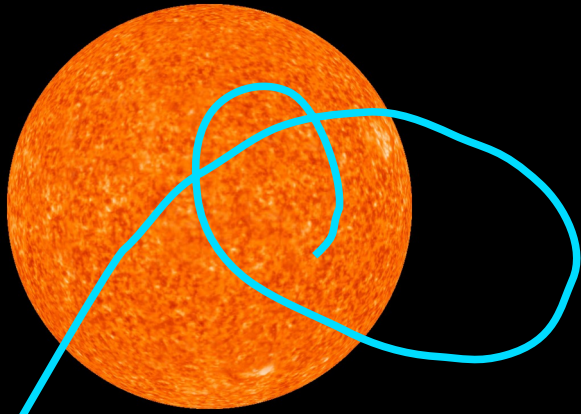


*No significant excess detected so far*



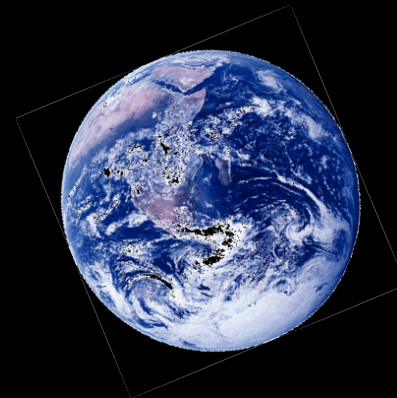
XENONnT'23



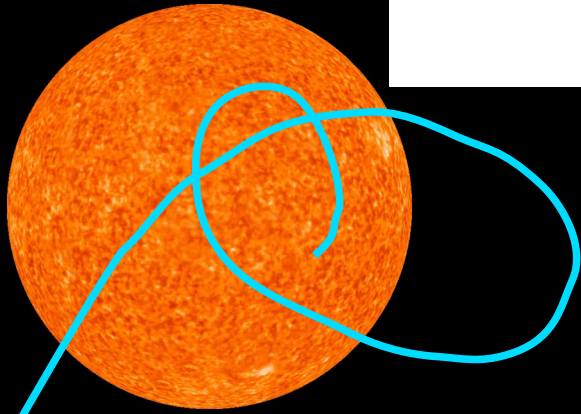


$C_c$  captures per second,

$$C_c \propto (\sigma v)$$

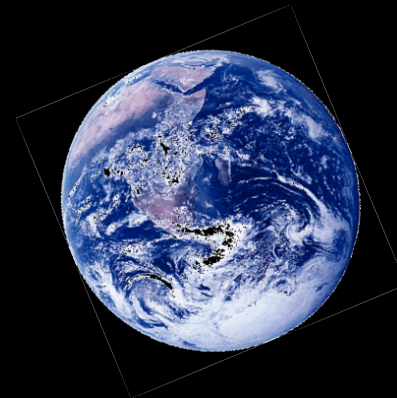


$$C = \int_0^{R_\odot} 4\pi r^2 dr \frac{\rho_{\text{loc}}}{m_{\text{DM}}} \int_{v \leq v_{\text{max}}^{(\text{Sun})}(r)} d^3v \frac{f(\vec{v})}{v} (v^2 + [v_{\text{esc}}(r)]^2) \times \int_{m_{\text{DM}}v^2/2}^{2\mu_A^2(v^2 + [v_{\text{esc}}(r)]^2)/m_A} dE_R \frac{d\sigma}{dE_R}$$



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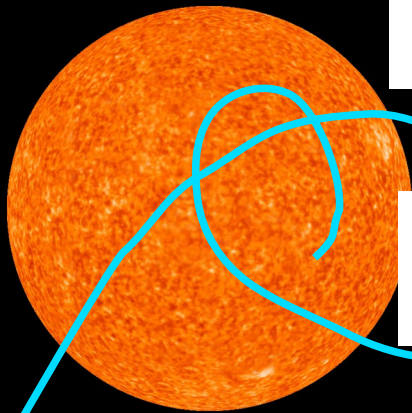
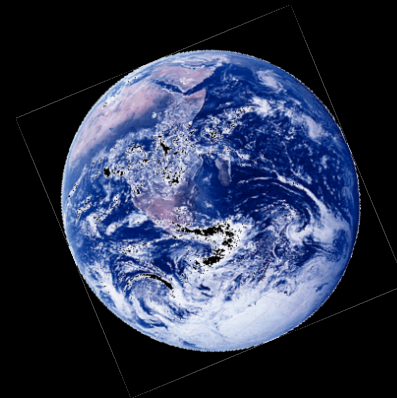


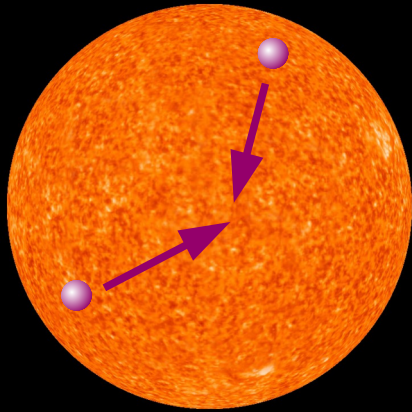
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$$C \simeq 10^{20} \text{ s}^{-1} \left( \frac{\rho_0}{0.3 \text{ GeV/cm}^3} \right) \left( \frac{1 \text{ TeV}}{m_{\text{DM}}} \right)^2 \frac{2.77 \sigma_p^{(\text{SD})} + 4.27 \cdot 10^3 \sigma_p^{(\text{SI})}}{10^{-40} \text{ cm}^2}$$

$C_c$  captures per second,

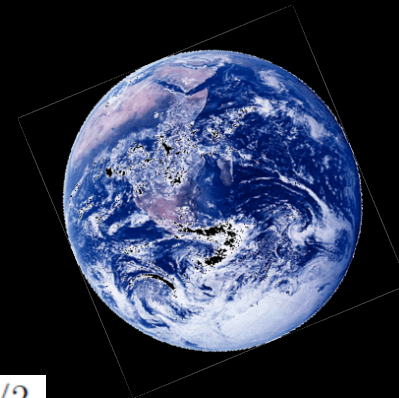
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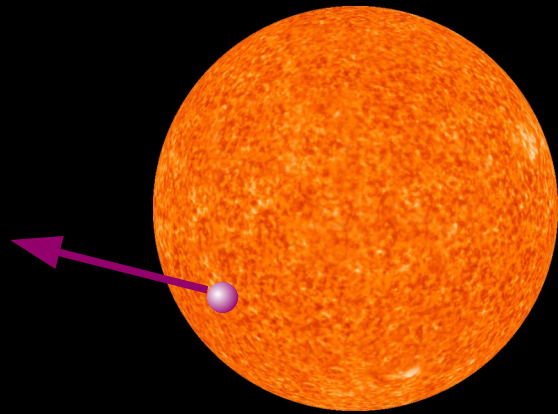


$\Gamma_A$  annihilations per second

$$\Gamma_A = \frac{1}{2} C_A N^2$$

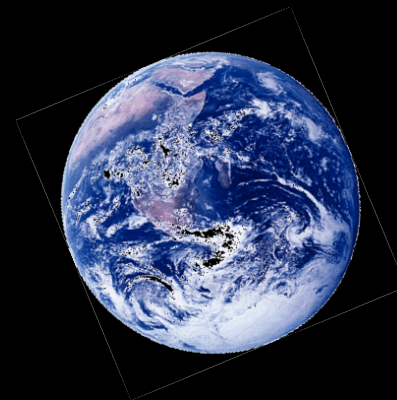


$$C_{A,\odot} \simeq 1.2 \cdot 10^{-52} \text{ s}^{-1} \left( \frac{\langle \sigma v \rangle}{2.2 \cdot 10^{-26} \text{ cm}^3 \text{ s}^{-1}} \right) \left( \frac{m_{\text{DM}}}{\text{TeV}} \right)^{3/2}$$



$C_e$  evaporations (or decays)  
per second

Negligible

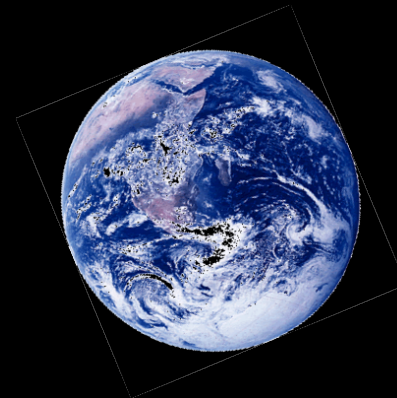
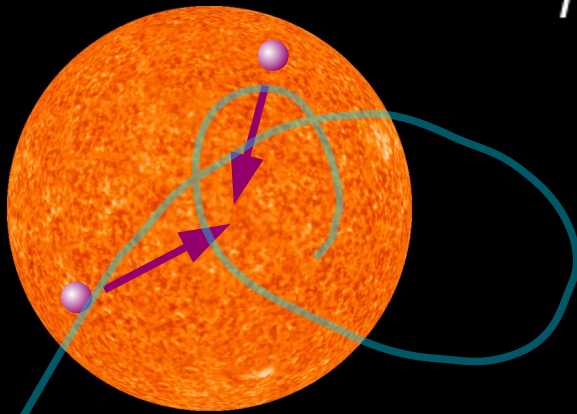


## Evolution of the number of WIMPs

$$\frac{dN}{dt} = C_c - C_a N^2$$

$$N(t) = \sqrt{\frac{C_c}{C_a}} \tanh\left(\frac{t}{\tau}\right), \text{ with } \tau = 1/\sqrt{C_c C_a}$$

$$\Gamma_A = \frac{C_a}{2} N^2 = \frac{C_c}{2} \tanh^2\left(\frac{t}{\tau}\right)$$

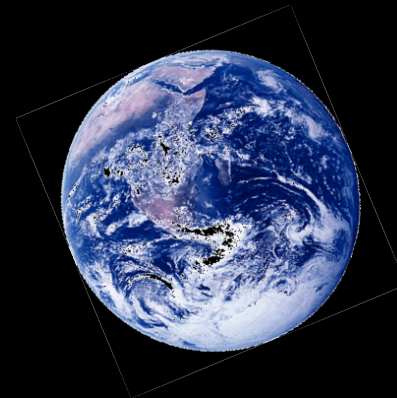
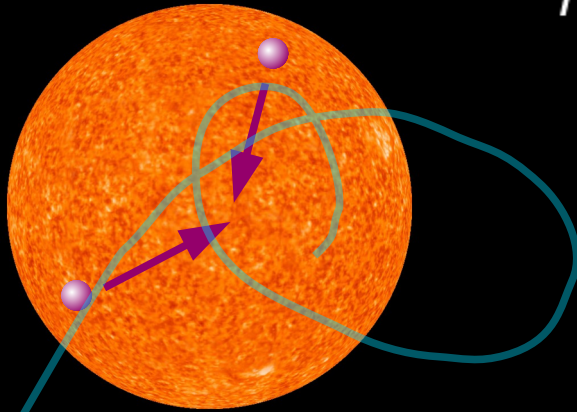


