

CNB essentials 3: Bounds on neutrino properties from cosmology



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Bounds on N_{eff}

Relativistic particles in the universe

At $T < m_e$, the radiation content of the Universe is

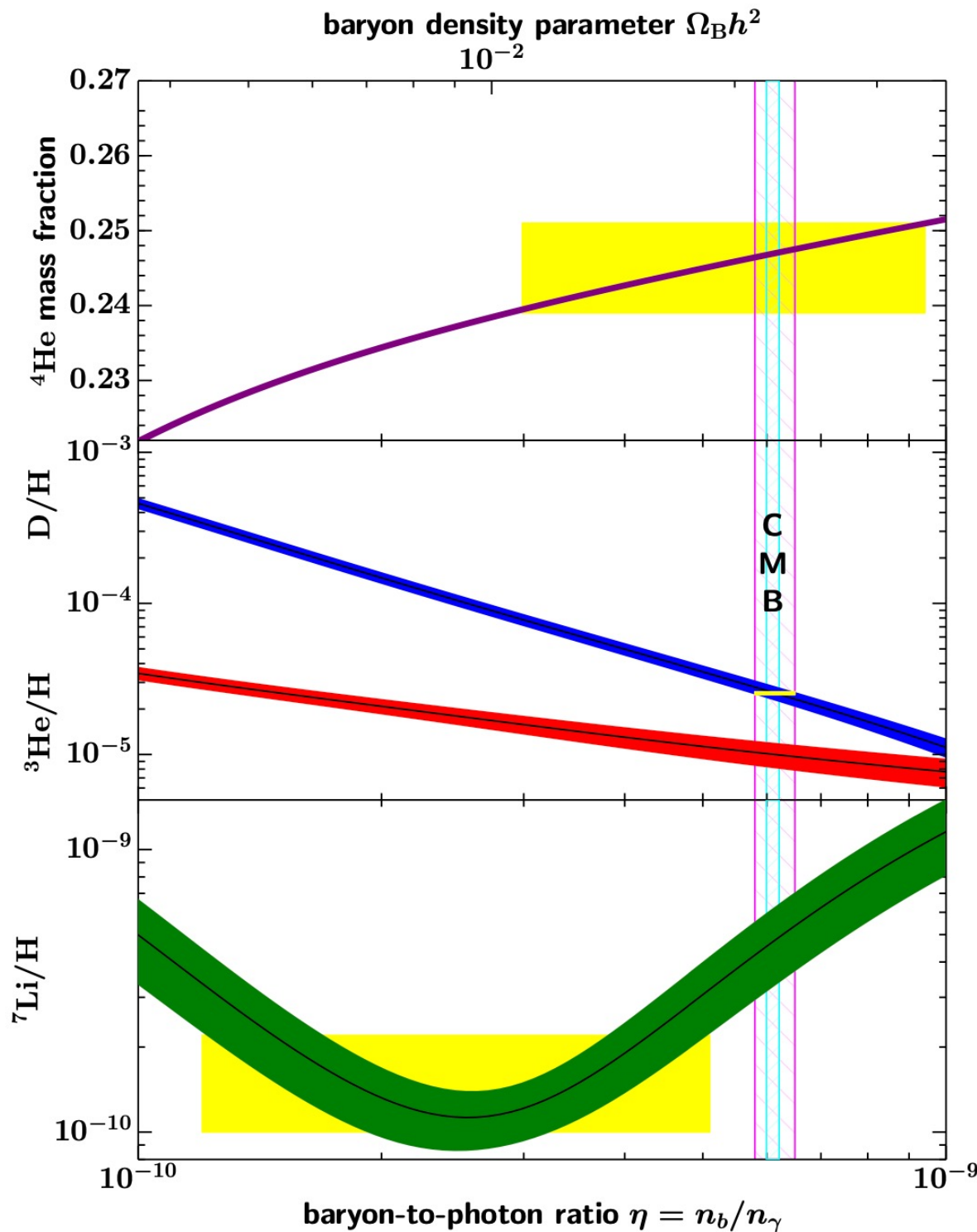
$$\rho_{\text{rad}} = \rho_{\gamma} + \rho_{\nu} + \rho_x = \rho_{\gamma} \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right]$$

effective number of relativistic neutrino species
(effective number of neutrinos)

N_{eff} is a way to measure the ratio $\frac{\rho_{\nu} + \rho_x}{\rho_{\gamma}}$

Number of light neutrino types (LEP data) $N_{\nu} = 2.984 \pm 0.008$

BBN: Predictions vs Observations



$$\eta_{10} = \frac{n_B/n_\gamma}{10^{-10}} \simeq 274 \Omega_B h^2$$

$$5.8 \leq \eta_{10} \leq 6.5$$

(95% CL)

Fields, Molaro & Sarkar,
PDG 2020

Effect of neutrinos on Primordial Nucleosynthesis

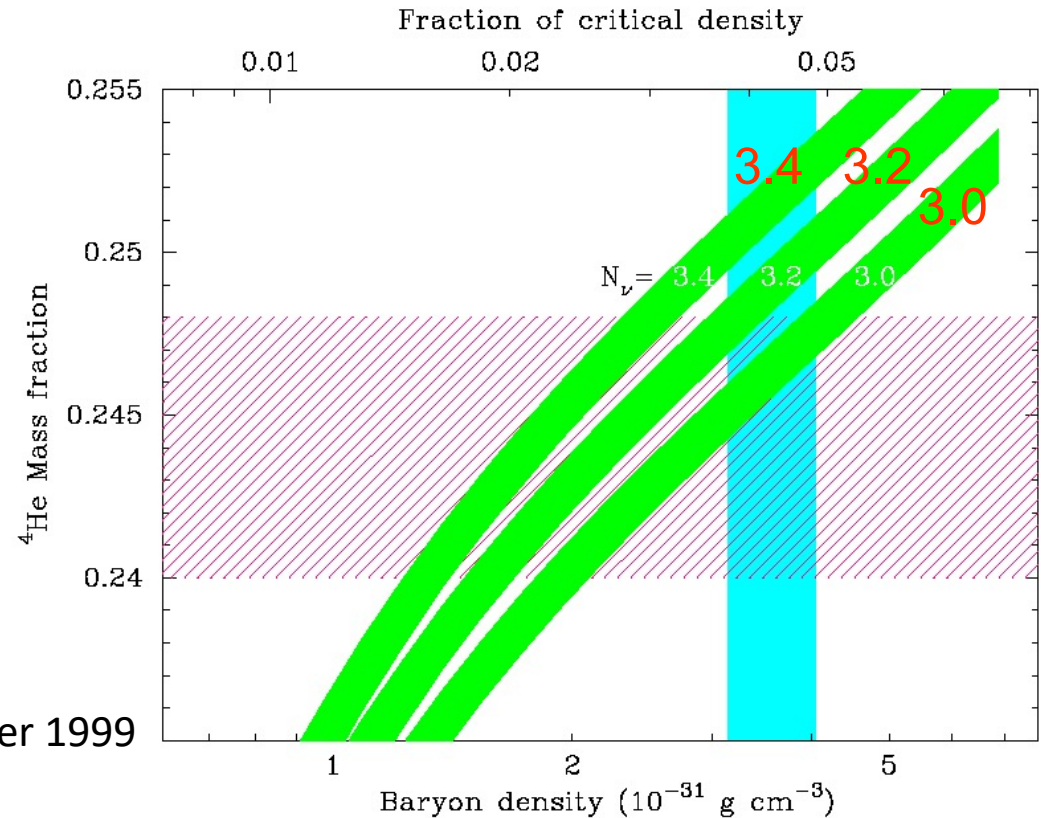
1. N_{eff} fixes the **expansion rate** during BBN

$$H = \sqrt{\frac{8\pi\rho}{3M_p^2}}$$

$$\rho(N_{\text{eff}}) > \rho_0 \rightarrow \uparrow {}^4\text{He}$$

$$\Delta Y_p \simeq 0.013 \Delta N_{\text{eff}}$$

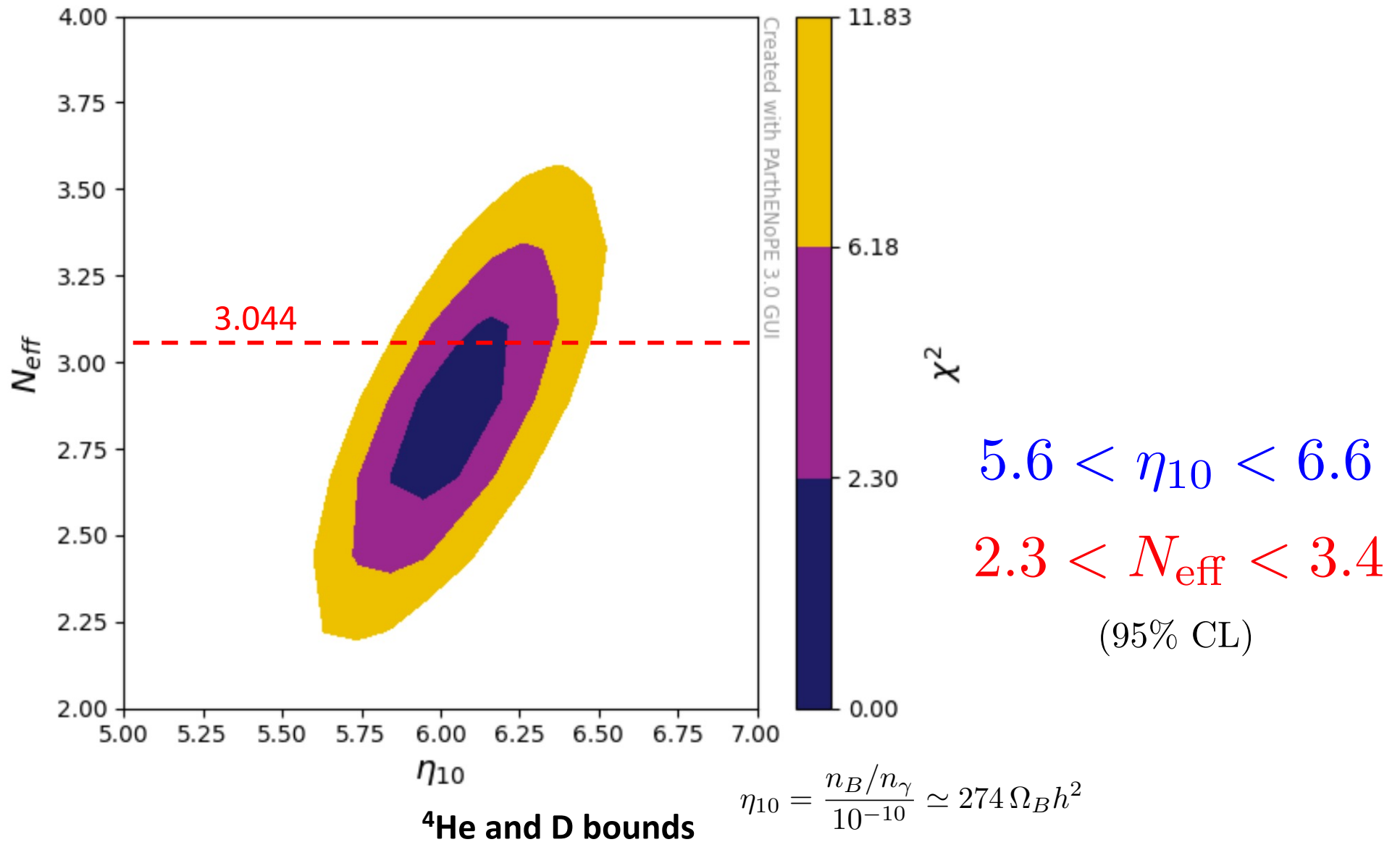
Burles, Nollett & Turner 1999



2. Direct effect of **electron** neutrinos and antineutrinos on the **n-p reactions**



BBN: allowed ranges for N_{eff}



ParthENoPE BBN code, S Gariazzo et al, arXiv:2103.05027

The minimal Λ CDM model fits very well Planck data

Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits
$\Omega_b h^2$	0.02212 ± 0.00022	0.02249 ± 0.00025	0.0240 ± 0.0012
$\Omega_c h^2$	0.1206 ± 0.0021	0.1177 ± 0.0020	0.1158 ± 0.0046
$100\theta_{MC}$	1.04077 ± 0.00047	1.04139 ± 0.00049	1.03999 ± 0.00089
τ	0.0522 ± 0.0080	0.0496 ± 0.0085	0.0527 ± 0.0090
$\ln(10^{10} A_s)$	3.040 ± 0.016	$3.018^{+0.020}_{-0.018}$	3.052 ± 0.022
n_s	0.9626 ± 0.0057	0.967 ± 0.011	0.980 ± 0.015
H_0 [km s ⁻¹ Mpc ⁻¹] . .	66.88 ± 0.92	68.44 ± 0.91	69.9 ± 2.7
Ω_Λ	0.679 ± 0.013	0.699 ± 0.012	$0.711^{+0.033}_{-0.026}$
Ω_m	0.321 ± 0.013	0.301 ± 0.012	$0.289^{+0.026}_{-0.033}$

Parameter	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits	TT,TE,EE+lowE+lensing+BAO 68% limits
$\Omega_b h^2$	0.02236 ± 0.00015	0.02237 ± 0.00015	0.02242 ± 0.00014
$\Omega_c h^2$	0.1202 ± 0.0014	0.1200 ± 0.0012	0.11933 ± 0.00091
$100\theta_{MC}$	1.04090 ± 0.00031	1.04092 ± 0.00031	1.04101 ± 0.00029
τ	$0.0544^{+0.0070}_{-0.0081}$	0.0544 ± 0.0073	0.0561 ± 0.0071
$\ln(10^{10} A_s)$	3.045 ± 0.016	3.044 ± 0.014	3.047 ± 0.014
n_s	0.9649 ± 0.0044	0.9649 ± 0.0042	0.9665 ± 0.0038
H_0 [km s ⁻¹ Mpc ⁻¹]	67.27 ± 0.60	67.36 ± 0.54	67.66 ± 0.42
Ω_Λ	0.6834 ± 0.0084	0.6847 ± 0.0073	0.6889 ± 0.0056
Ω_m	0.3166 ± 0.0084	0.3153 ± 0.0073	0.3111 ± 0.0056

