

# CNB essentials 2: Basics of neutrino cosmology, neutrino decoupling



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EuCAPT AstroNu  
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Prague, 21 Sep 2021



# Where do neutrinos come from?

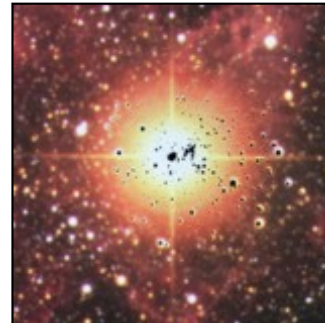
✓ Nuclear reactors



Sun



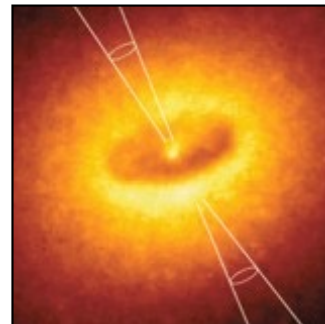
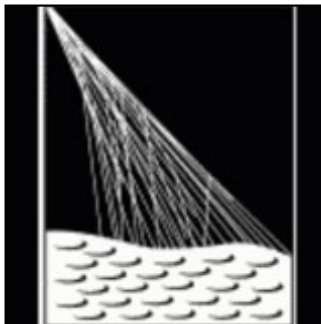
✓ Particle accelerators



Supernovae

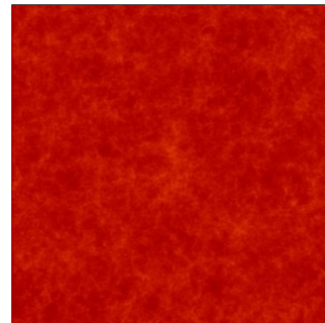
SN 1987A ✓

✓ Earth Atmosphere  
(Cosmic rays)



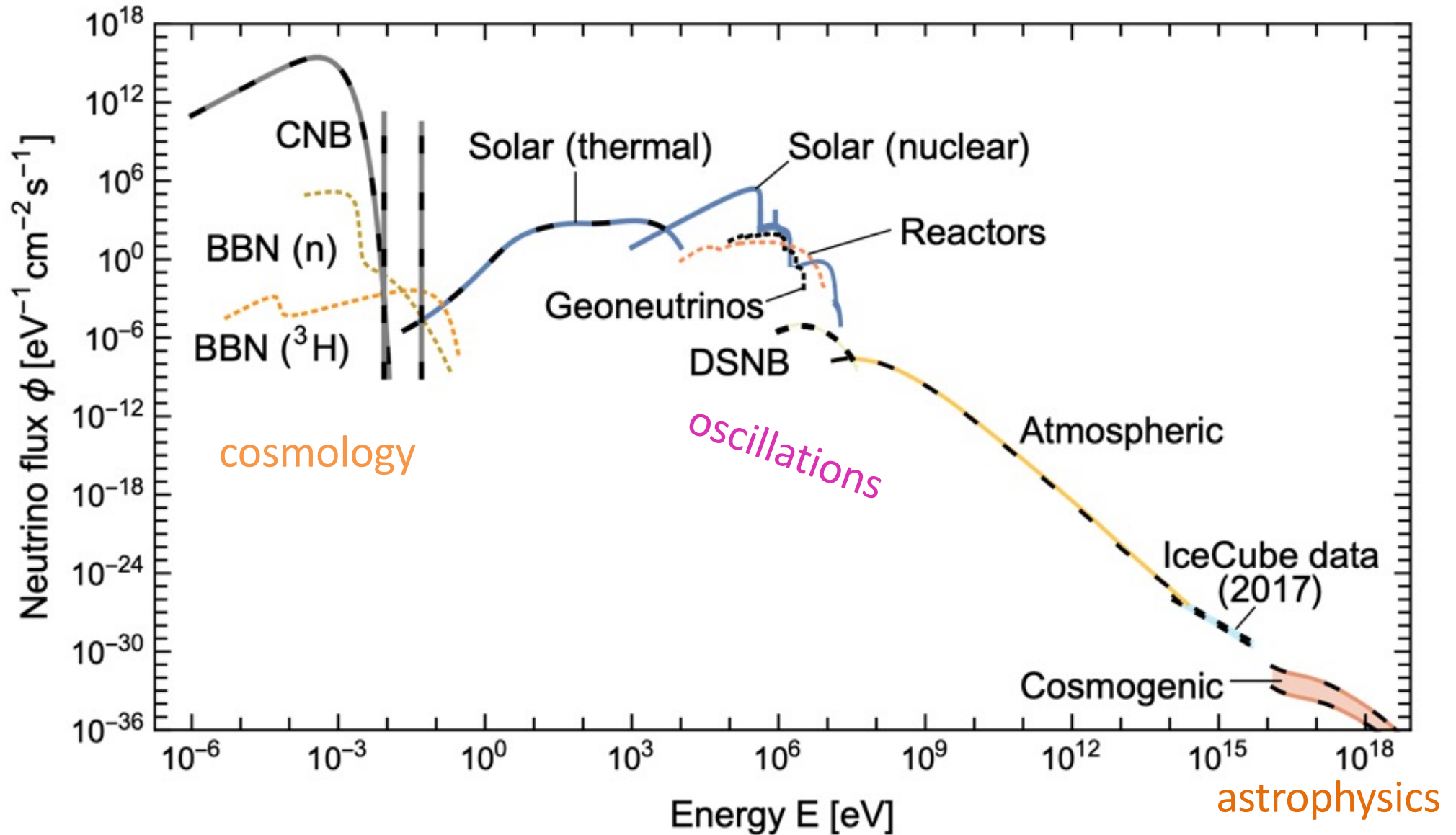
Accelerators in  
astrophysical sources ? ✓