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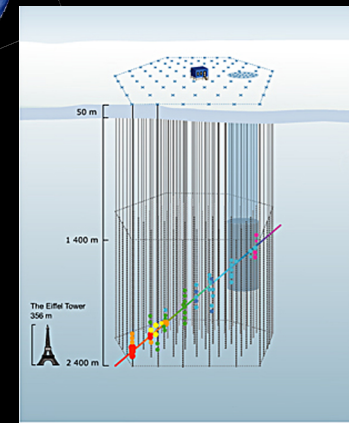
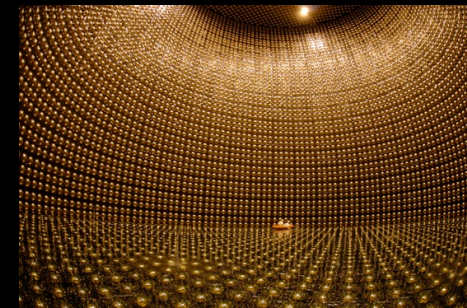
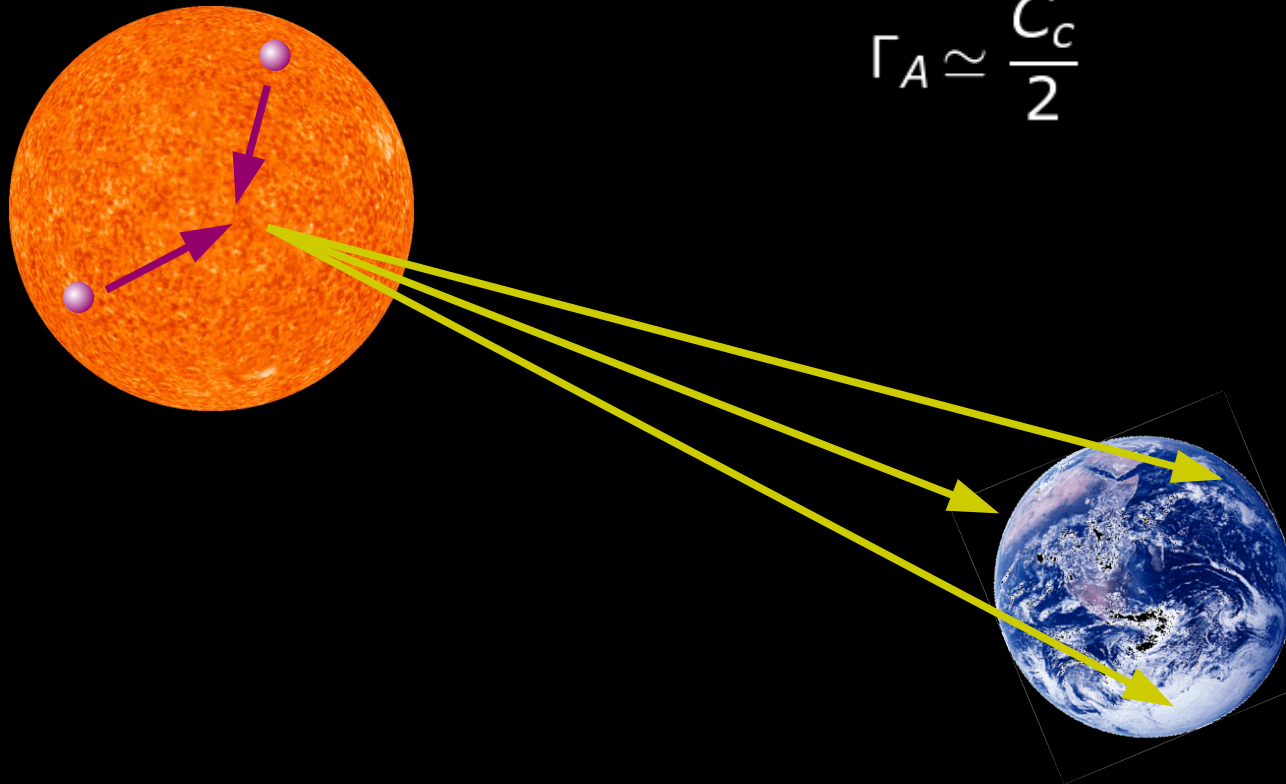
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$$\Gamma_A = \frac{C_a}{2} N^2 \simeq \frac{C_c}{2}, \text{ if } t \gg \tau$$



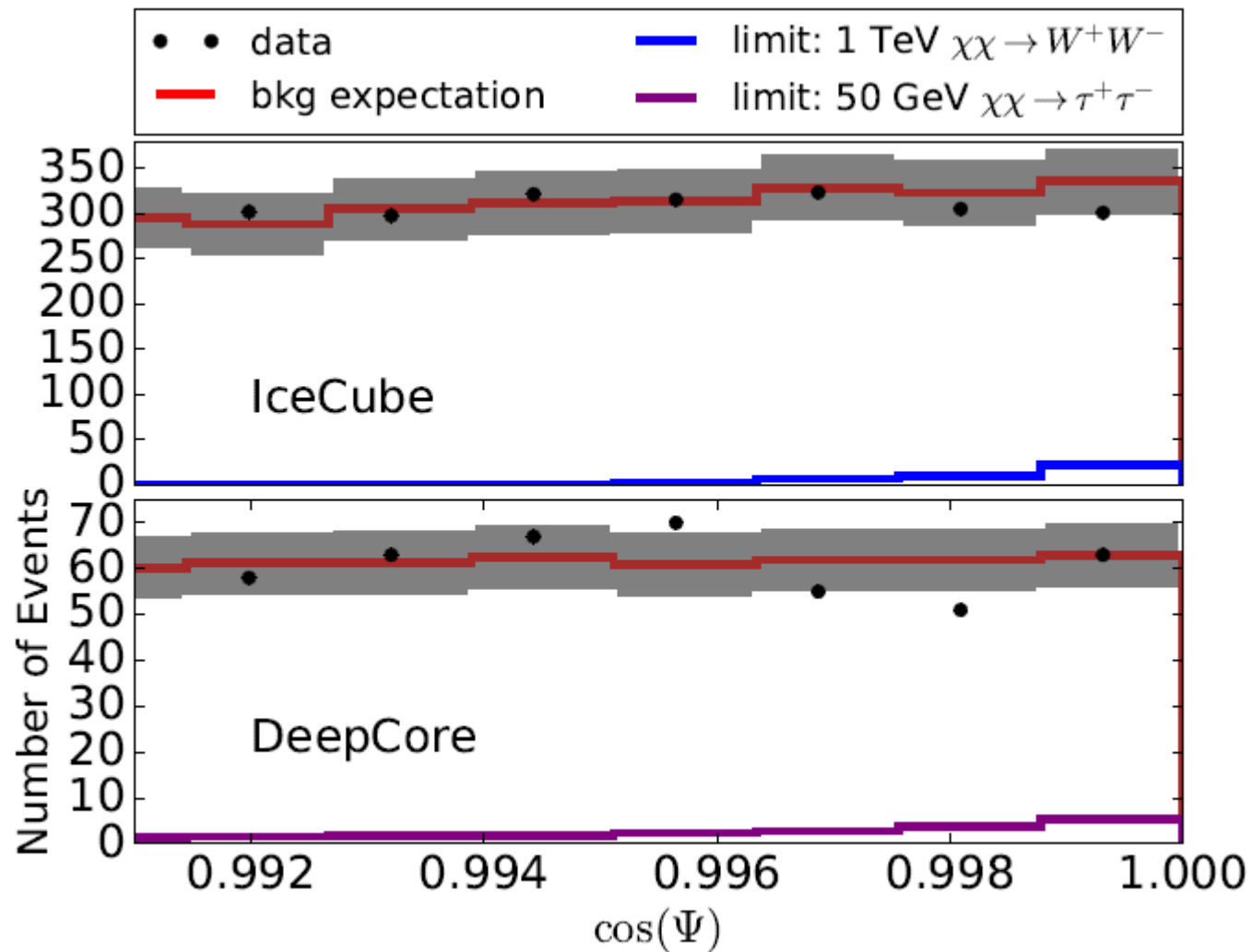
- The annihilation of dark matter particles produces Standard Model particles.
- Most of them are absorbed in their propagation inside the Sun.
- Only neutrinos can escape.
- The neutrino flux measures in fact the *capture rate*

$$\Gamma_A \simeq \frac{C_C}{2}$$



# Neutrinos from annihilations in the Sun

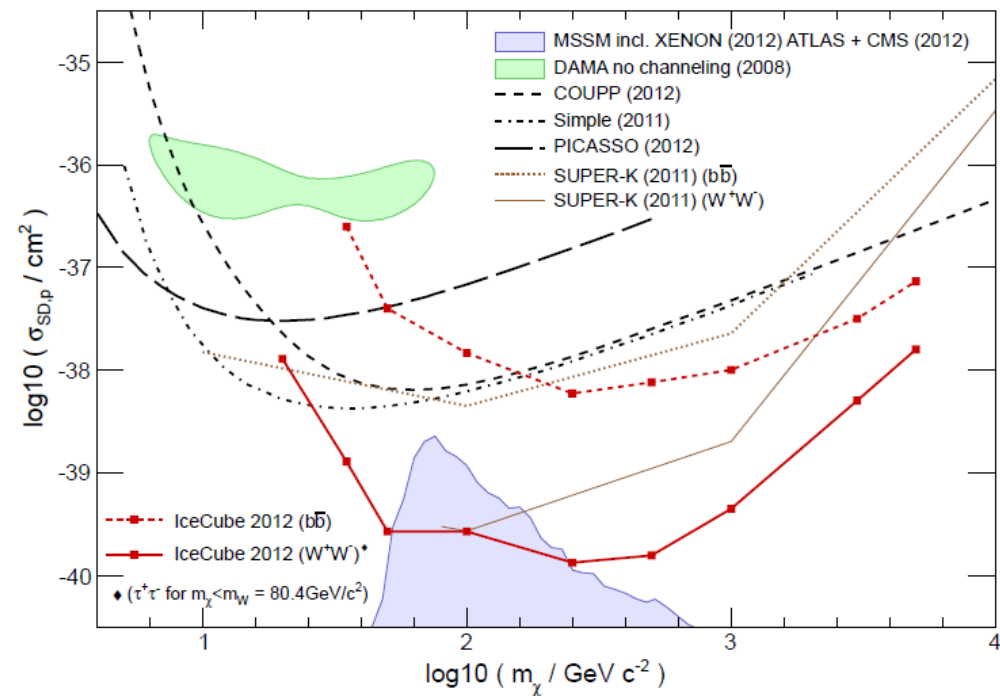
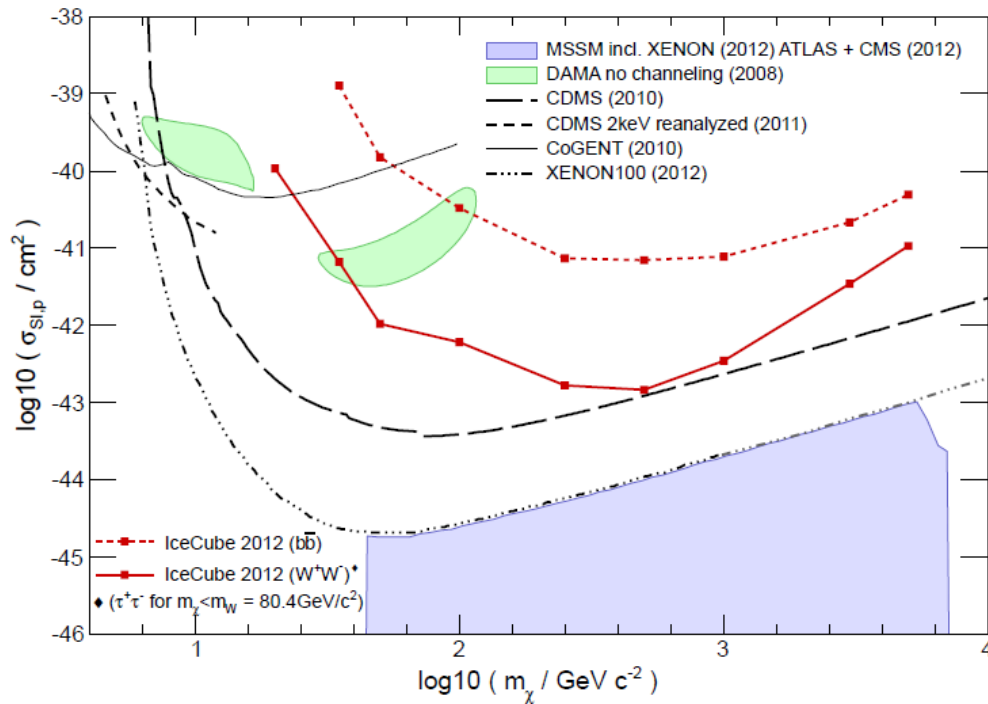
Observations consistent with the background-only hypothesis





# Limits on the scattering cross-section

Limits on the spin-dependent and spin-independent scattering cross section of dark matter particles with protons.