CMB anisotropies + other data

$$N_{\text{eff}} \lesssim 17$$

(2001) early CMB data

$$N_{\text{eff}} = 4.2^{+1.2}_{-1.7}$$

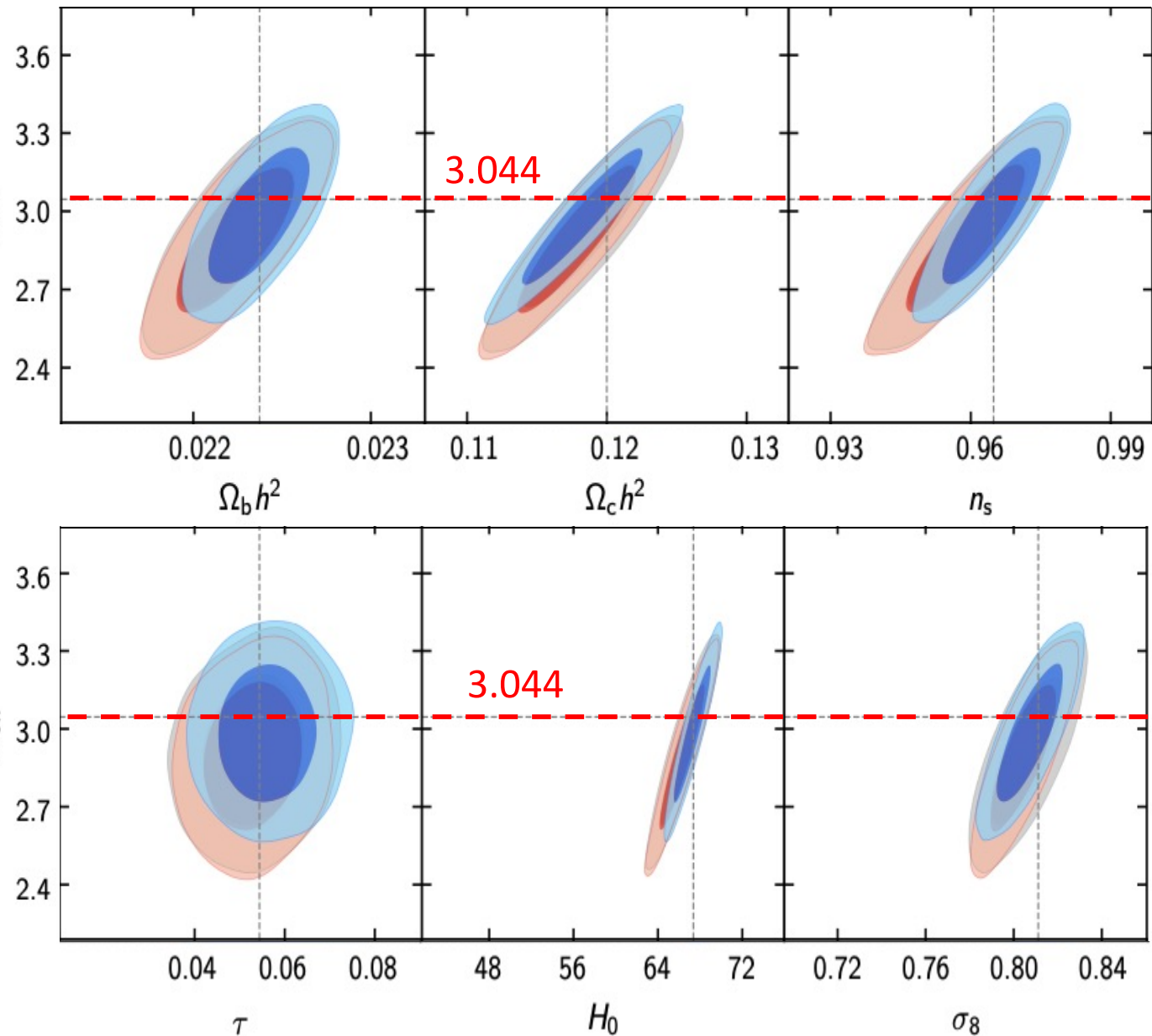
(2005) WMAP+...

Planck: 1-parameter extensions of the Λ CDM model

68+95%
Conf regions

Planck TT, TE, EE + lowE + lensing + lensing + BAO

N_{eff}



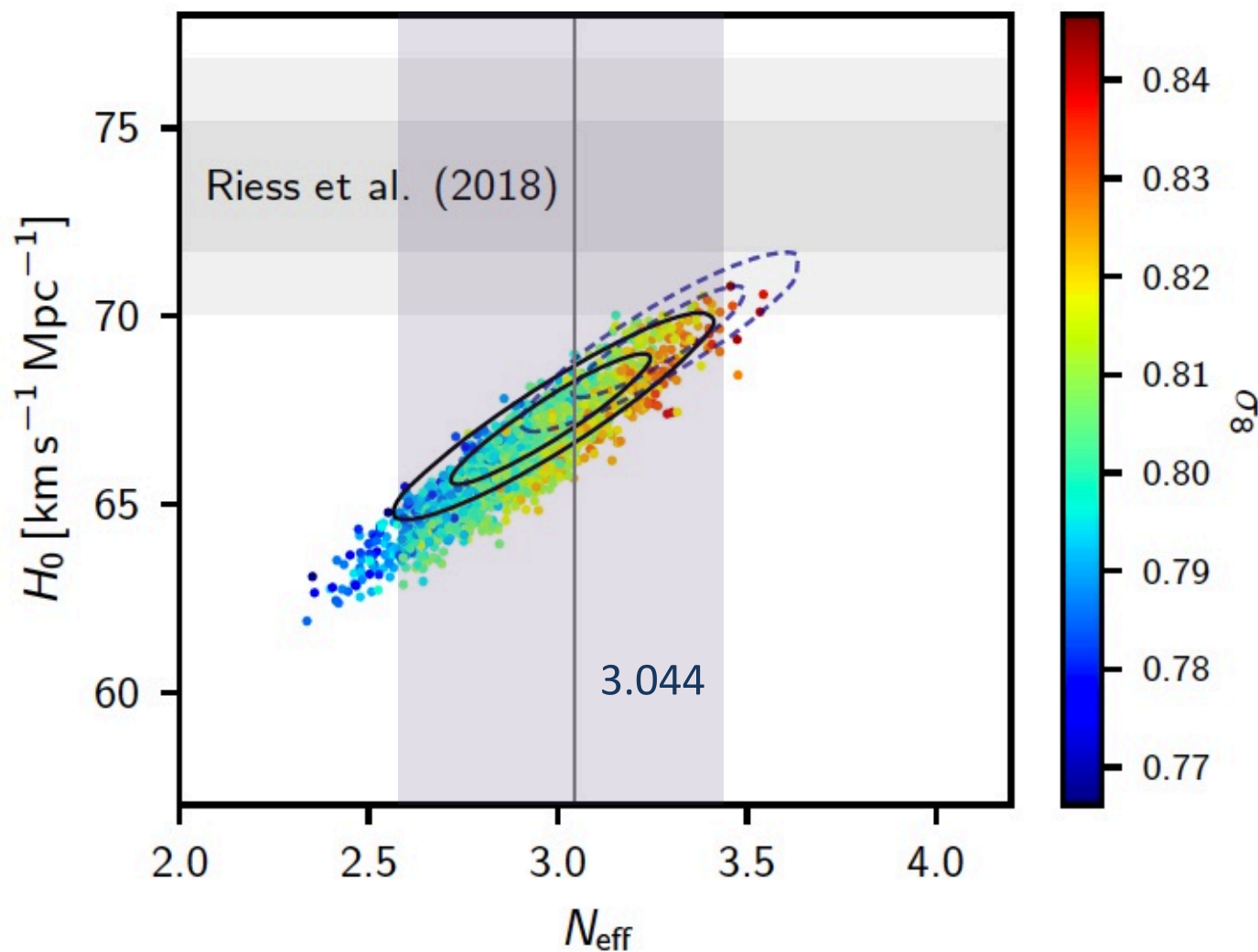
CMB anisotropies + other data

$$N_{\text{eff}} \lesssim 17 \quad \text{(2001) early CMB data}$$

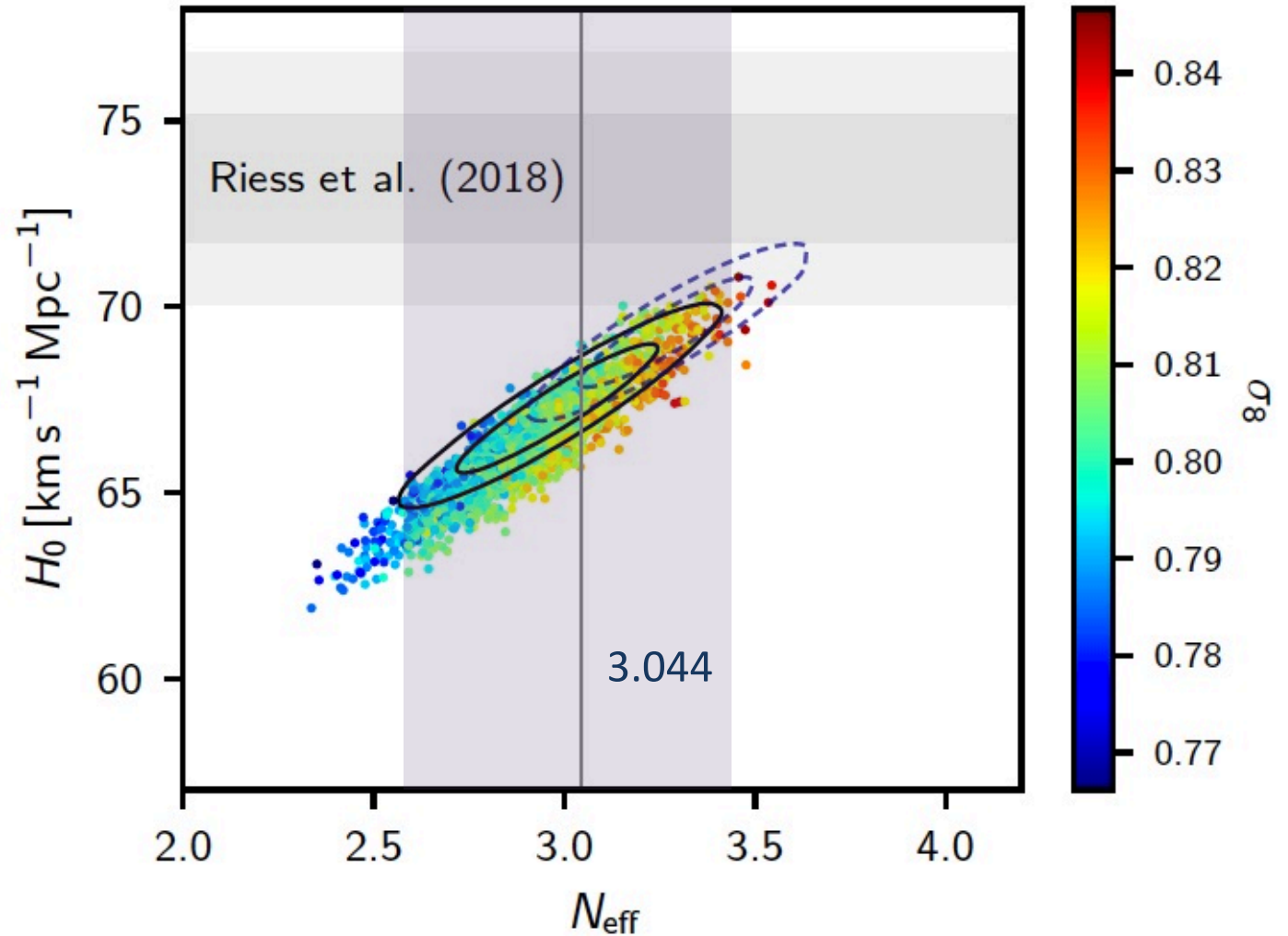
$$N_{\text{eff}} = 4.2^{+1.2}_{-1.7} \quad \text{(2005) WMAP+...}$$

$$N_{\text{eff}} = 2.99^{+0.34}_{-0.33} \quad \text{(2018) Planck}$$

(95%, TT,TE,EE+lowE+lensing+BAO)