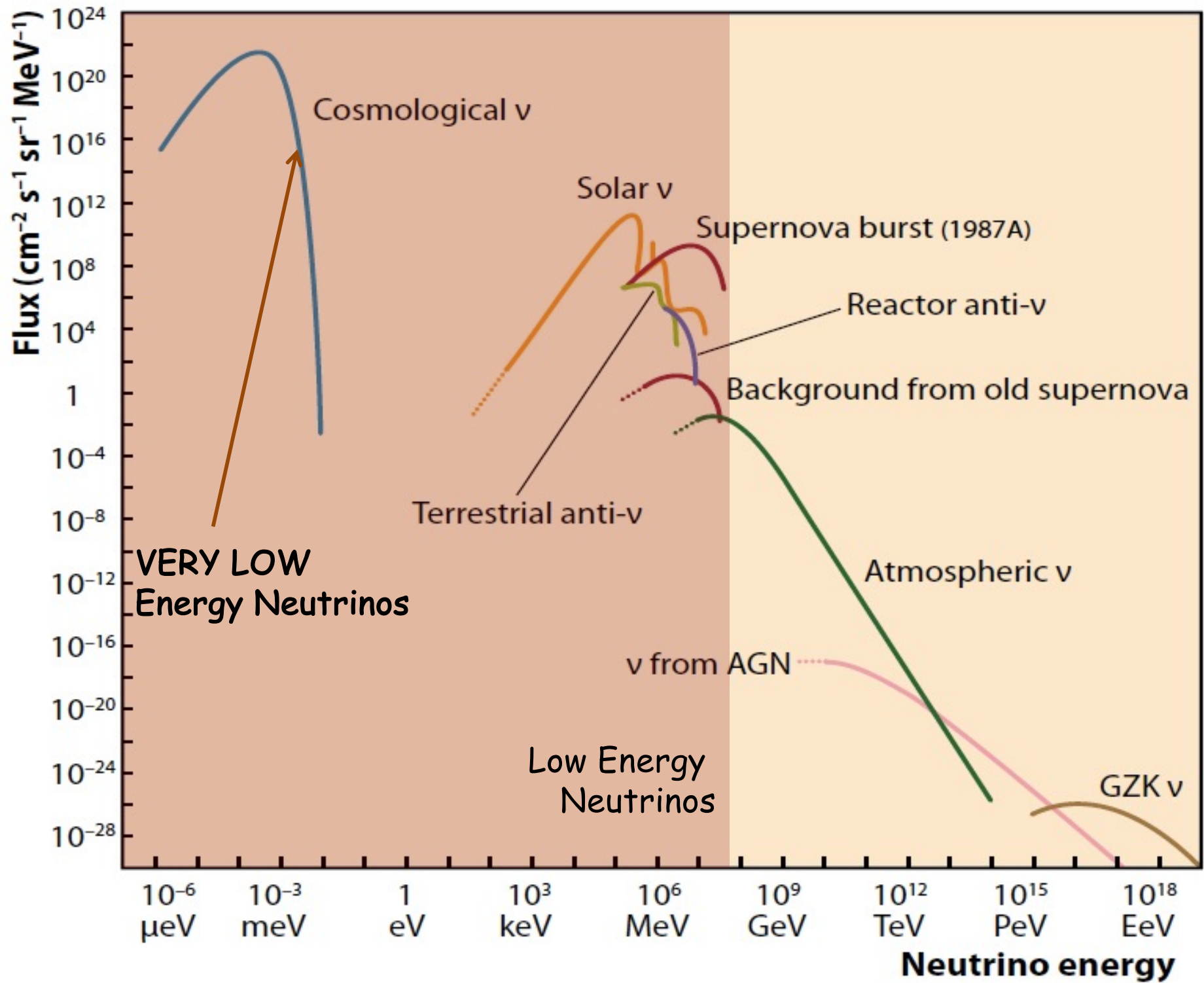
✓ Earth interior  
(Natural Radioactivity)