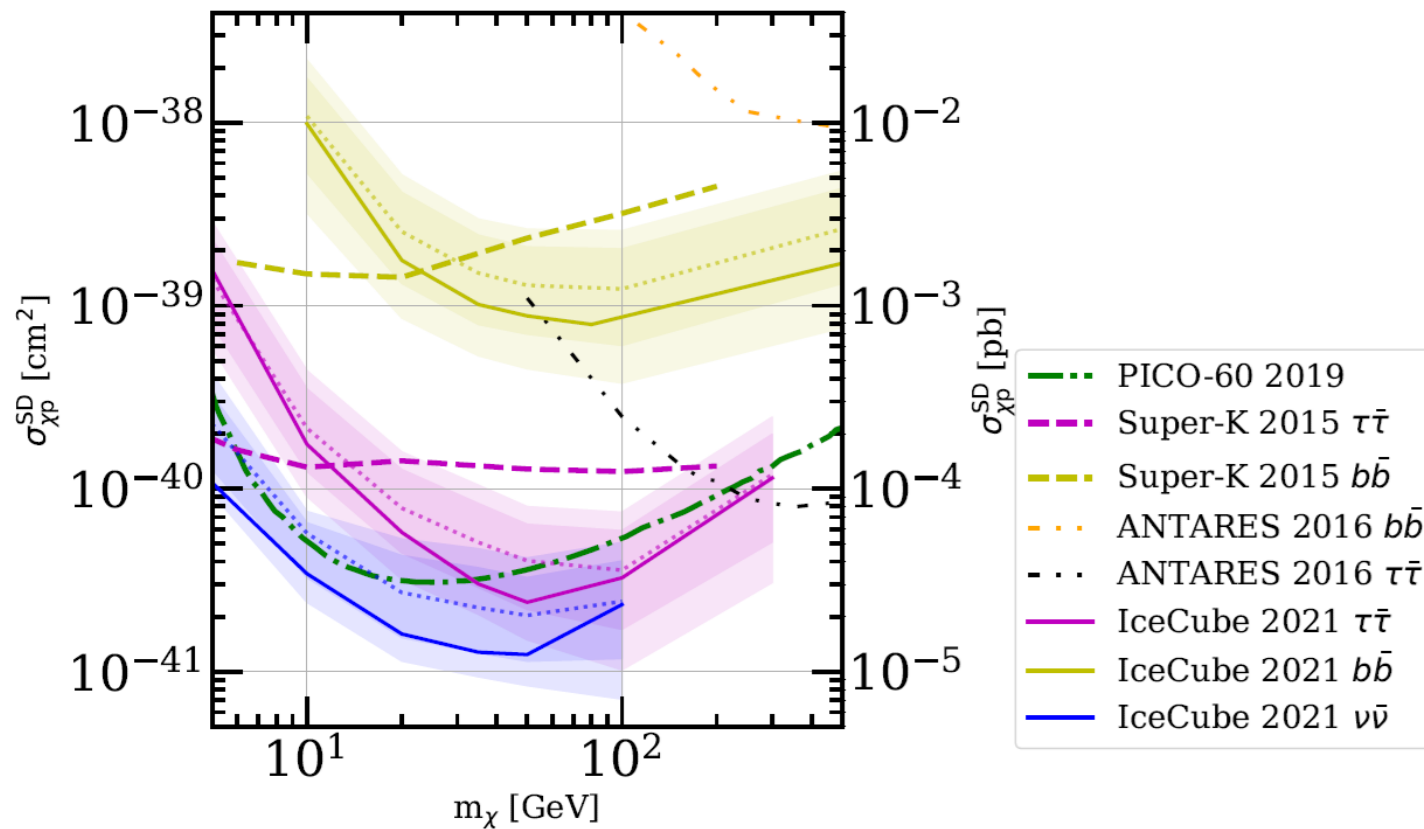


IceCube Collaboration  
arXiv:1212.4097



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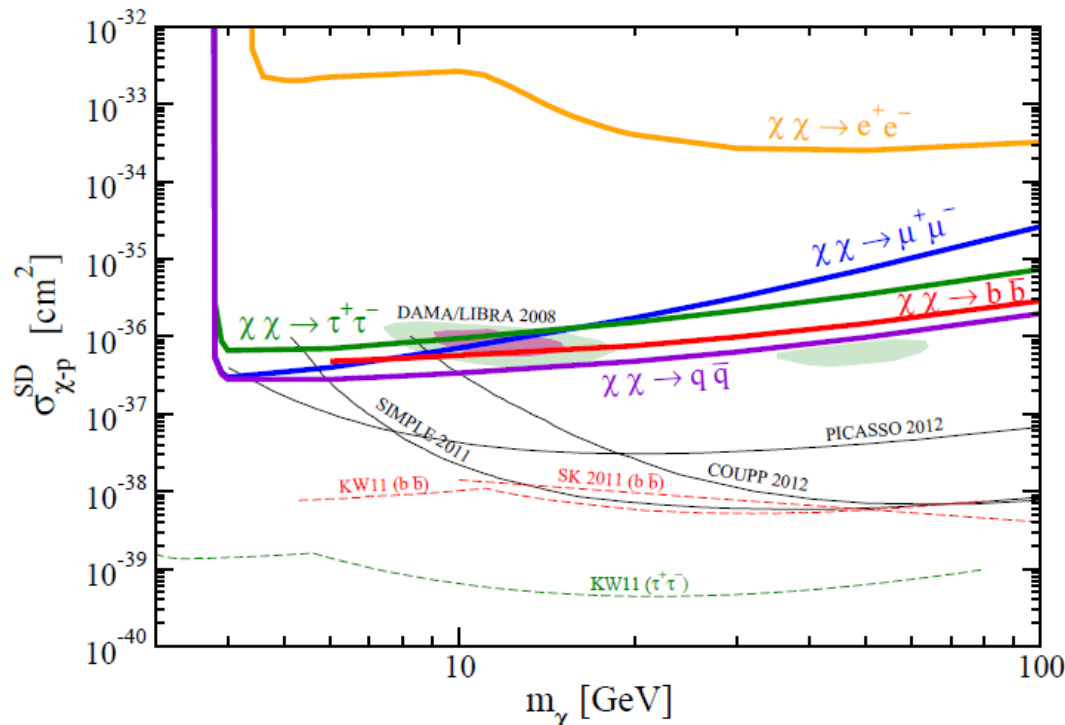
IceCube Collaboration'21

# Limits on the scattering cross-section

## Limits on annihilations channels into light fermions.

The annihilation  $\text{DM DM} \rightarrow q \bar{q}$ , with  $q$  a light quark, **does not produce high energy neutrinos**. The light quark produces pions which are quickly stopped in the solar interior before decaying. This annihilation channel produces only MeV neutrinos.

The MeV neutrinos could be detected at Super-Kamiokande



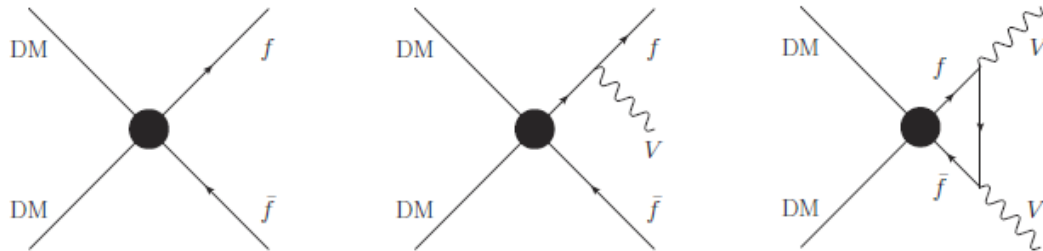
Bernal, Martín-Albo, Palomares-Ruiz, arXiv:1208.0834

See also Rott, Siegal-Gaskins, Beacom, arXiv:1208.0827

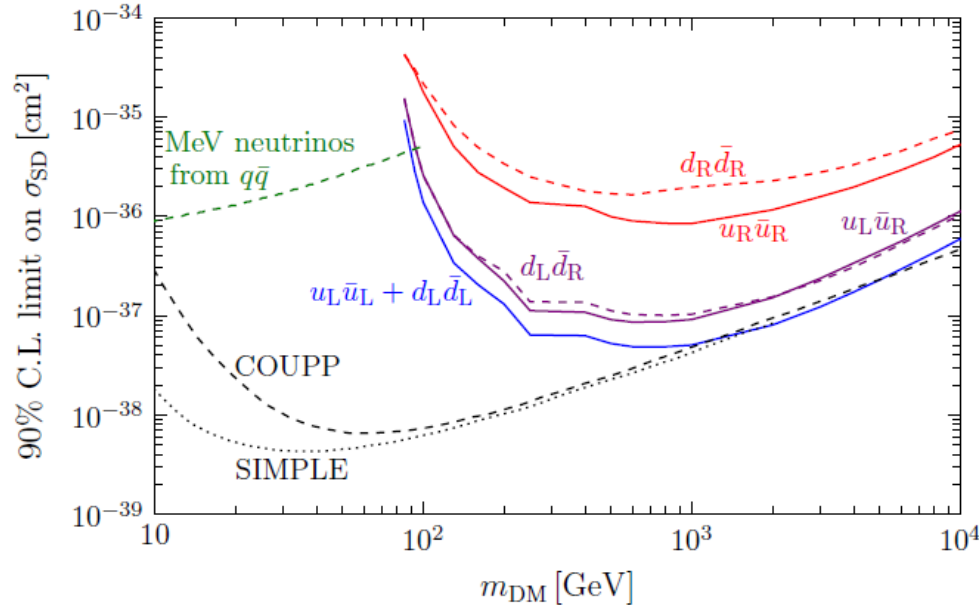
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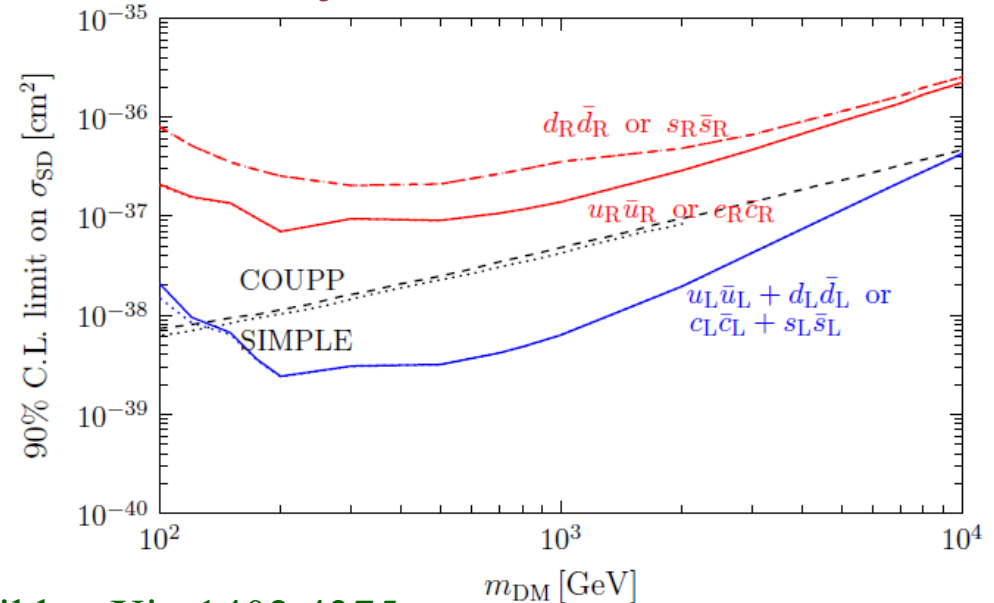
The higher order annihilations  $DM DM \rightarrow q \bar{q} Z$  or  $DM DM \rightarrow Z Z$  do produce high energy neutrinos via the decay of the Z boson