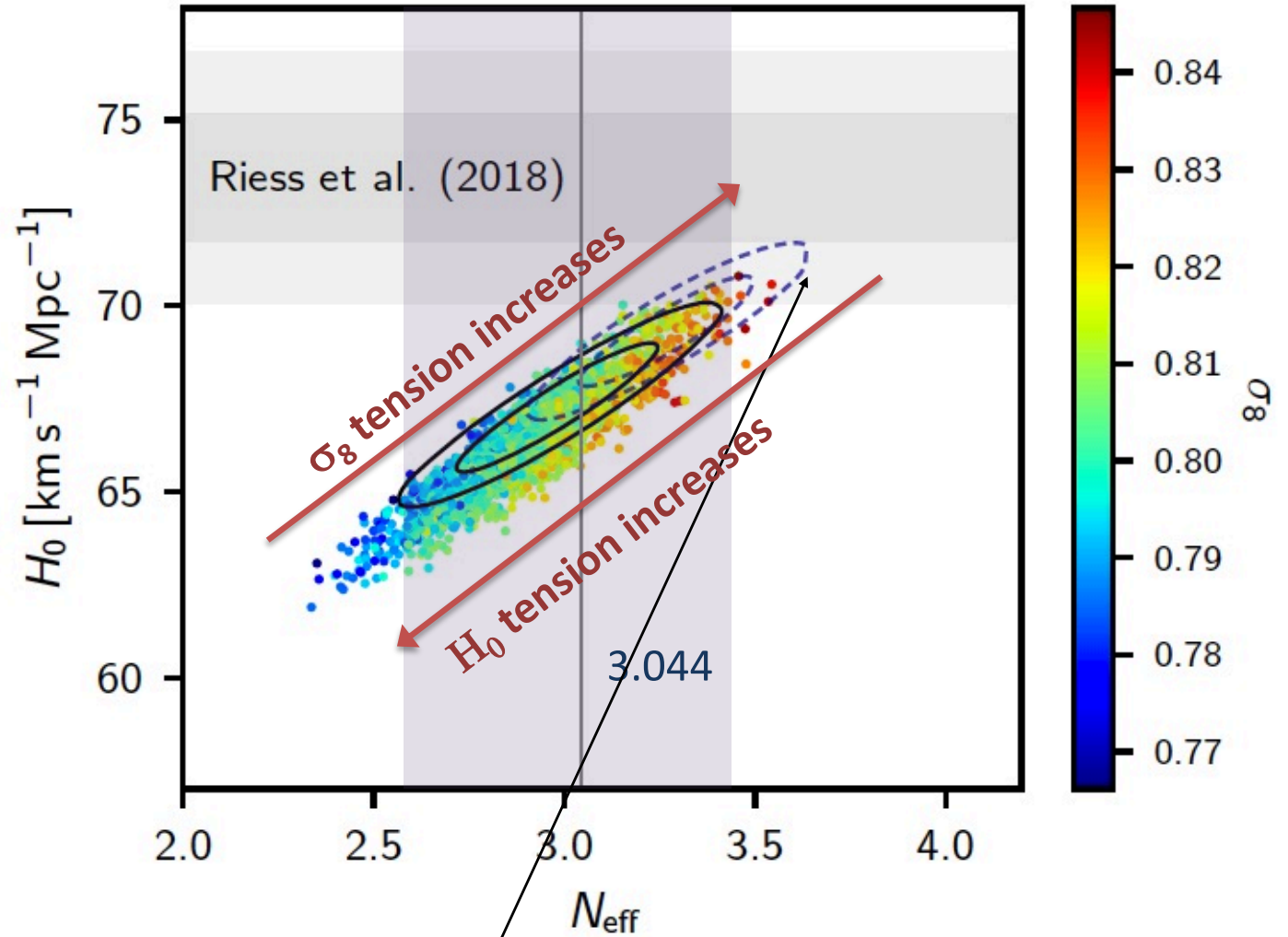
CMB anisotropies + other data



$$N_{\text{eff}} = 2.92^{+0.36}_{-0.37} \quad (95\%, \text{Planck TT,TE,EE+lowE})$$

$$N_{\text{eff}} = 2.99^{+0.34}_{-0.33} \quad (95\%, \text{TT,TE,EE+lowE+lensing+BAO})$$

CMB anisotropies + other data

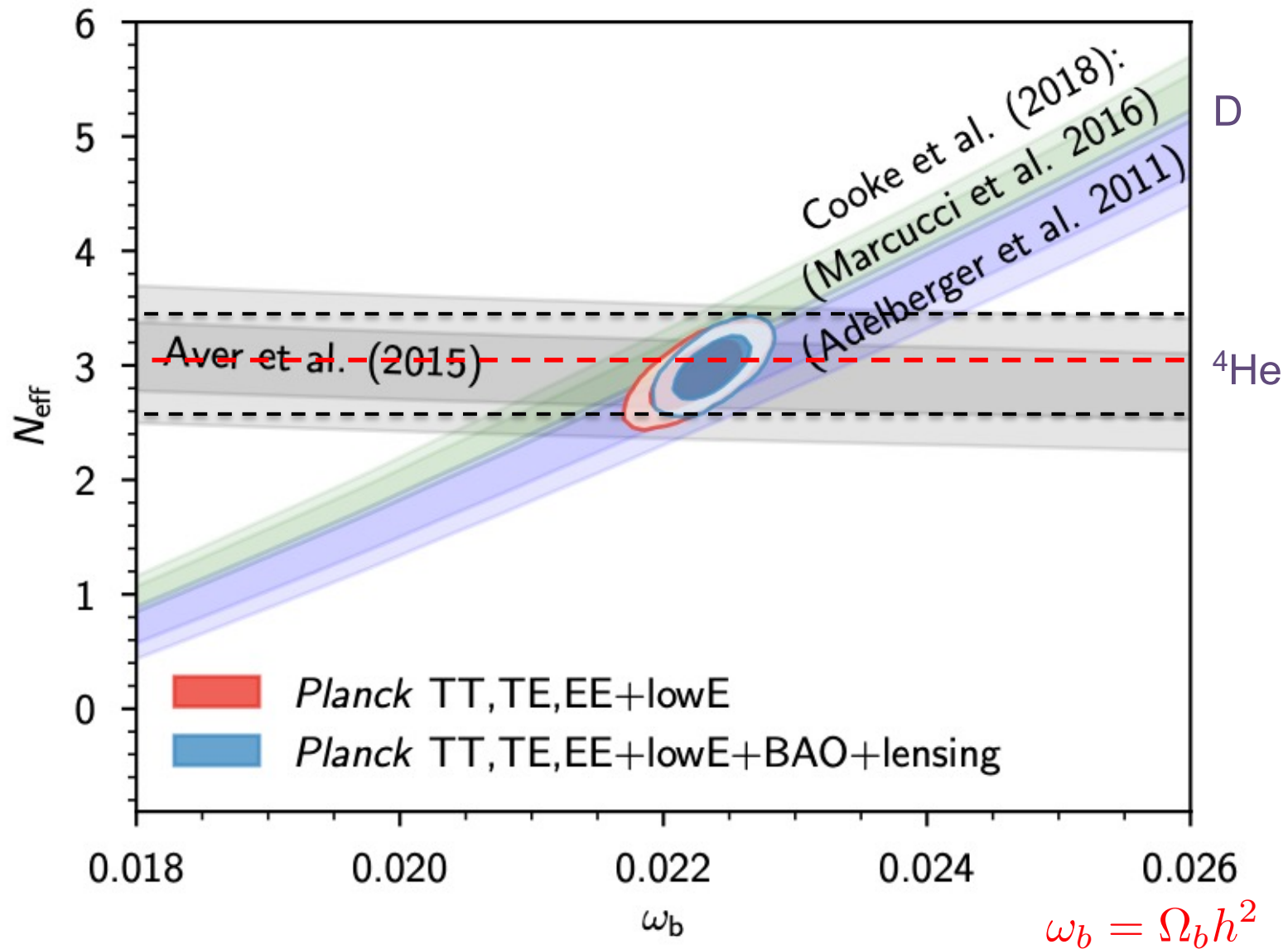


$$N_{\text{eff}} = 3.27 \pm 0.15$$

$$H_0 = (69.32 \pm 0.97) \text{ km s}^{-1} \text{ Mpc}^{-1}$$

} 68%, TT,TE,EE+lowE+lensing+BAO+R18.

Comparison: allowed ranges for N_{eff} and BBN



**N_{eff} with non-standard
neutrino-electron interactions**

Non-standard neutrino-electron interactions

Non-standard interactions (NSI) between **neutrinos** and **electrons** can be parametrised as follows:

$$\mathcal{L}_{\text{NSI}e} = -2\sqrt{2}G_F \sum_{\alpha,\beta} \varepsilon_{\alpha\beta}^X (\bar{\nu}_\alpha \gamma^\mu P_L \nu_\beta) (\bar{e} \gamma_\mu P_X e)$$

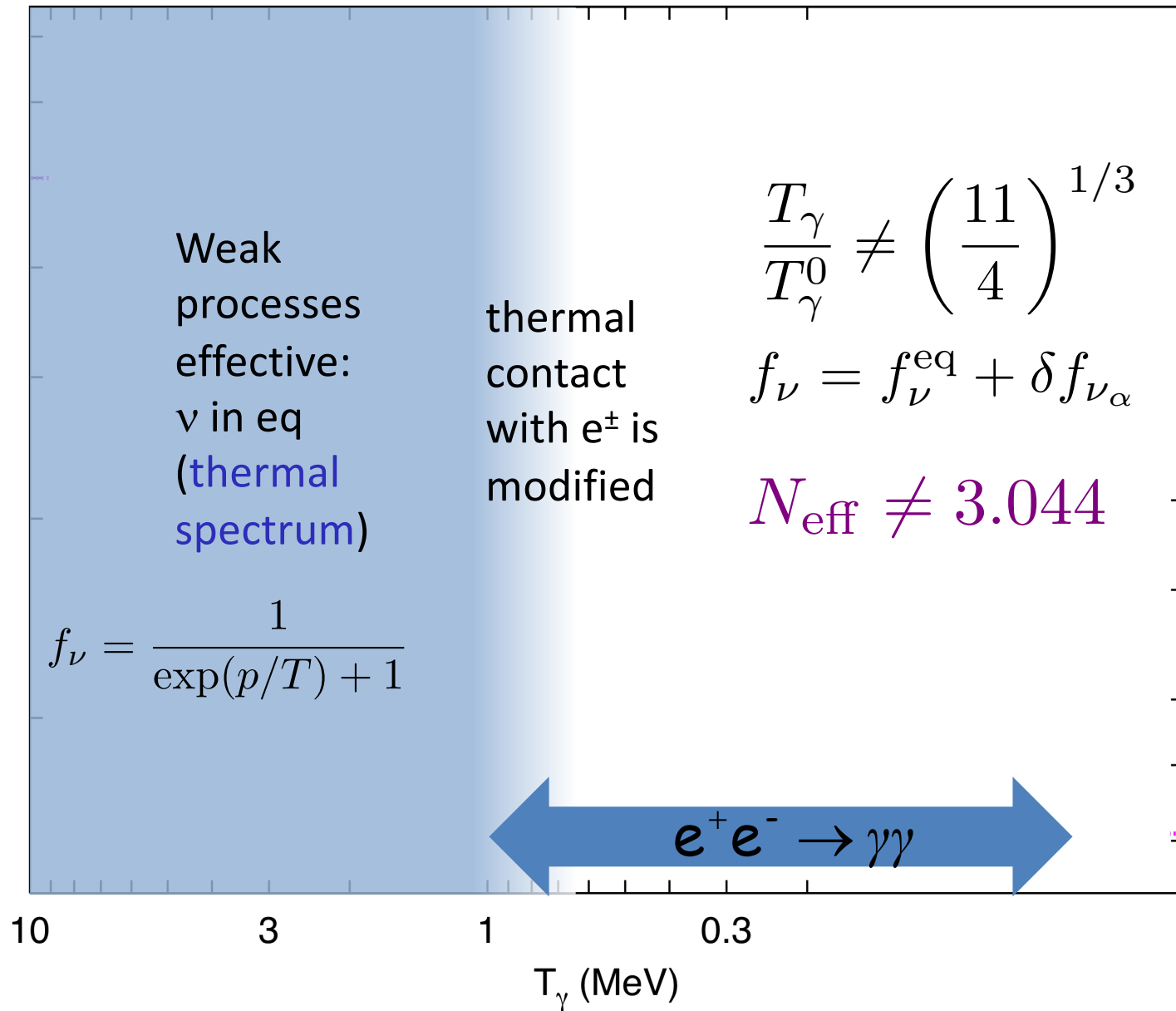
$$\text{with } X \in \{L, R\} \\ \alpha, \beta \in \{e, \mu, \tau\}$$

Dimensionless coefficients $\varepsilon_{\alpha\beta}^X$ quantify the strength of the interactions with respect to the SM

$\varepsilon_{\alpha\alpha}^X$ **Non-universal NSI**

$\varepsilon_{\alpha\beta}^X$ (with $\alpha \neq \beta$) **Flavour-changing ($\alpha \neq \beta$) NSI**

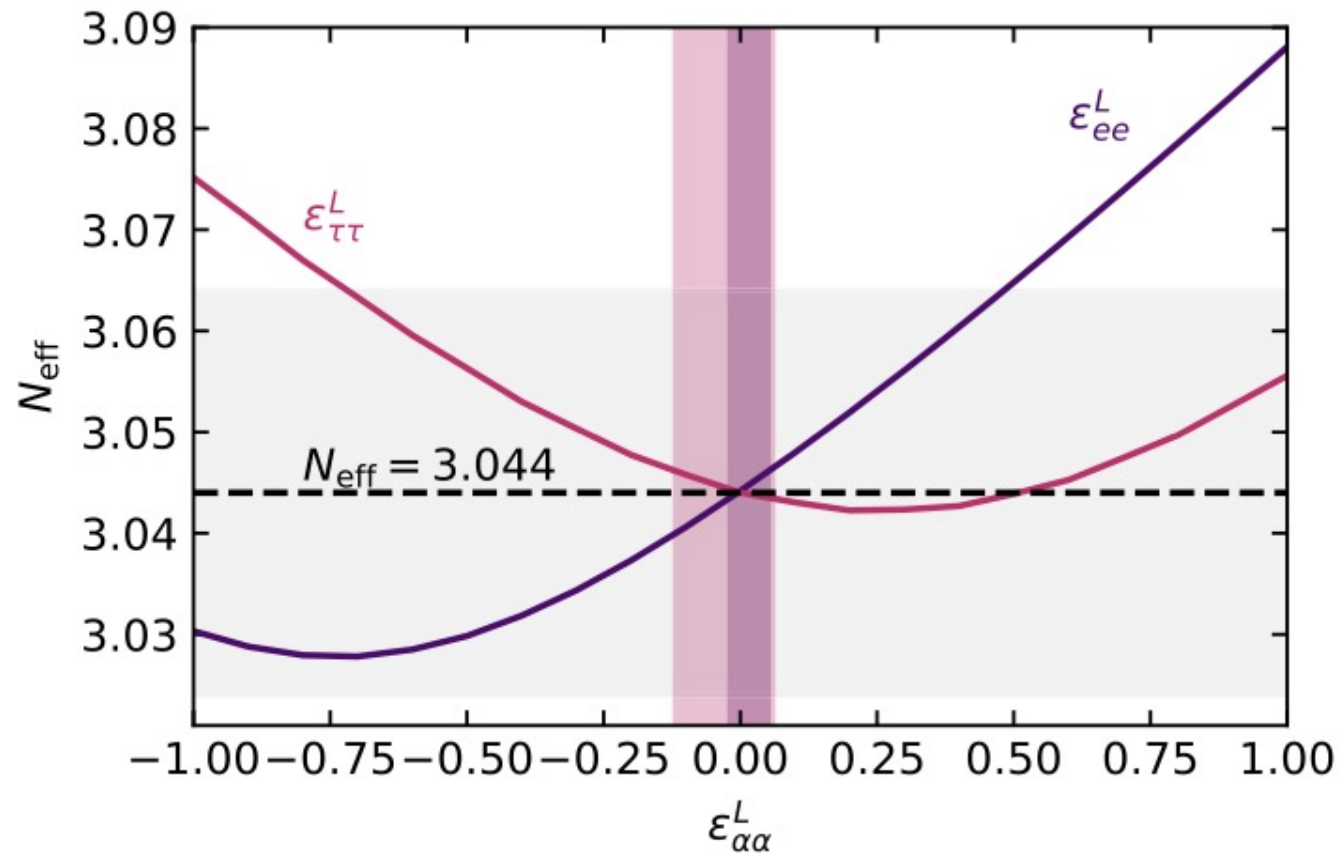
N_{eff} with non-standard neutrino-electron interactions