Dirac dark matter



Majorana/scalar dark matter

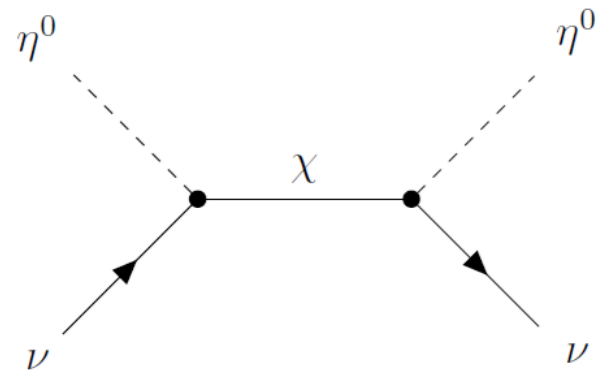
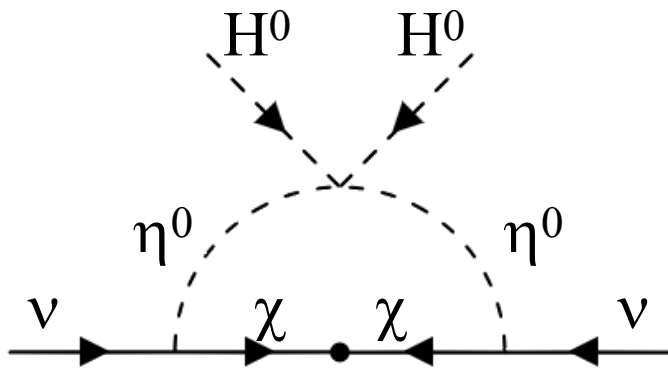


**Probing the DM–neutrino  
scattering cross–section  
using neutrinos**

# DM-neutrino scattering

Much less studied. But natural in some models

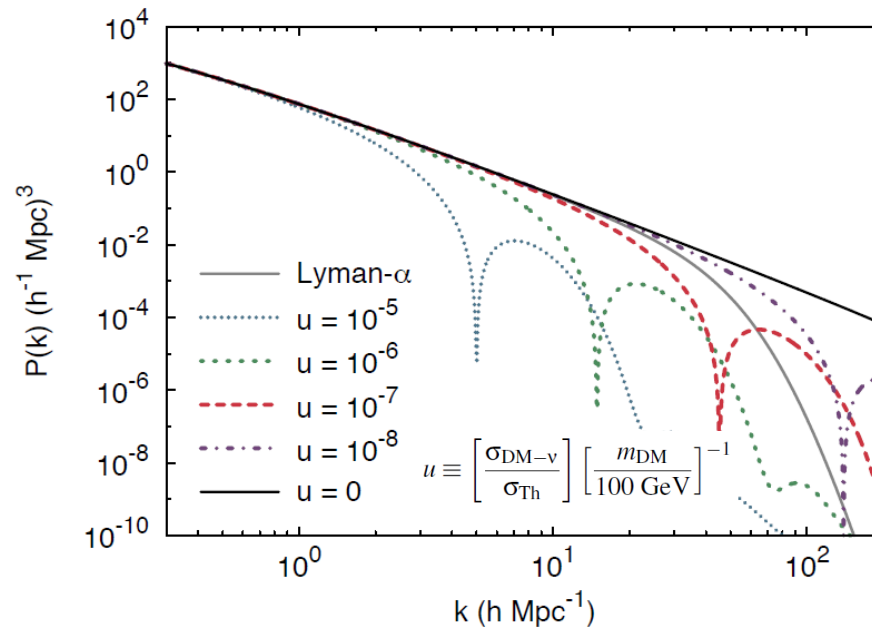
◆ “Scotogenic” neutrino mass models: neutrino masses arise from interactions with the DM.



# DM-neutrino scattering

- DM scattering with neutrinos?

## Constraints from cosmology



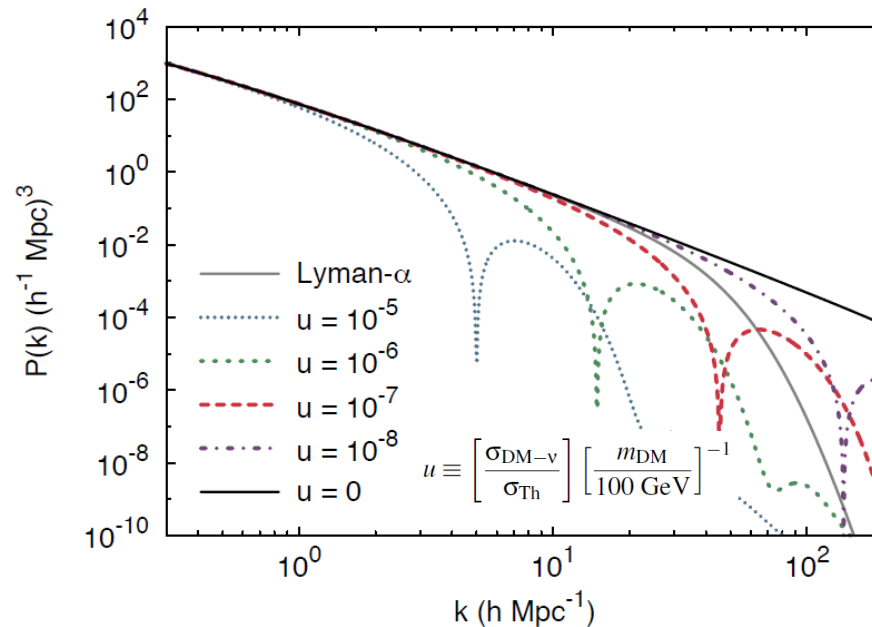
Wilkinson, Boehm, Lesgourgues'13,14

$$\sigma_{\text{DM}-\nu} \lesssim 10^{-33} (m_{\text{DM}}/\text{GeV}) \text{ cm}^2$$

# DM-neutrino scattering

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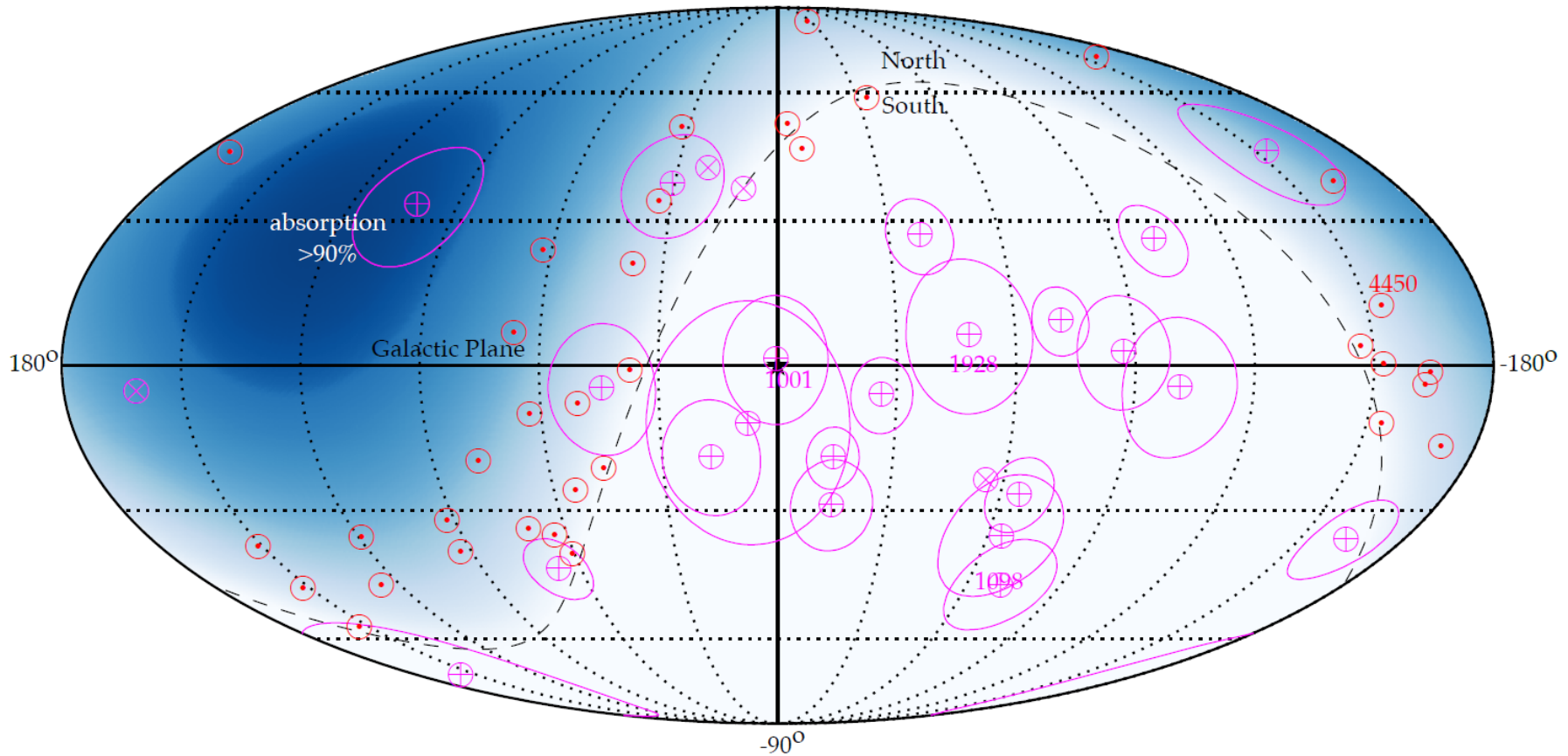
Wilkinson, Boehm, Lesgourgues'13,14

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Other constraints from Astroparticle Physics?

# A new era in neutrino astronomy

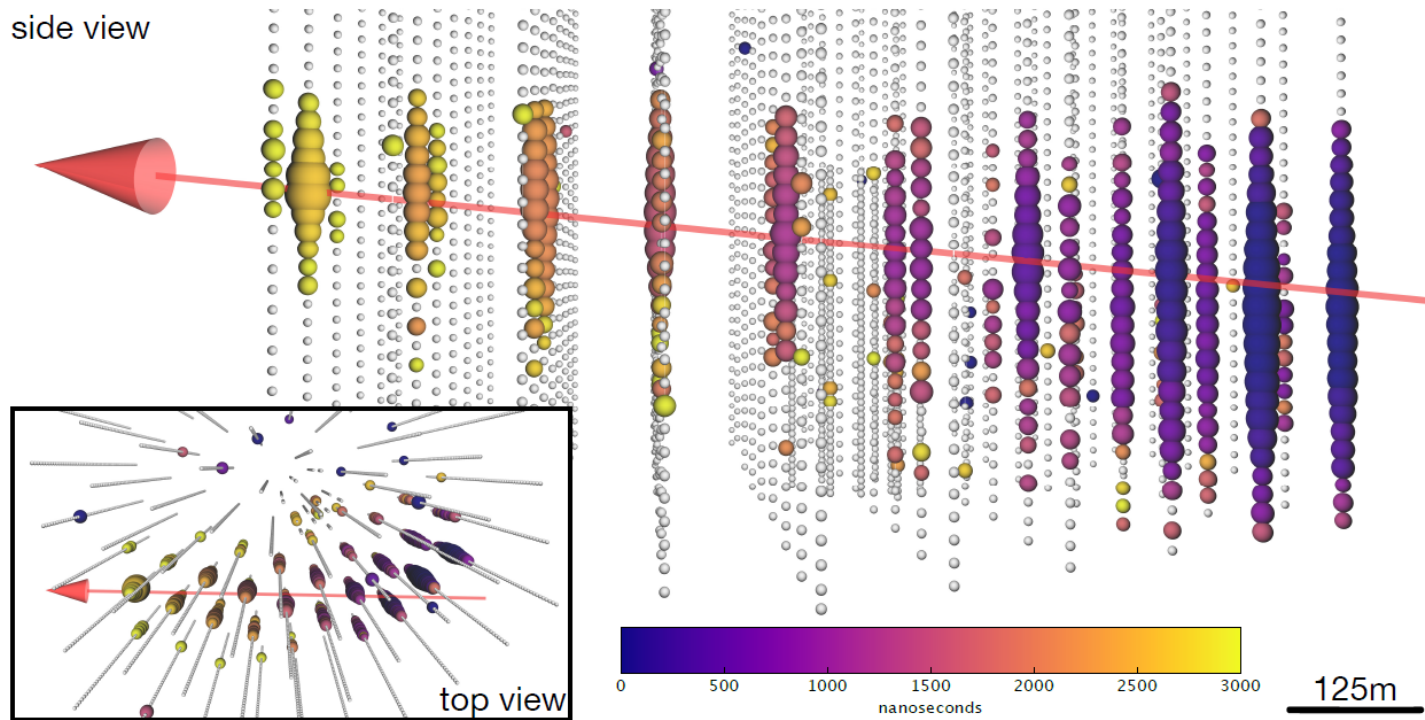
Arrival directions of most energetic neutrino events (HESE 6yr (magenta) &  $\nu_\mu + \bar{\nu}_\mu$  8yr (red))





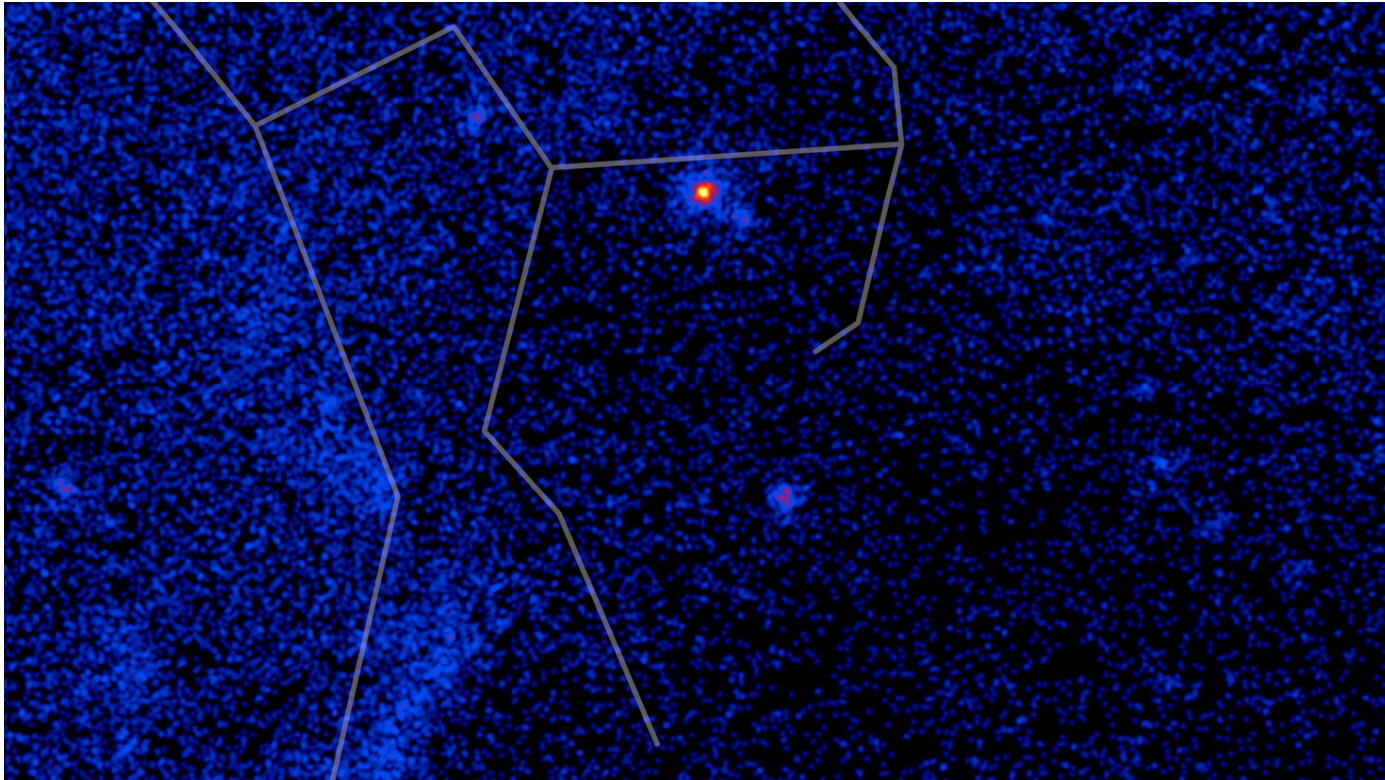
# A new era in neutrino astronomy

- 22 September 2017: detection of the neutrino event IceCube-170922A, with an energy of 290 TeV



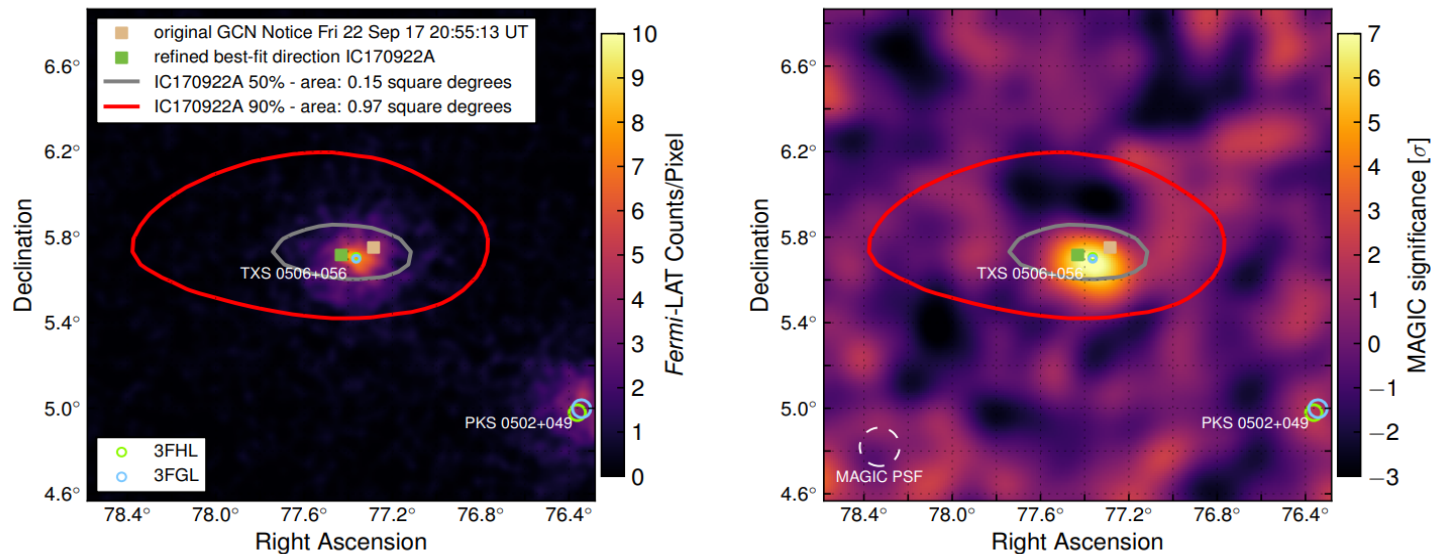
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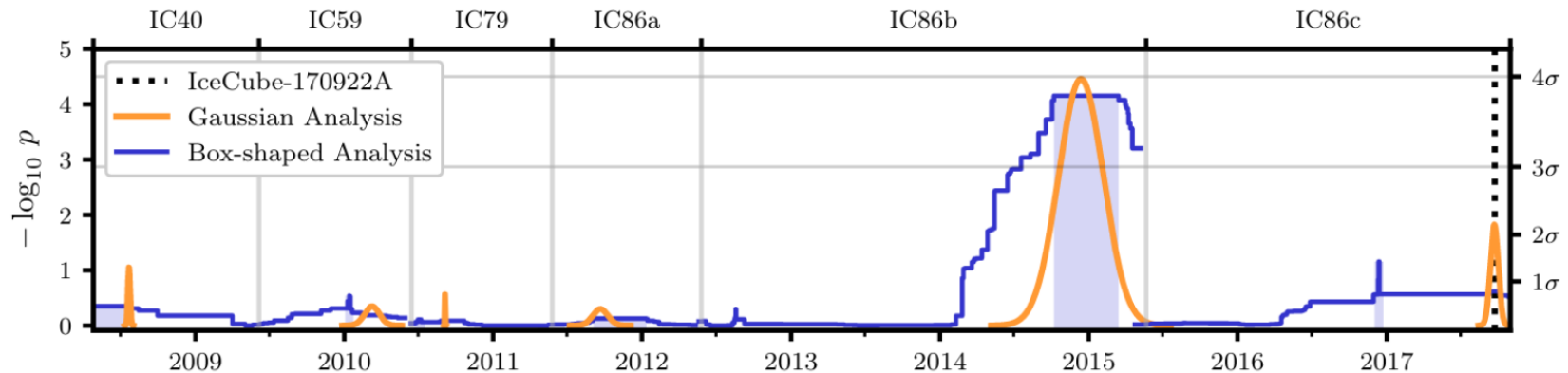
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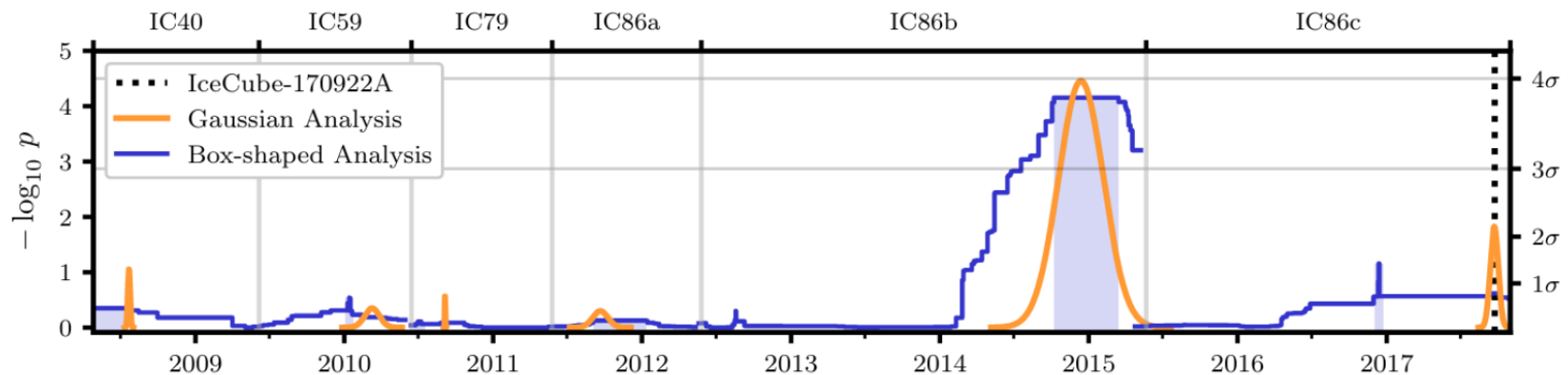
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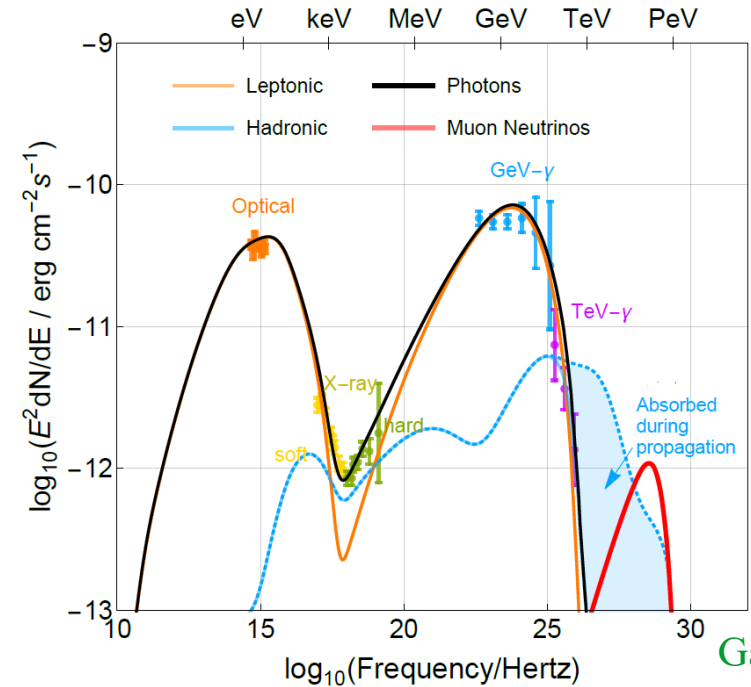
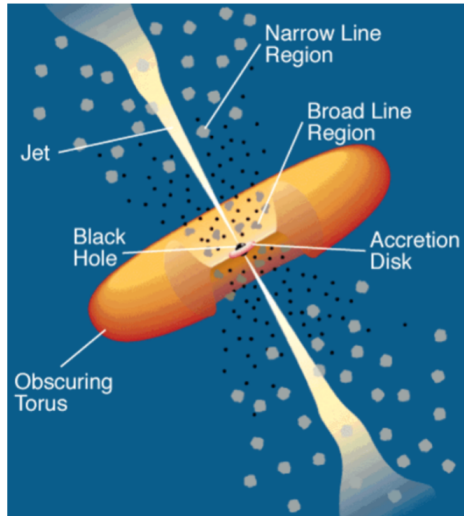