Expansion of the universe

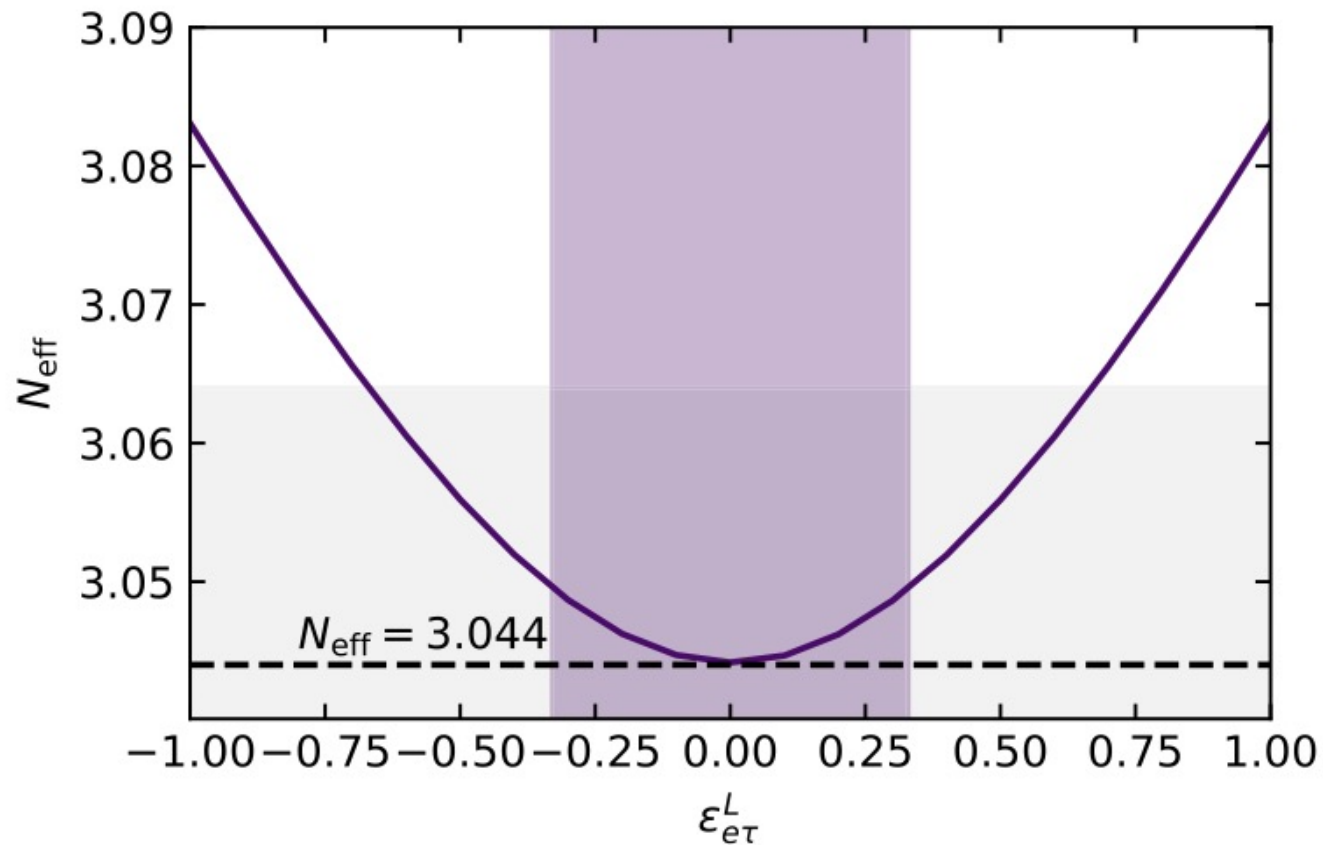
N_{eff} with only one NSI parameter

N_{eff} for non-universal NSI

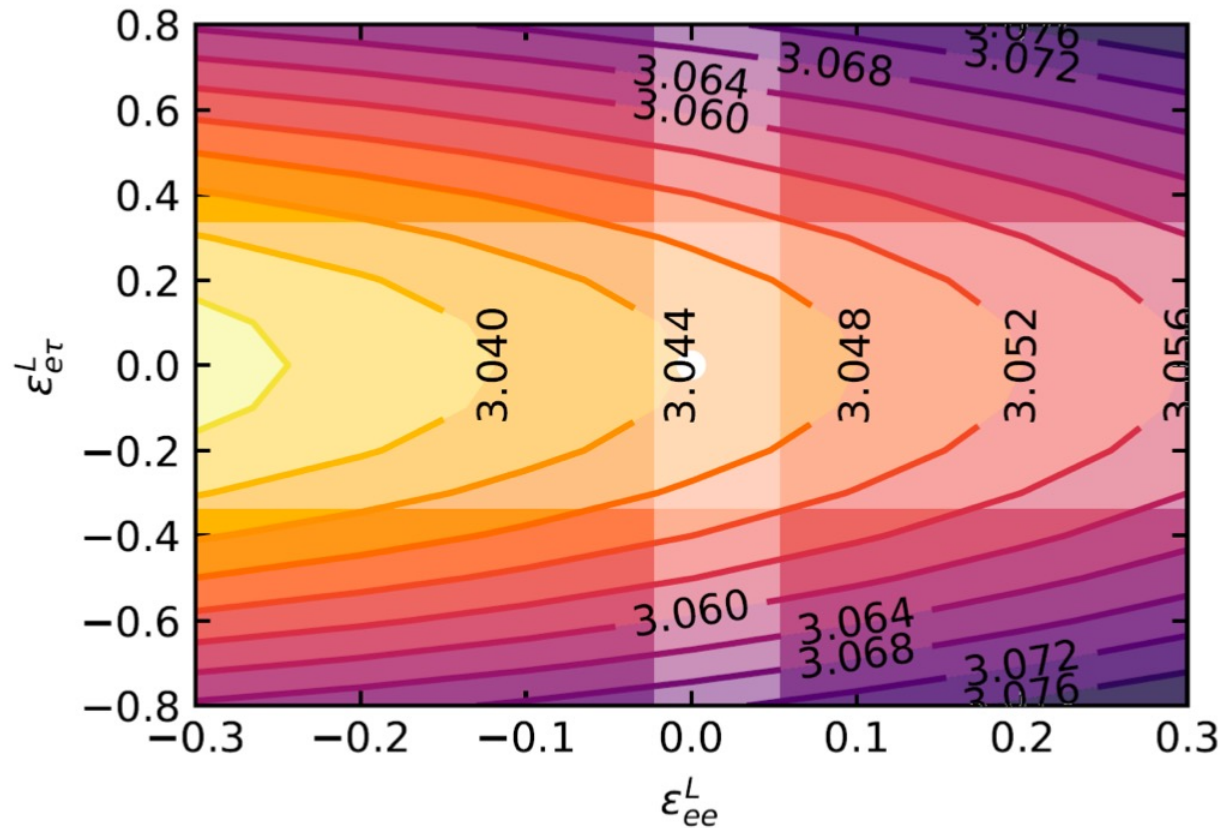


N_{eff} with only one NSI parameter

N_{eff} for flavour-changing NSI



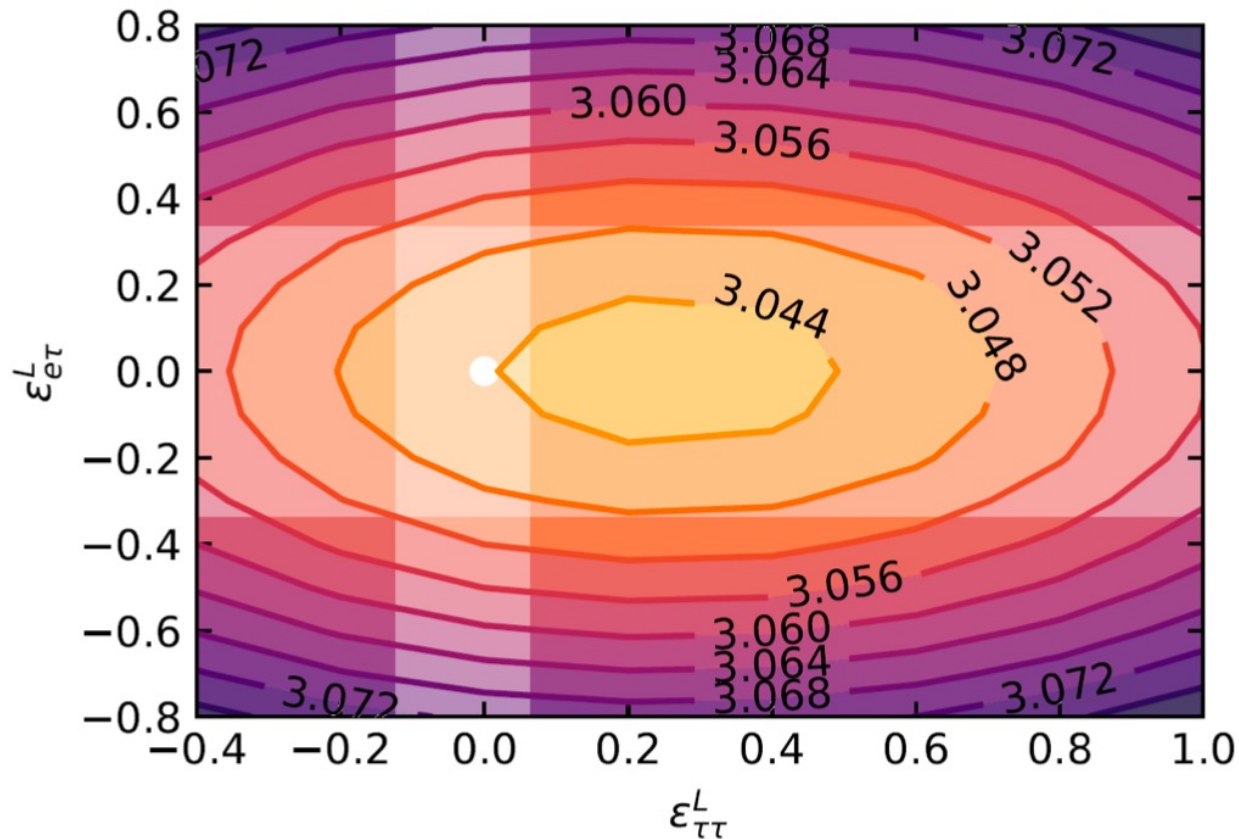
N_{eff} varying 2 NSI parameters



Future determinations of N_{eff} are expected to have an error of 0.02-0.03

White shaded bands correspond to terrestrial bounds on NSI.

N_{eff} varying 2 NSI parameters



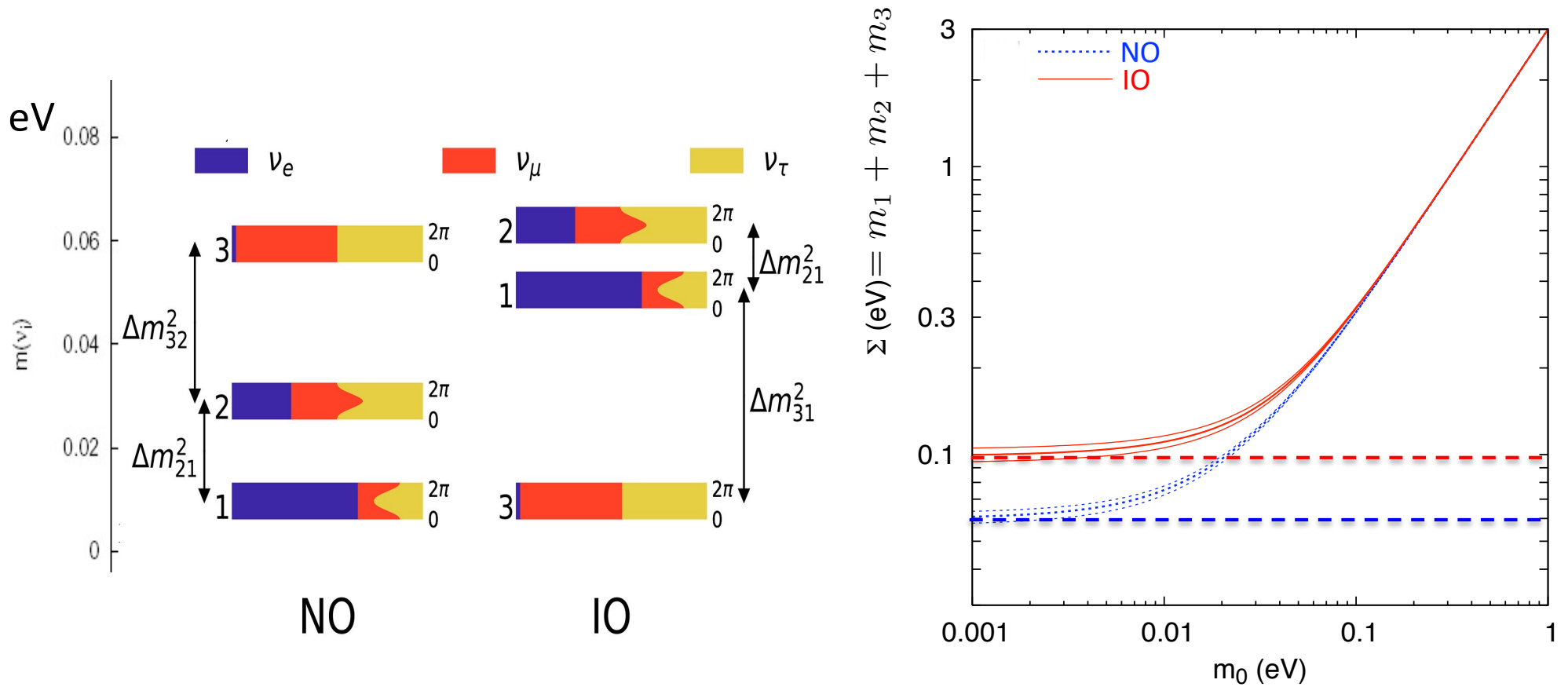
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Bounds on neutrino masses

Neutrino masses

Data on flavour oscillations do not fix the absolute scale of neutrino masses





$$0.06 \text{ (0.1) eV} \lesssim \sum_i m_{\nu_i} \lesssim 2.4 \text{ eV}$$

Neutrinos as Dark Matter

- Neutrinos are natural **DM candidates**

$$\Omega_\nu h^2 = \frac{\sum_i m_i}{93.2 \text{ eV}} \quad \Omega_\nu < 1 \rightarrow \sum_i m_i \lesssim 46 \text{ eV}$$

$$\Omega_\nu < \Omega_m \simeq 0.3 \rightarrow \sum_i m_i \lesssim 15 \text{ eV}$$

- They stream freely until non-relativistic (collisionless phase mixing) 
Neutrinos are HOT Dark Matter (large thermal motion)
- First structures to be formed when Universe became matter –dominated are very large
- Ruled out by structure formation  CDM

Massive Neutrinos can still be subdominant DM: **limits on m_ν from Structure Formation (combined with other cosmological data)**

Cosmological bounds on neutrino mass(es)

A unique cosmological bound on m_ν DOES NOT exist !