First identified source of high energy astrophysical neutrinos

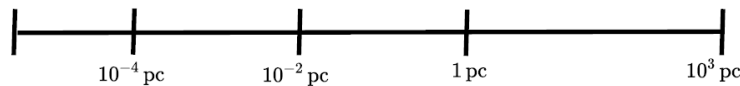
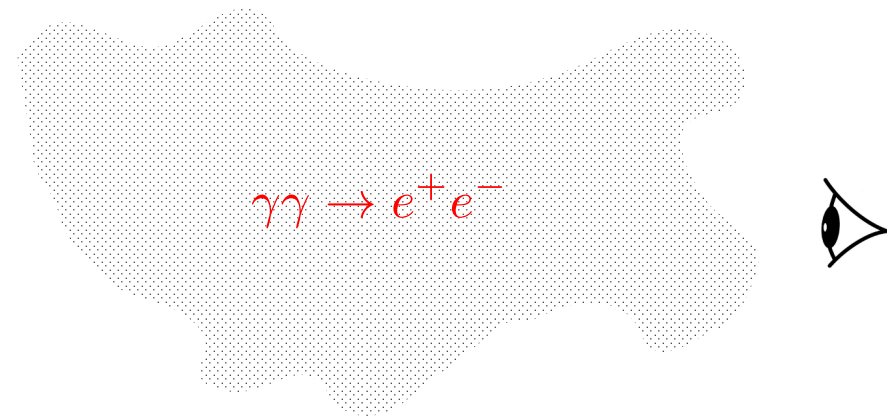
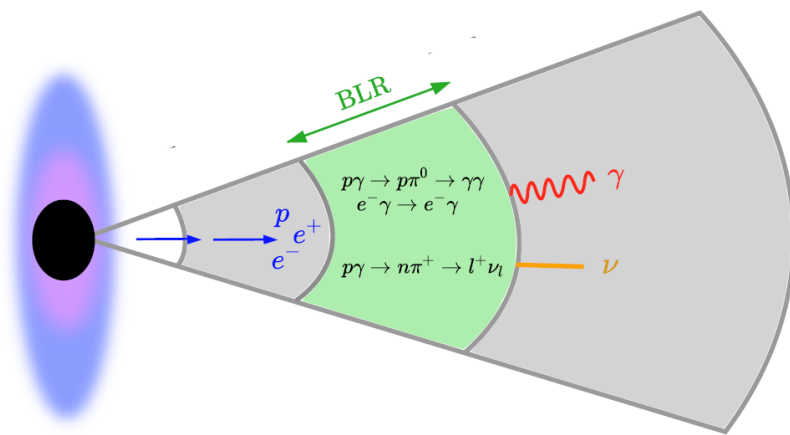


# Photon and neutrino production in blazars

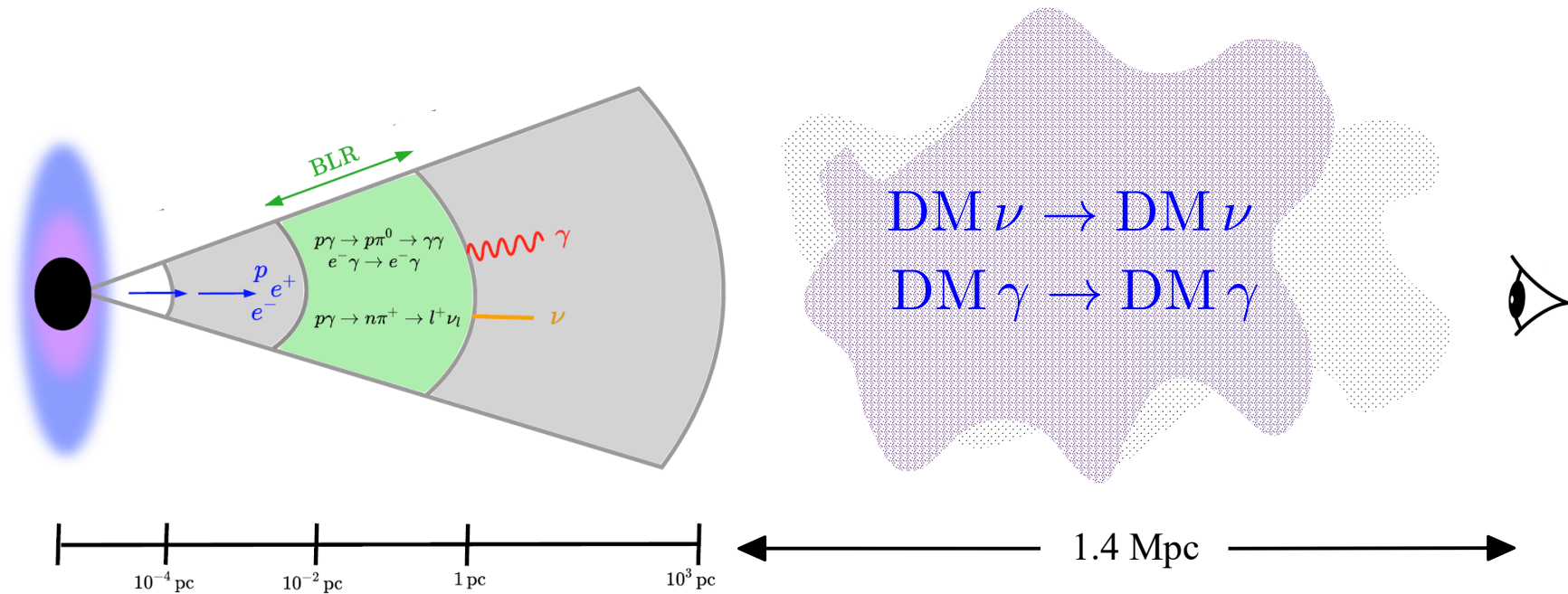
The neutrino and photon fluxes can be qualitatively well reproduced in leptohadronic models.