Different analyses have found upper bounds on neutrino masses, since they depend on

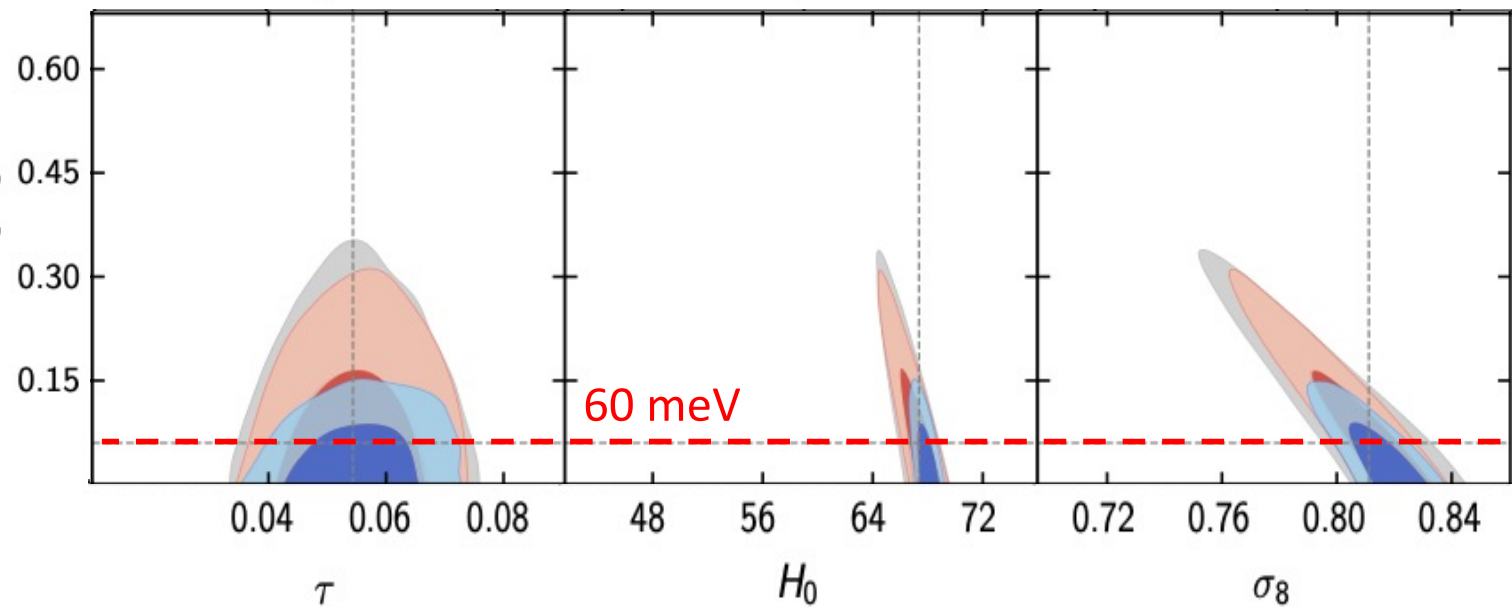
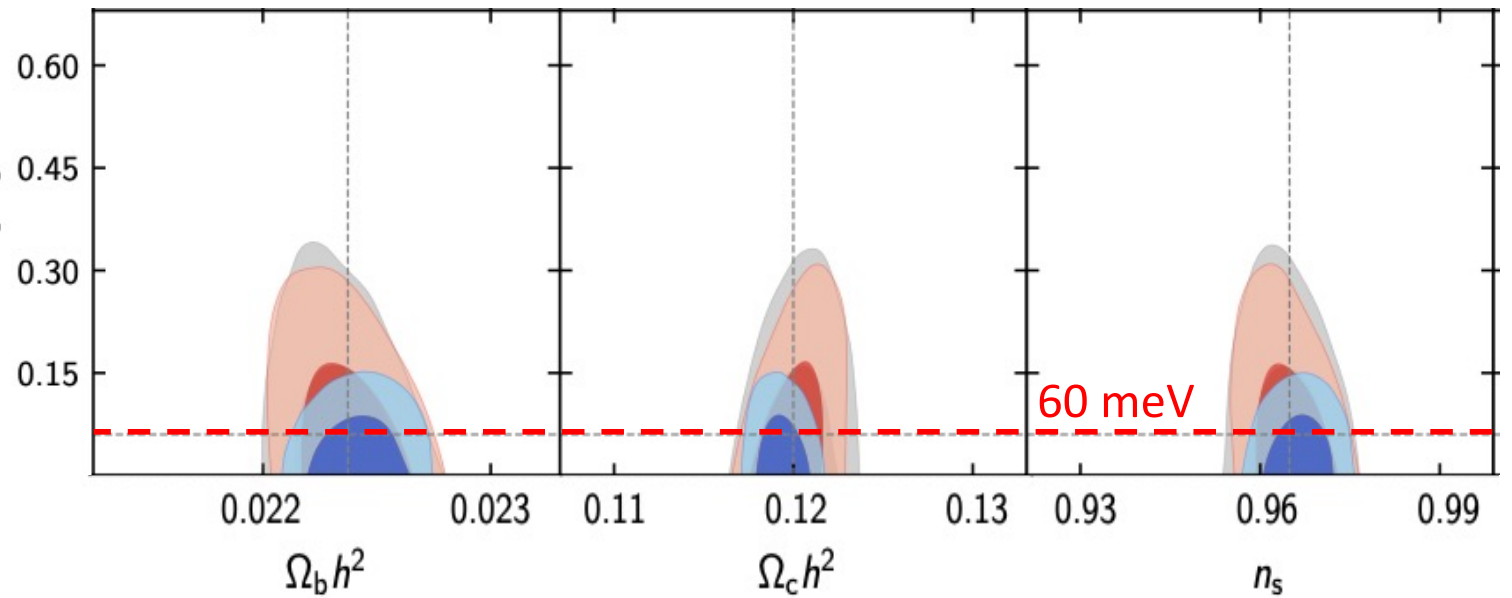
- The combination of **cosmological data** used
- The assumed **cosmological model**: number of parameters (problem of parameter degeneracies)
- The **properties of relic neutrinos**

Planck: 1-parameter extensions of the Λ CDM model

68+95%
Conf regions

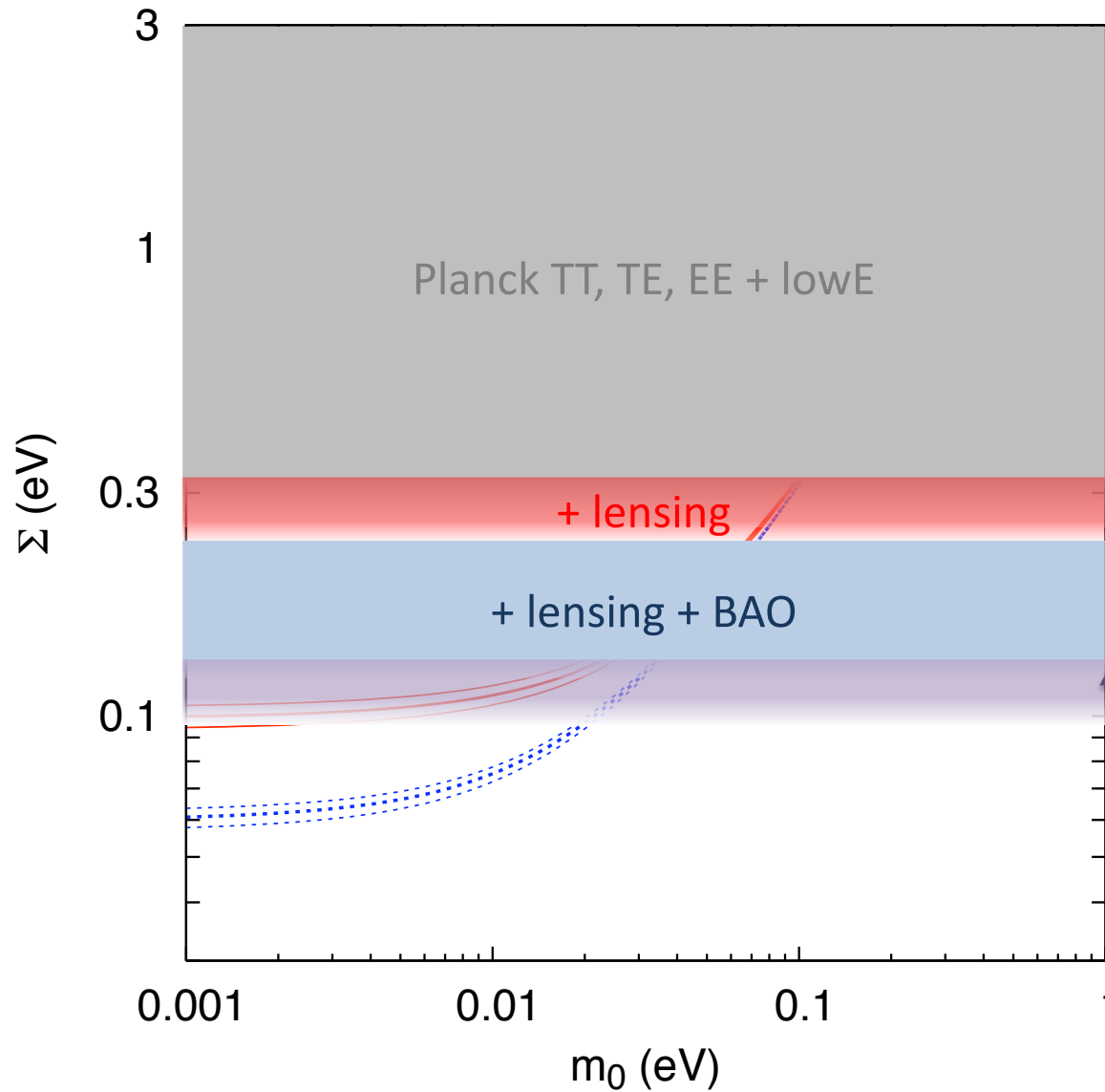
Planck TT, TE, EE + lowE + lensing + lensing + BAO

Σm_ν
(eV)



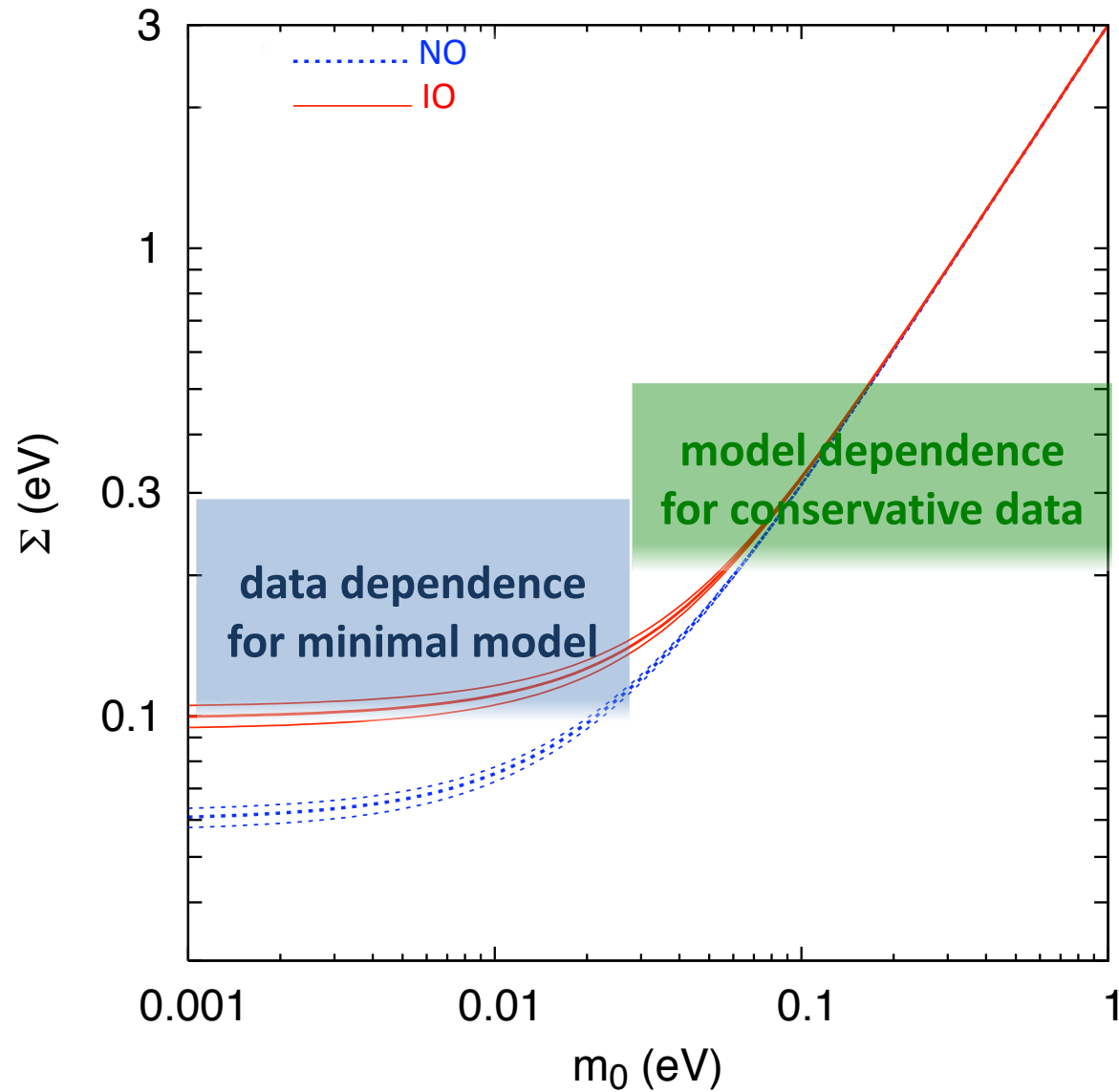
Bounds on Σm_ν from Planck (+other cosmo data)