Gao et al'18

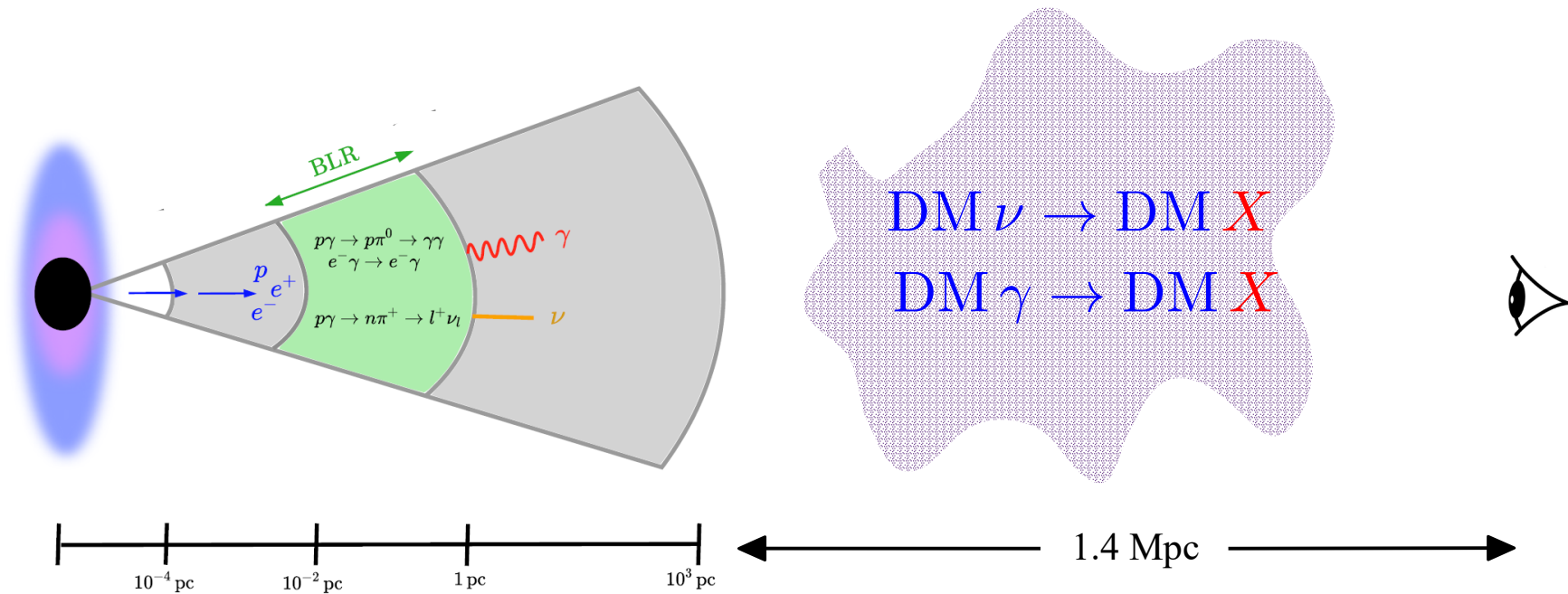


# Blazars as probes of the intergalactic medium



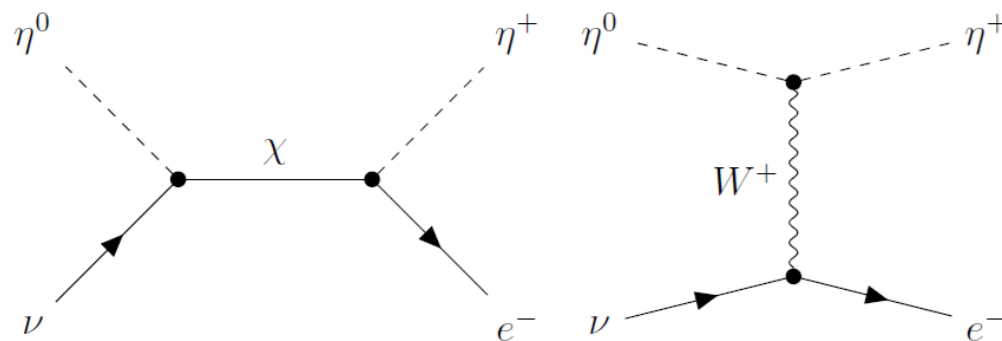
Neutrinos propagate through the intergalactic medium and through the Milky Way before reaching us. If the dark matter neutrino cross-section is large, the neutrino flux would be attenuated. Same with photons.

# Blazars as probes of the intergalactic medium



Neutrinos propagate through the intergalactic medium and through the Milky Way before reaching us. If the dark matter neutrino cross-section is large, the neutrino flux would be attenuated. Same with photons.

**Note: the absorption could also be due to inelastic scatterings**





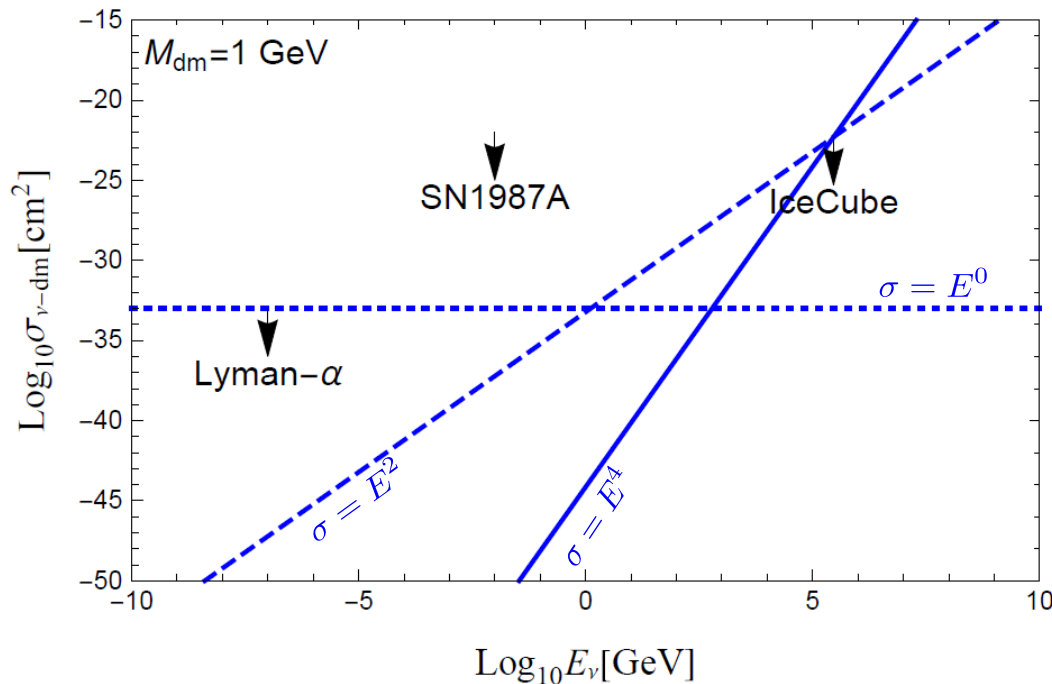
# Blazars as probes of the intergalactic medium

- Flux attenuation  $\frac{\Phi^{\text{obs}}}{\Phi^{\text{em}}} = \exp \left[ - \sigma_{\text{DM}-\nu} \Sigma_{\text{DM}} \right]$

$$\Sigma_{\text{DM}} = \frac{1}{m_{\text{DM}}} \int_{\text{path}} dr \rho(\vec{r}) \quad \text{is the number of DM particles along the path}$$

- Require that the attenuation of the neutrino flux due to DM-neutrino scatterings is less than 90%:

$$\frac{\sigma_{\text{DM}-\nu}}{m_{\text{DM}}} \lesssim \frac{2.3}{\Sigma_{\text{DM}}}$$



Choi, Kim, Rott'19  
Kelly, Machado '19

# The dark matter spike around TXS 0506+056

In the center of the blazar it is located a supermassive black hole, with mass  $\sim 3 \times 10^8 M_{\text{sun}}$ .

The growth of the black hole produces a “spike” in the dark matter distribution Gondolo, Silk'99, Peebles '72, Quinlan, Hernquist, Sigurdsson '95

$$\rho(r) = \rho_0 \left( \frac{r_0}{r} \right)^\gamma \longrightarrow \rho_{\text{sp}} \sim \rho_R \left( \frac{R_{\text{sp}}}{r} \right)^{\frac{9-2\gamma}{4-\gamma}}$$

For TXS 0506+056,  $R_{\text{sp}} \sim 1 \text{ pc}$

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$$\rho(r) = \rho_0 \left( \frac{r_0}{r} \right)^\gamma \longrightarrow \rho_{\text{sp}} \sim \rho_R \left( \frac{R_{\text{sp}}}{r} \right)^{\frac{9-2\gamma}{4-\gamma}}$$

For TXS 0506+056,  $R_{\text{sp}} \sim 1 \text{ pc}$

Dark matter annihilations soften the spike to a saturation density, proportional to  $\langle \sigma v \rangle$

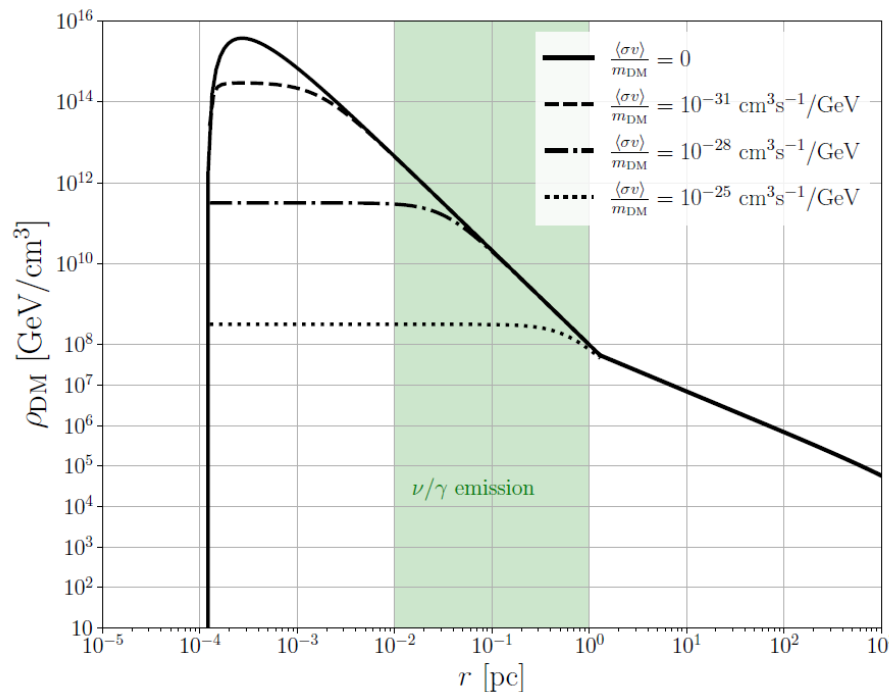
$$\rho(r) = \begin{cases} 0 & r \leq 4R_S \\ \frac{\rho_{\text{sp}}(r)\rho_{\text{sat}}}{\rho_{\text{sp}}(r) + \rho_{\text{sat}}} & 4R_S \leq r \leq R_{\text{sp}} \\ \rho_0 \left( \frac{r}{r_0} \right)^{-\gamma} \left( 1 + \frac{r}{r_0} \right)^{-3+\gamma} & r \geq R_{\text{sp}} \end{cases}$$

# The dark matter spike around TXS 0506+056

In the center of the blazar it is located a supermassive black hole, with mass  $\sim 3 \times 10^8 M_{\text{sun}}$ .