Cosmological upper limits on the sum of neutrino masses



Bounds on Σm_ν from Planck (+other cosmo data)

Cosmological upper limits on the sum of neutrino masses



Probing the absolute neutrino mass scale

Searching for non-zero neutrino mass in laboratory experiments

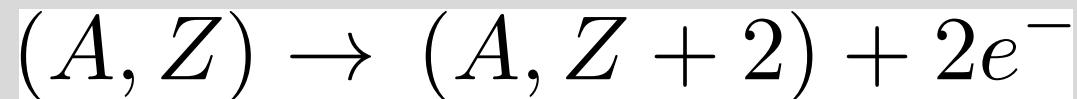
- **Tritium beta decay**: measurements of endpoint energy



$m_{\beta} < 2.2 \text{ eV}$ (95% CL) Mainz

Current experiment (KATRIN) $m(\nu_e) < 0.8 \text{ eV}$ (90% CL)

- **Neutrinoless double beta decay**: if Majorana neutrinos

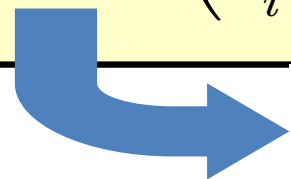


experiments with ${}^{76}\text{Ge}$, ${}^{130}\text{Te}$, ${}^{136}\text{Xe}$ and other isotopes:

$m_{\beta\beta} < 60\text{-}600 \text{ meV}$, depending on NME

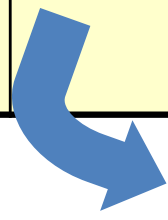
Probing the absolute neutrino mass scale

Tritium β decay	$m_\beta = \left(\sum_i U_{ei} ^2 m_i^2 \right)^{1/2}$	80 meV
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$$[c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

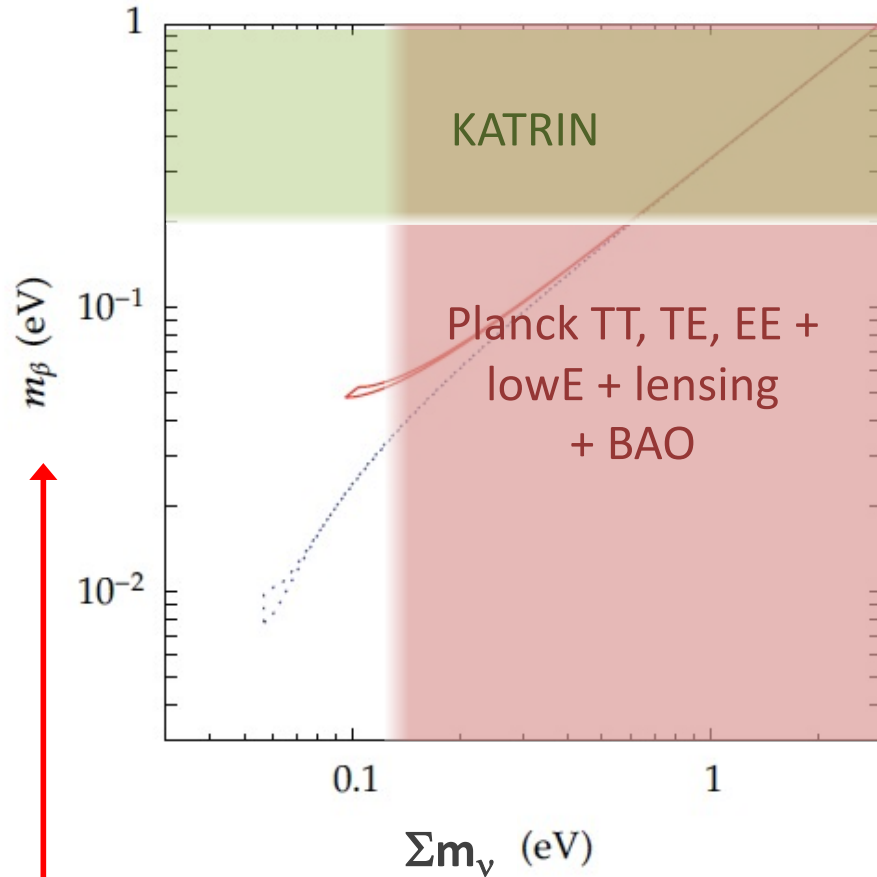
Neutrinoless double beta decay	$m_{\beta\beta} = \left \sum_i U_{ei}^2 m_i \right $	< 60-600 meV
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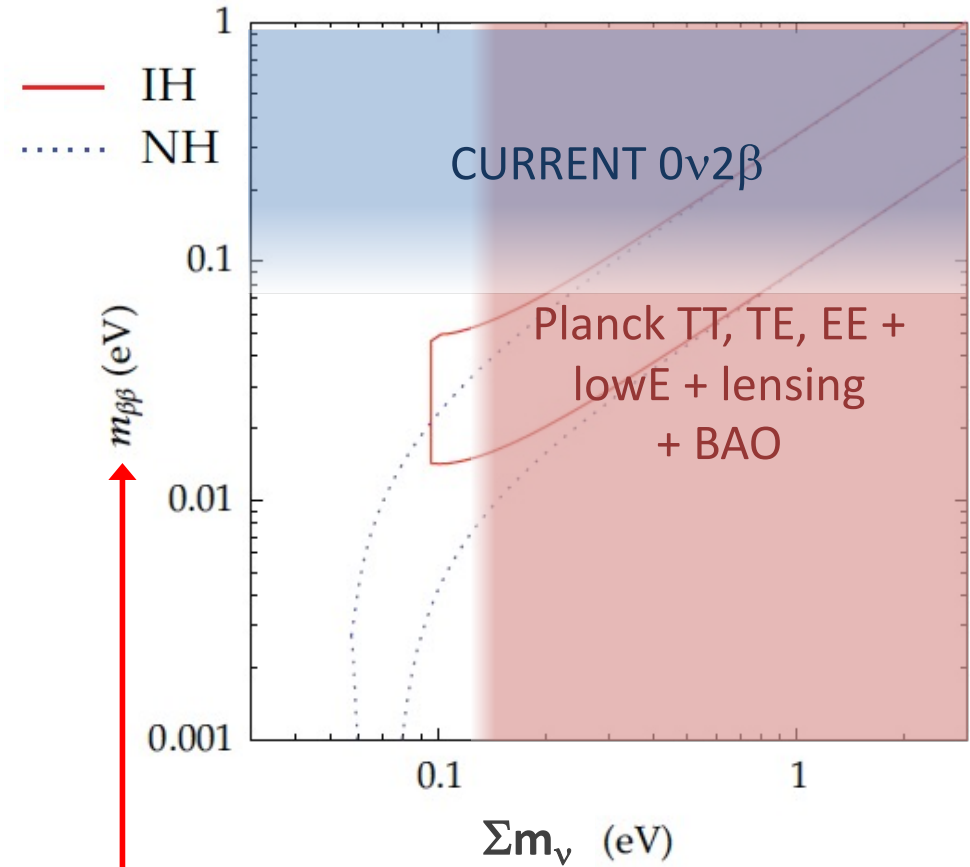
$$|c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

Cosmology	$\sim \sum_i m_i$	< 90-500 meV
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Tritium β decay, $0\nu 2\beta$ and Cosmology



$$[c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$



$$|c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

Bounds on active-sterile oscillations (3+1 case)

Mixing of four neutrino states?

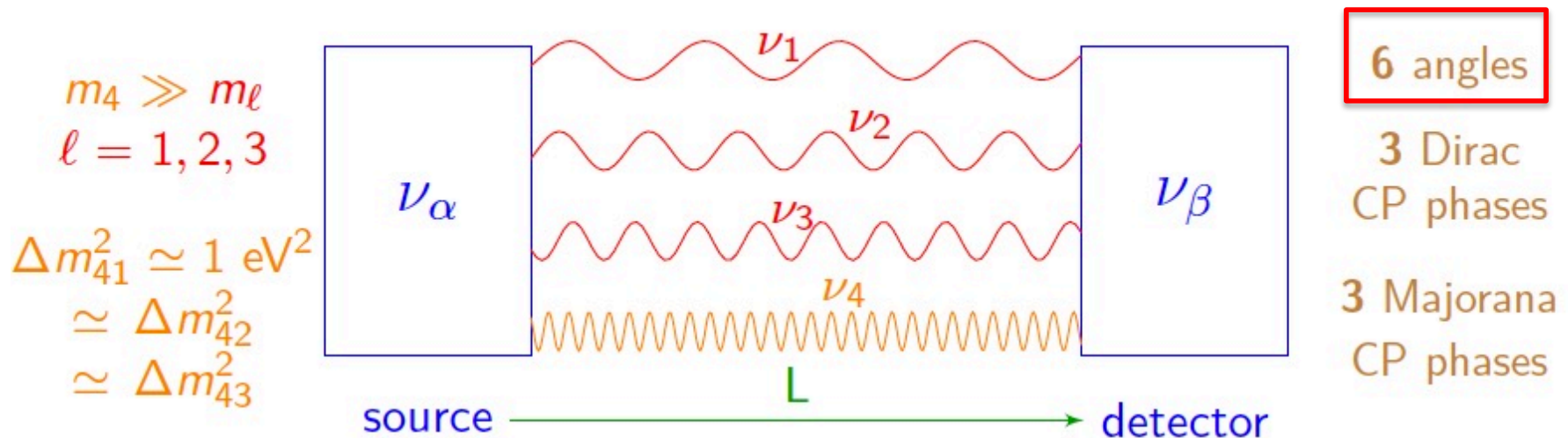
Additional neutrino (**sterile**) states introduced in order to explain some anomalies in experimental data

4 flavour neutrinos, 4 massive neutrinos

4x4 mixing matrix

$$\begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} & U_{\mu 4} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} & U_{\tau 4} \\ U_{s11} & U_{s12} & U_{s13} & U_{s14} \end{pmatrix}$$

We consider **3 (active) + 1 (sterile)**, a perturbation of the 3-neutrino case



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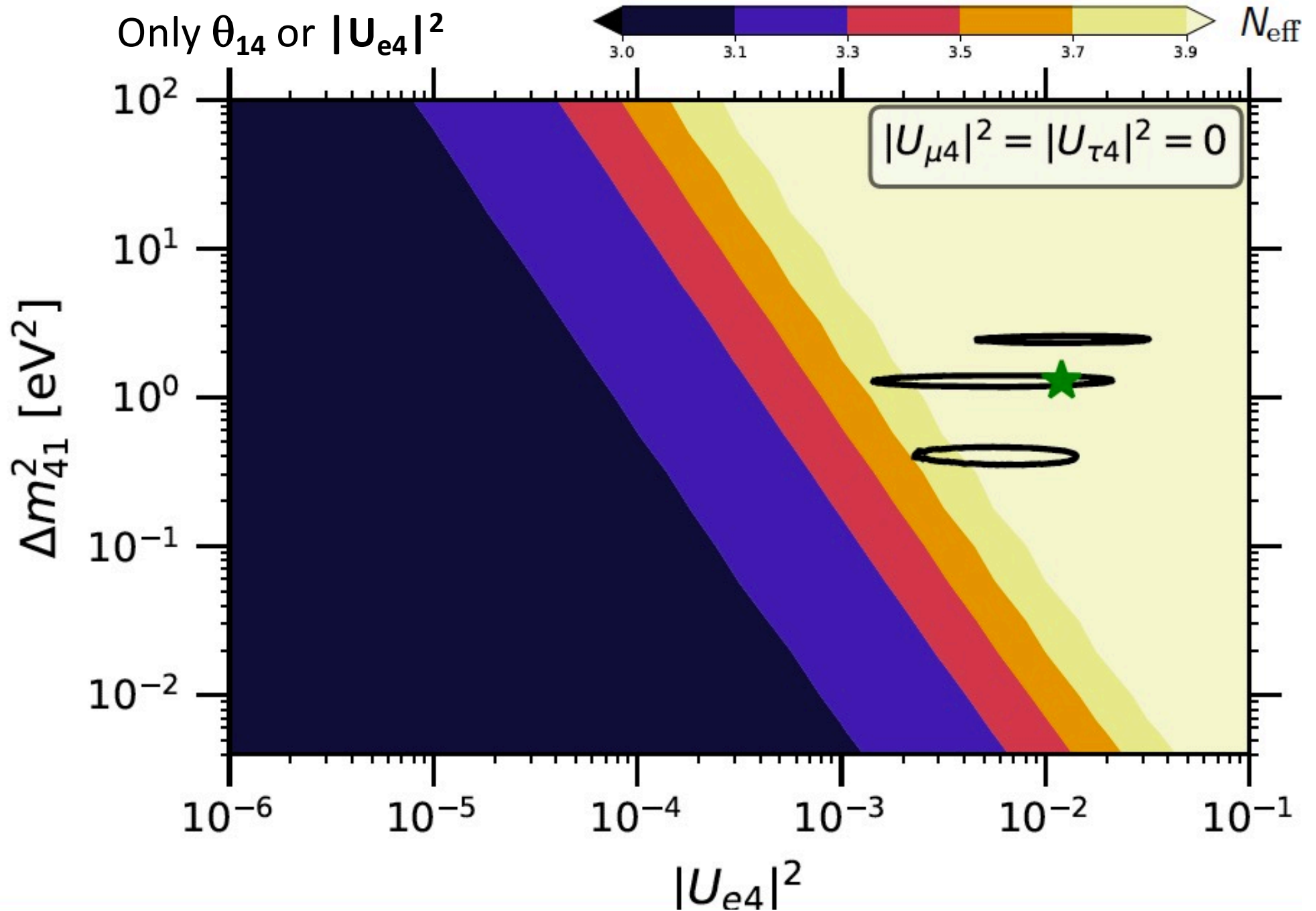
$$|U_{e4}|^2 = \sin^2 \theta_{14},$$