The growth of the black hole produces a “spike” in the dark matter distribution Gondolo, Silk'99, Peebles '72, Quinlan, Hernquist, Sigurdsson '95

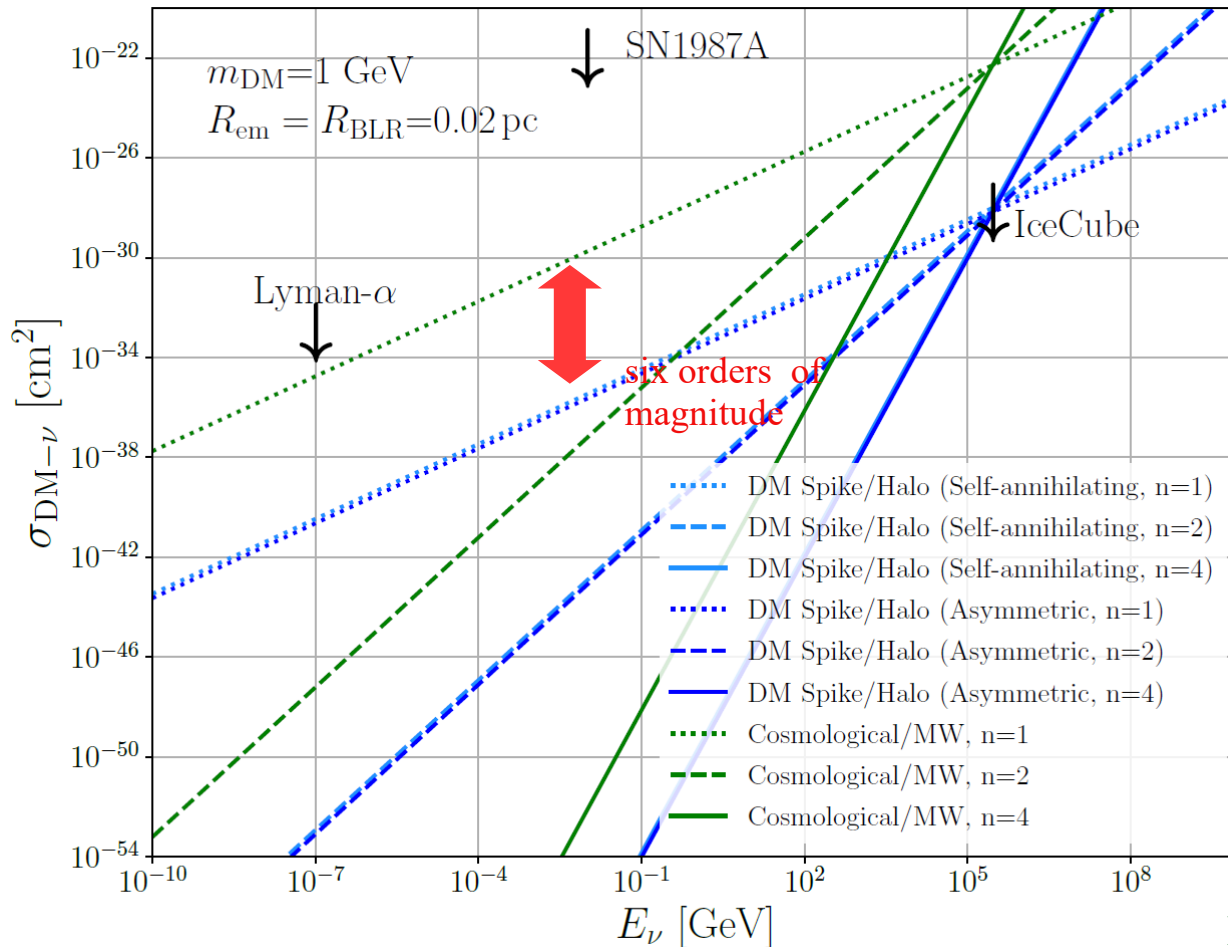
$$\rho(r) = \rho_0 \left( \frac{r_0}{r} \right)^\gamma \longrightarrow \rho_{\text{sp}} \sim \rho_R \left( \frac{R_{\text{sp}}}{r} \right)^{\frac{9-2\gamma}{4-\gamma}}$$



Ferrer, Herrera, AI'22

The region of emission of neutrinos and photons lies inside the dark matter spike  $\Rightarrow$  Increase in the column density.

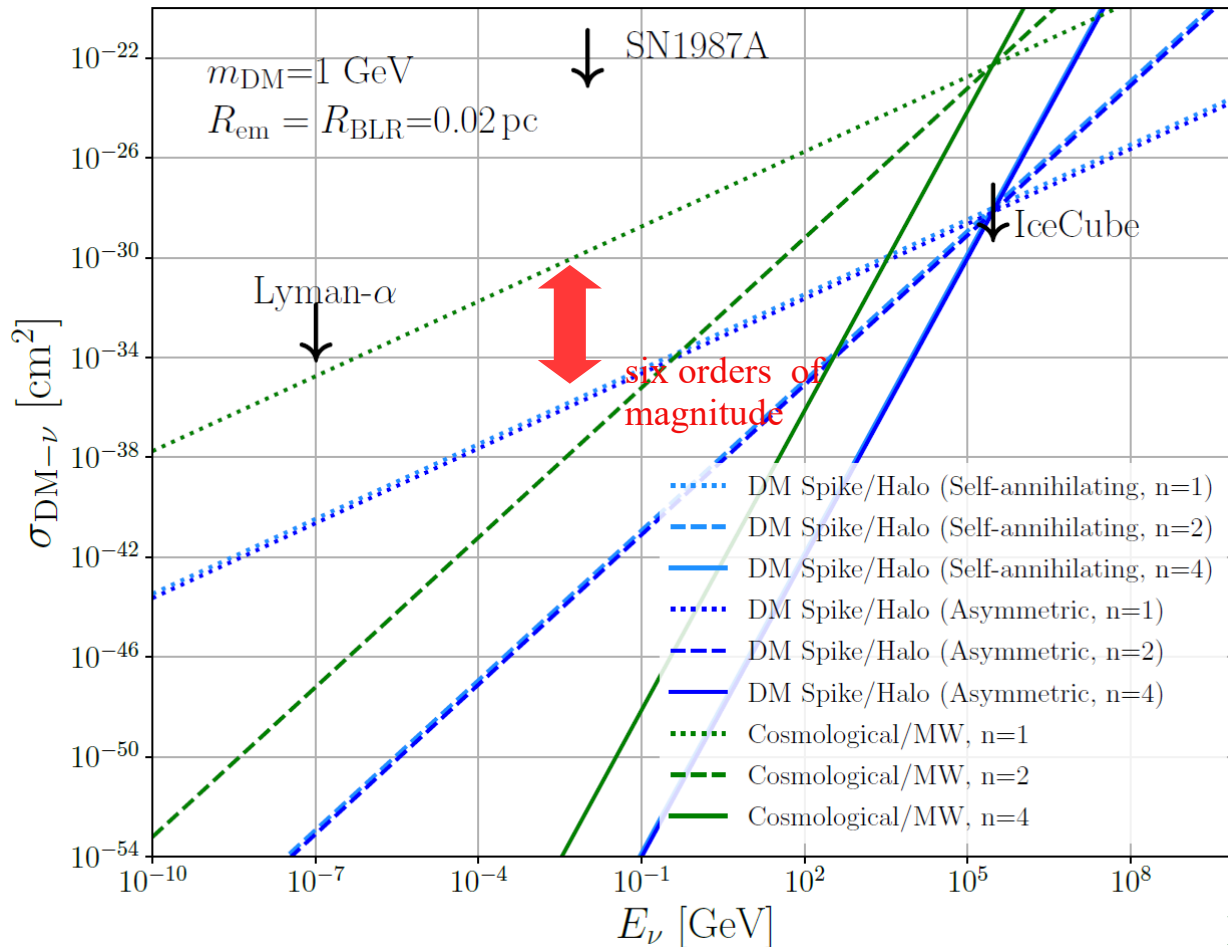
# Probing the scattering cross-section with neutrinos



Ferrer, Herrera, AI'22  
See also Cline et al'22

Complementarity with other searches?

# Probing the scattering cross-section with neutrinos

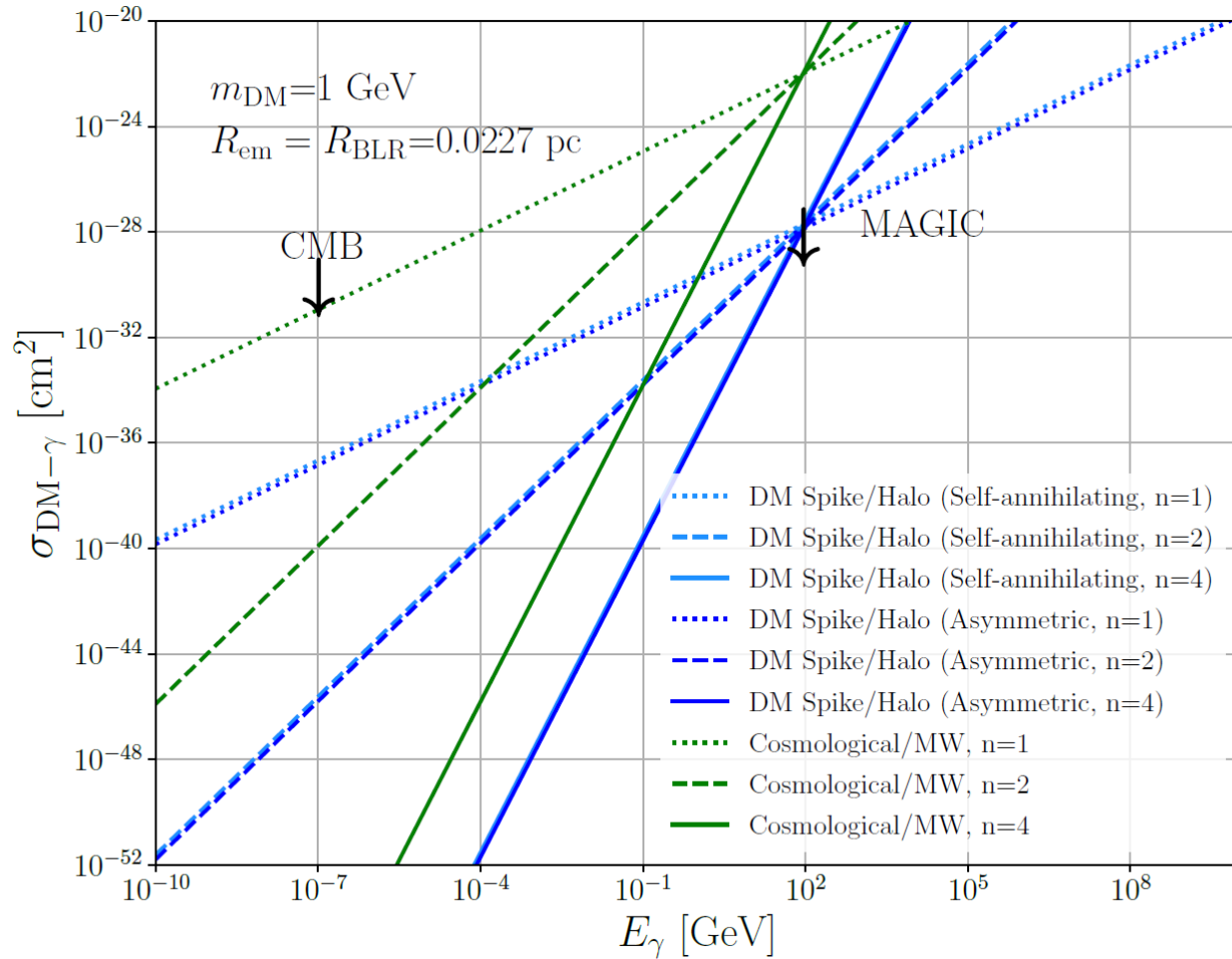


Ferrer, Herrera, AI'22  
See also Cline et al'22

Complementarity with other searches?

Better and more robust limits could be obtained as more and more high energy  $\nu$  sources are discovered.

# Probing the scattering cross-section with photons



Ferrer, Herrera, AI'22

# Conclusions

- After 40+ years of search, there is still no concluding evidence that dark matter is made of elementary particles.
- Neutrino telescopes could provide hints for dark matter, especially if dark matter annihilates mostly into neutrinos, and when dark matter interacts with nucleons via the spin-dependent interaction.
- The recent discovery of high energy astrophysical neutrino sources by IceCube opens new opportunities to probe dark matter properties, and new opportunities to probe neutrino mass models.