$$|U_{\mu4}|^2 = \cos^2 \theta_{14} \sin^2 \theta_{24},$$

$$|U_{\tau4}|^2 = \cos^2 \theta_{14} \cos^2 \theta_{24} \sin^2 \theta_{34},$$

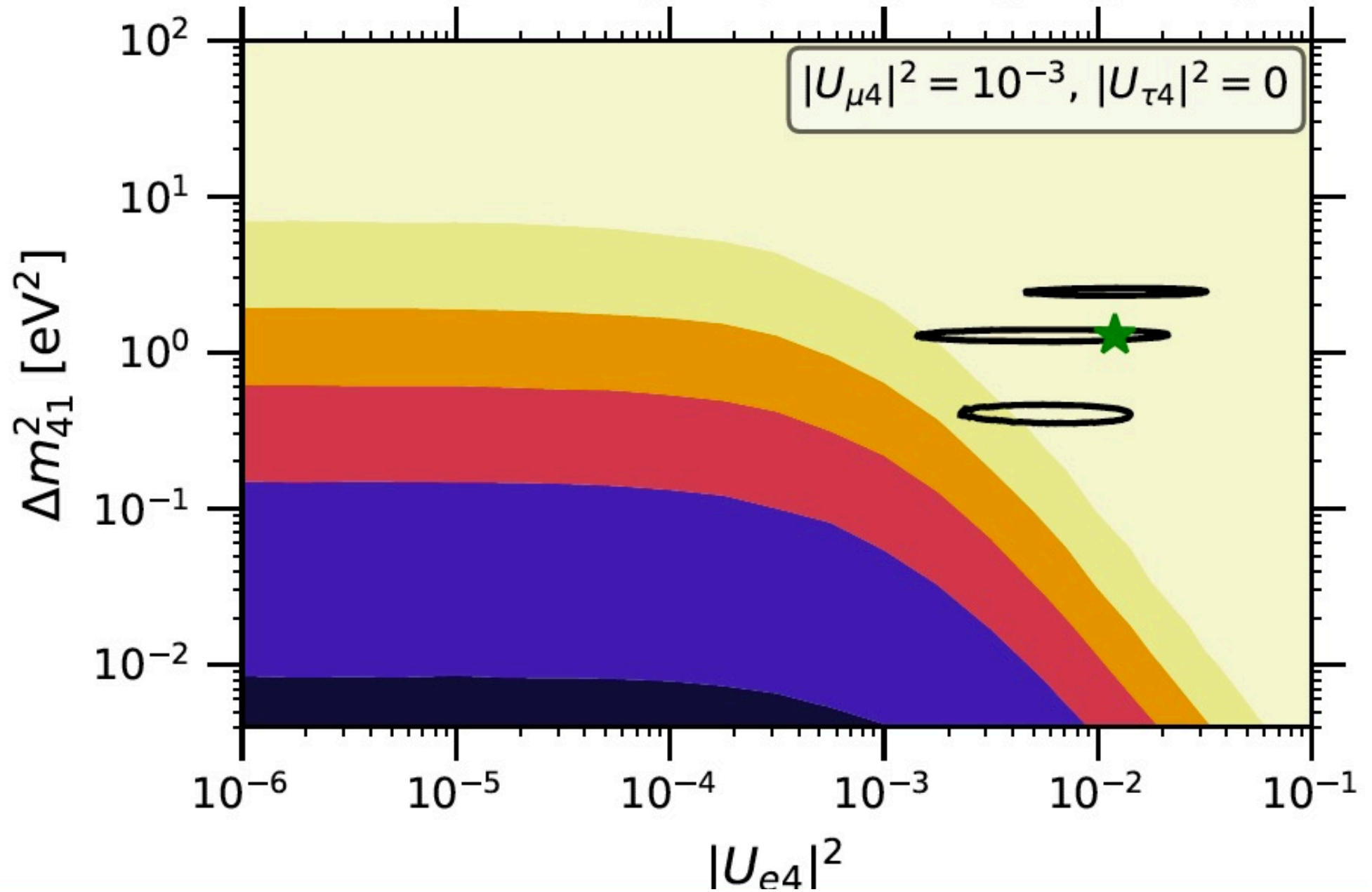
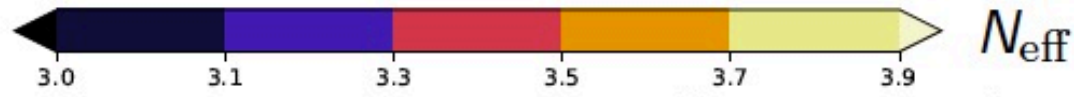
$$|U_{s4}|^2 = \cos^2 \theta_{14} \cos^2 \theta_{24} \cos^2 \theta_{34}.$$

Results: final value of N_{eff} and sterile mixing parameters



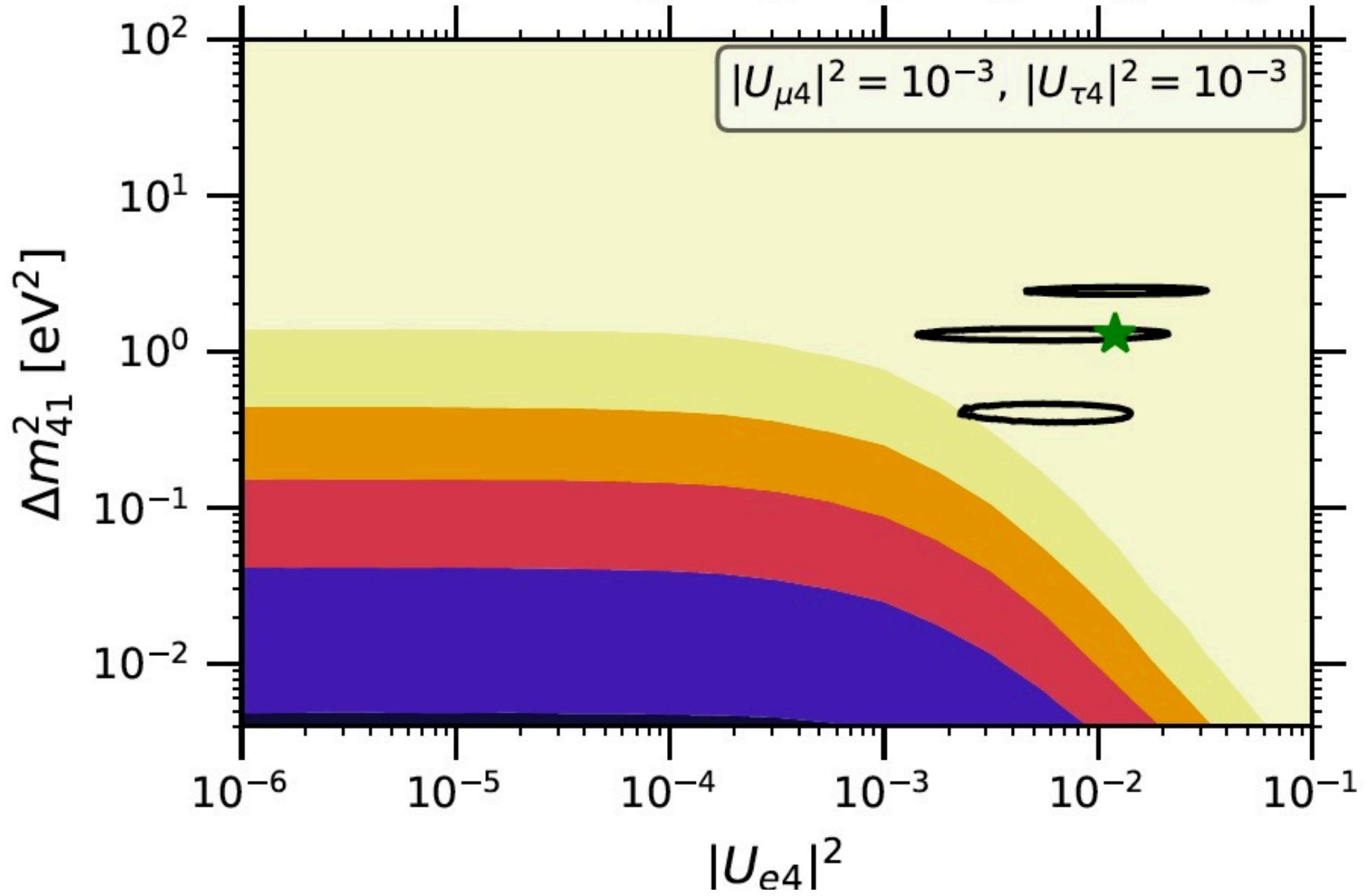
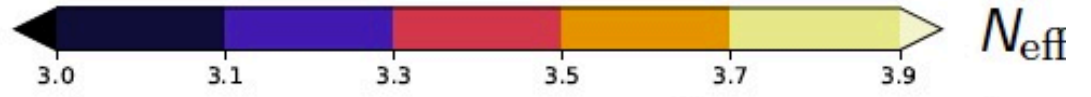
Results: final value of N_{eff} and sterile mixing parameters

We can vary more than one angle:



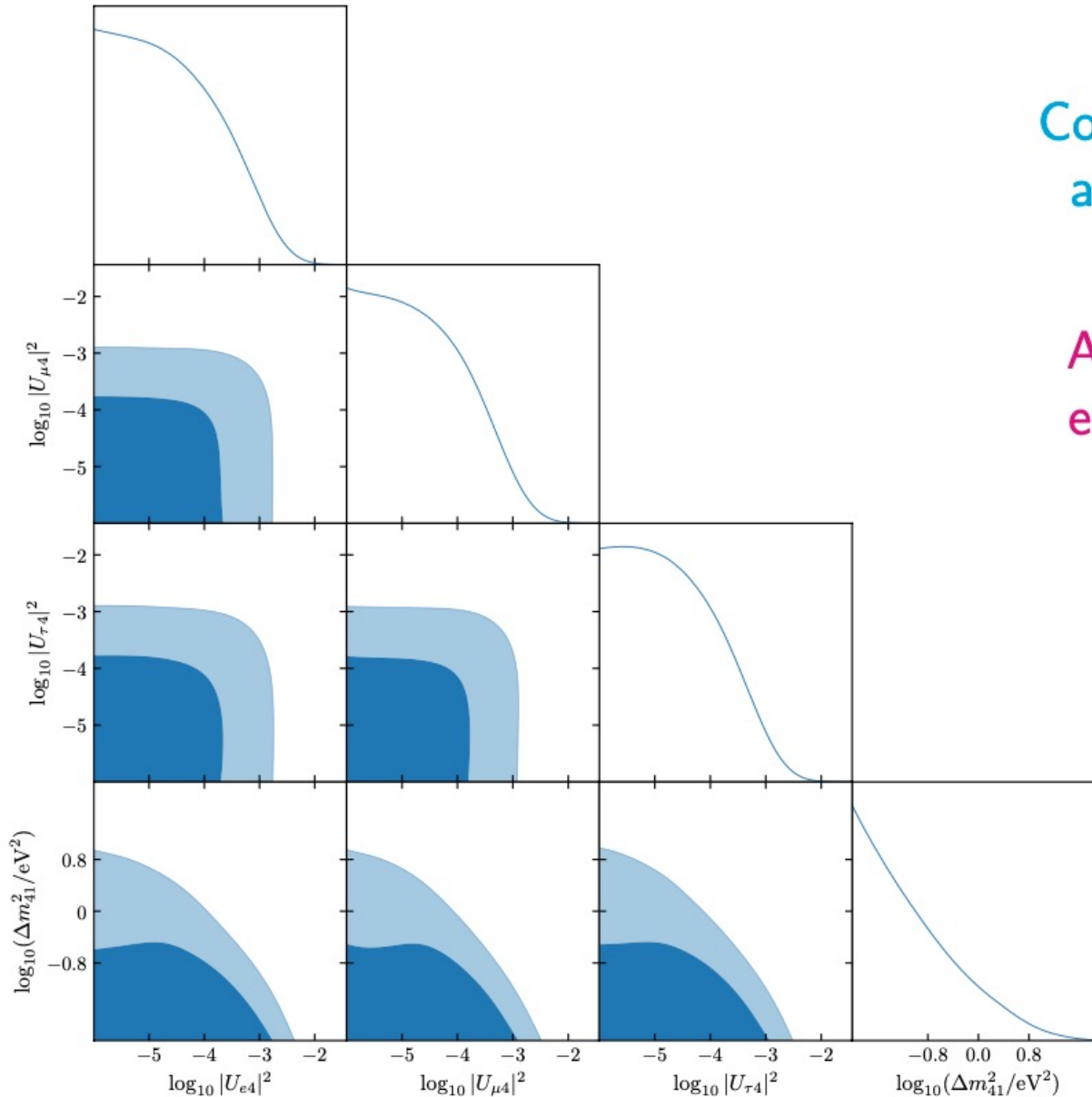
Results: final value of N_{eff} and sterile mixing parameters

We can vary more than one angle:



Cosmological bounds on active-sterile mixing parameters

Use multi-angle results from FortEPiA_{NO} to derive constraints on $|U_{\alpha 4}|^2$:



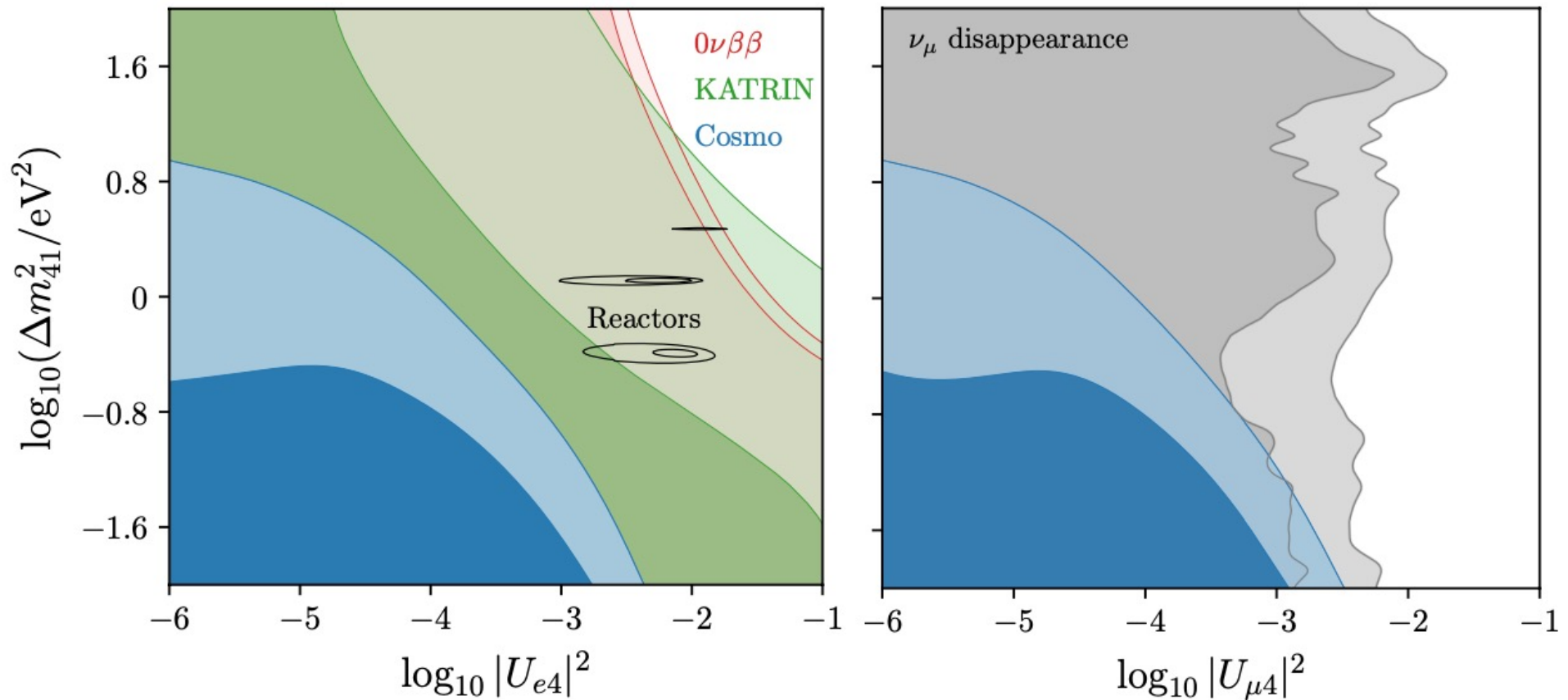
Constraints come from N_{eff}
and late-time density Ω_s

Angles $|U_{\alpha 4}|^2$ are almost
equivalent for cosmology

Bounds on active-sterile mixing parameters

Cosmological constraints are stronger than most other probes

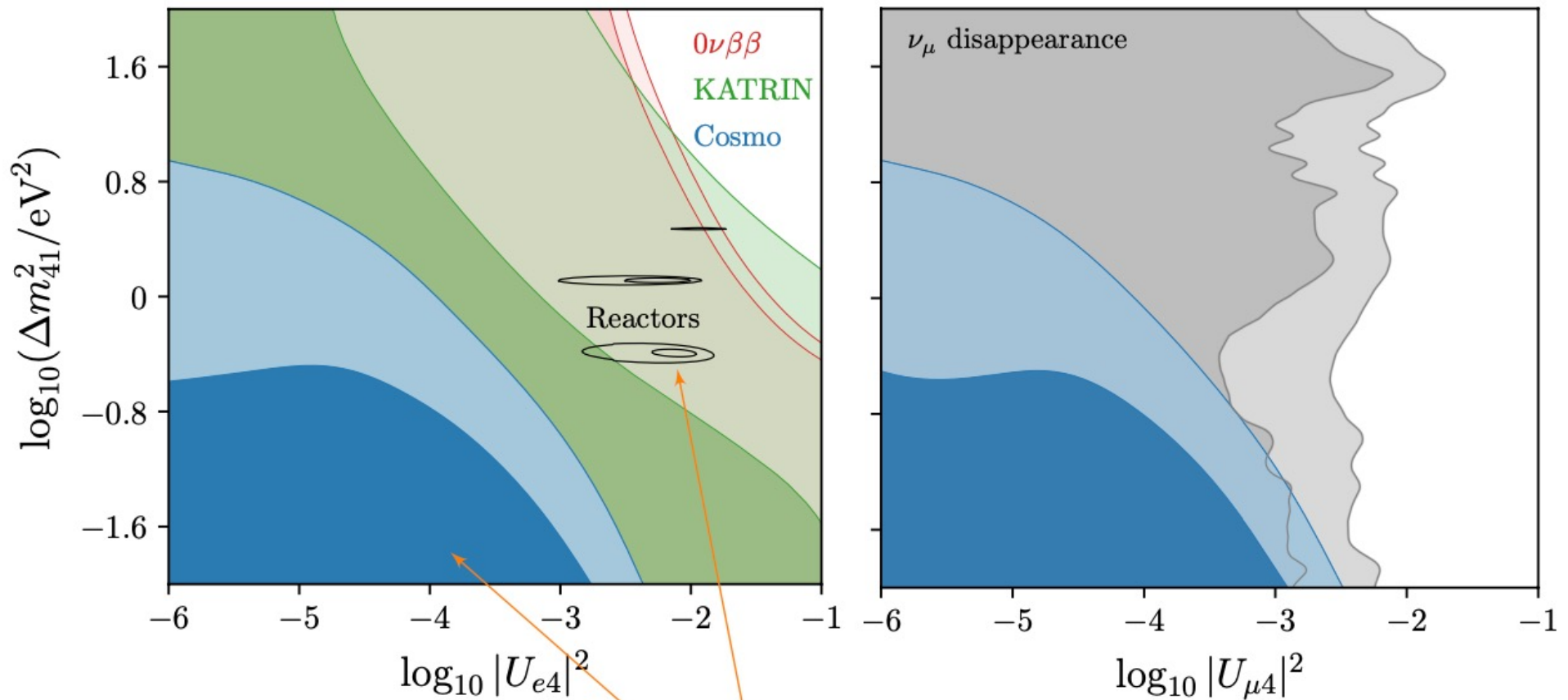
But much more model dependent (as all the cosmological constraints)!



Bounds on active-sterile mixing parameters

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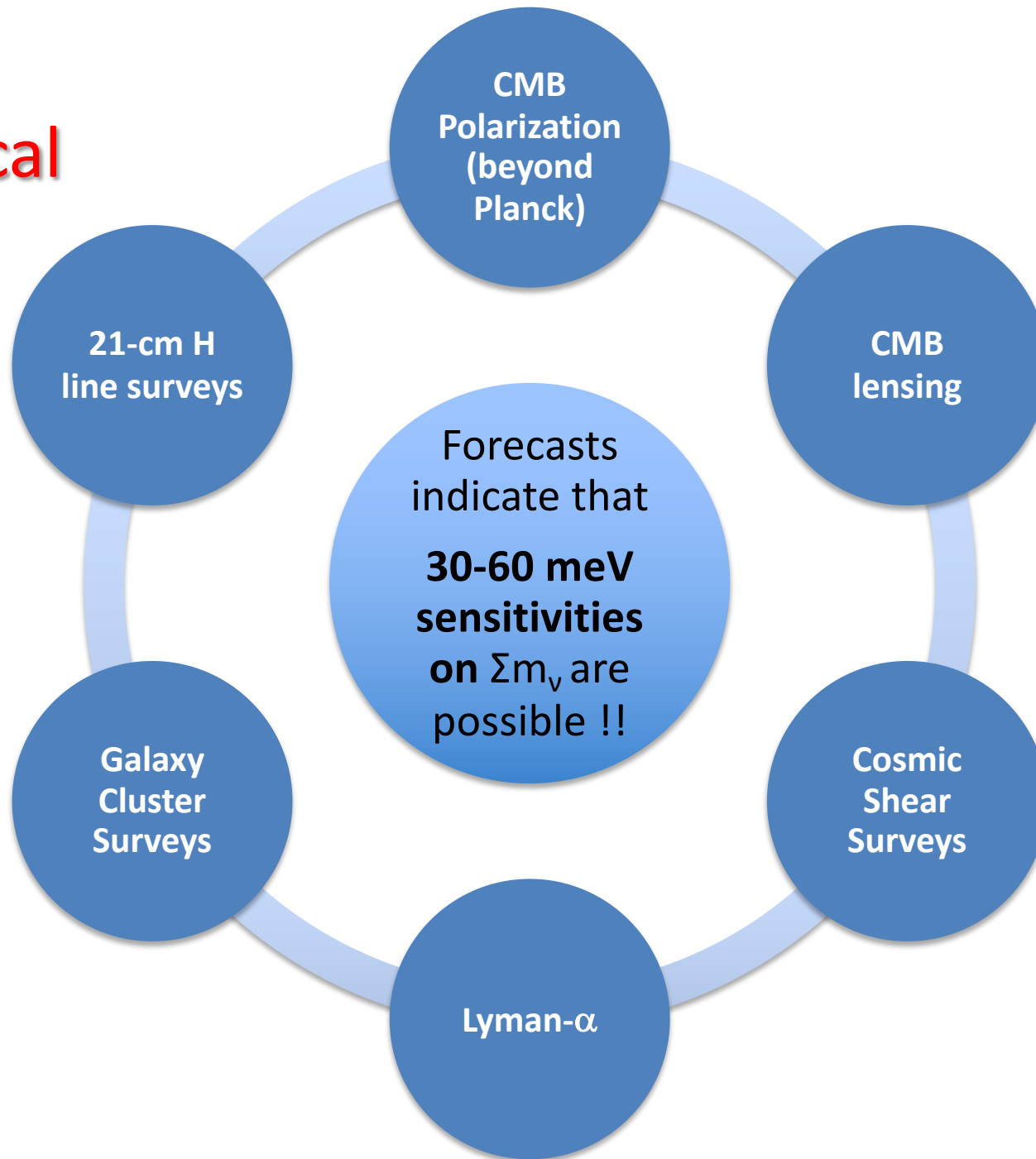


S Hagstotz et al, arXiv:2003.02289

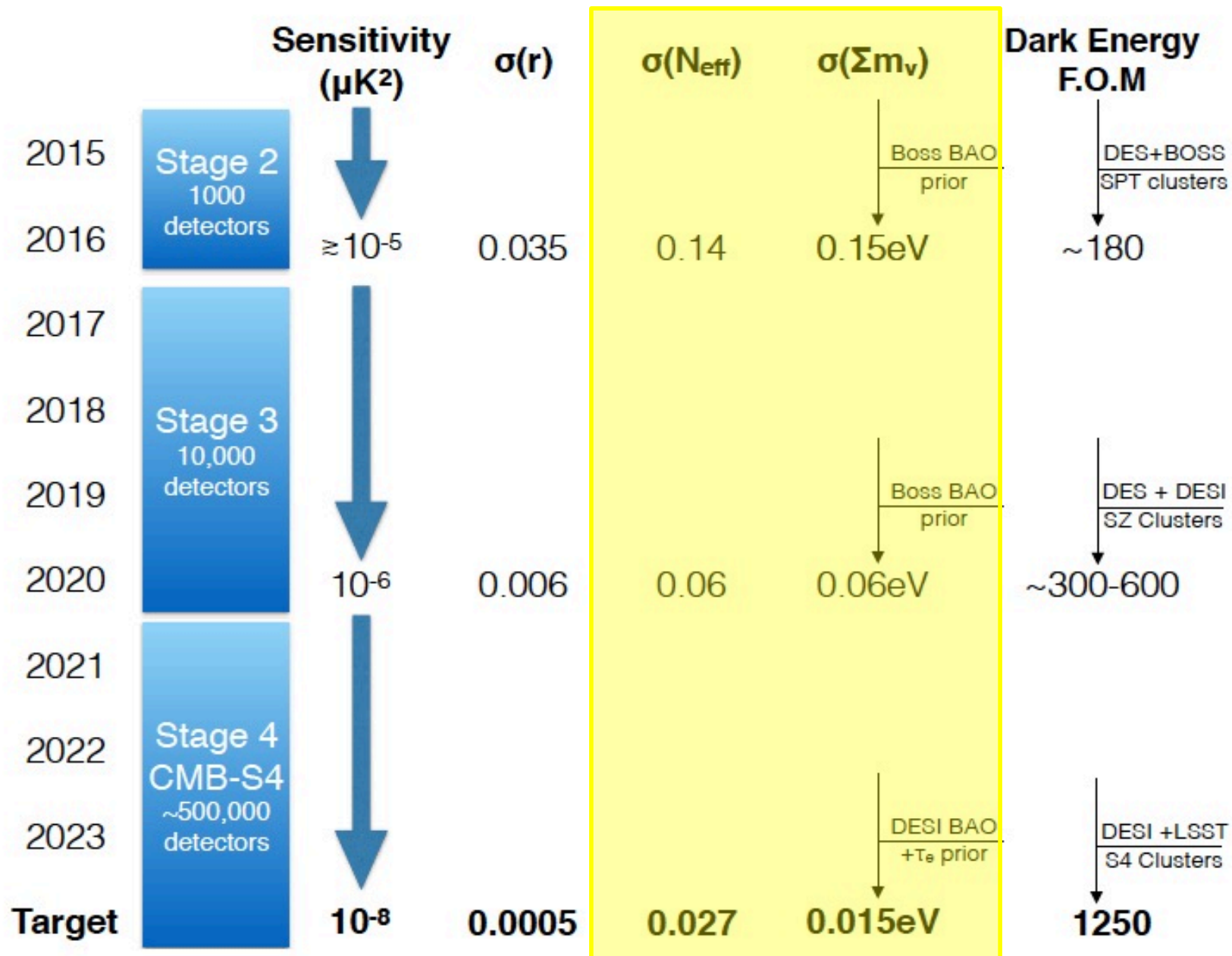
Warning: tension between reactor experiments and CMB bounds!

Future sensitivities on neutrino physics from cosmology

Future cosmological data



Future sensitivities on N_{eff} and neutrino masses



End