Early Universe  
(today  $336 \nu/\text{cm}^3$ )  
Indirect evidence

# Grand Unified Neutrino Spectrum at Earth





# Outline



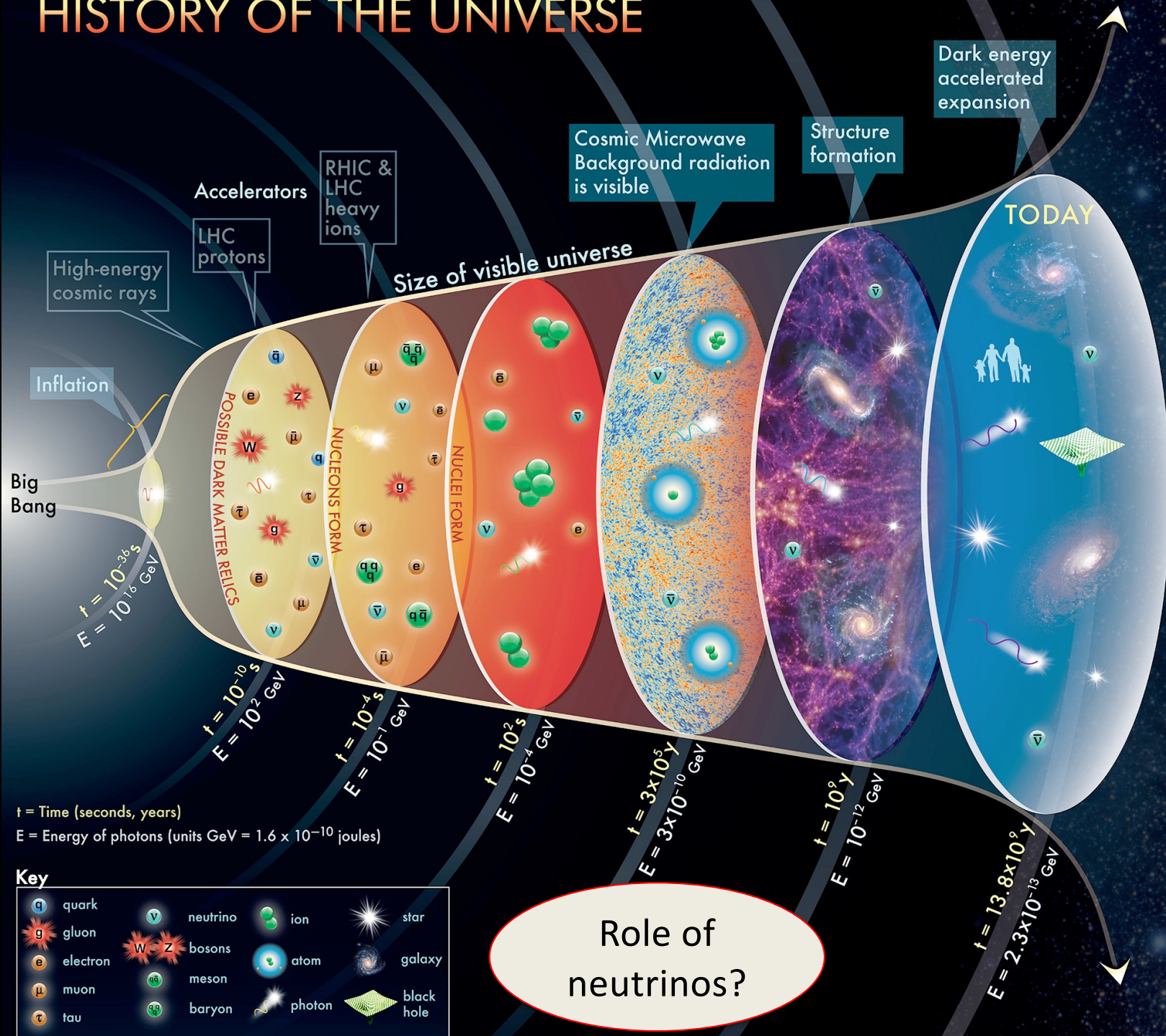
**Introduction: neutrinos and the history of the Universe**

**Production and decoupling of relic neutrinos**

**The radiation content of the Universe ( $N_{\text{eff}}$ )**

# **Introduction: neutrinos and the history of the Universe**

# HISTORY OF THE UNIVERSE



t = Time (seconds, years)  
 E = Energy of photons (units GeV =  $1.6 \times 10^{-10}$  joules)

**Key**

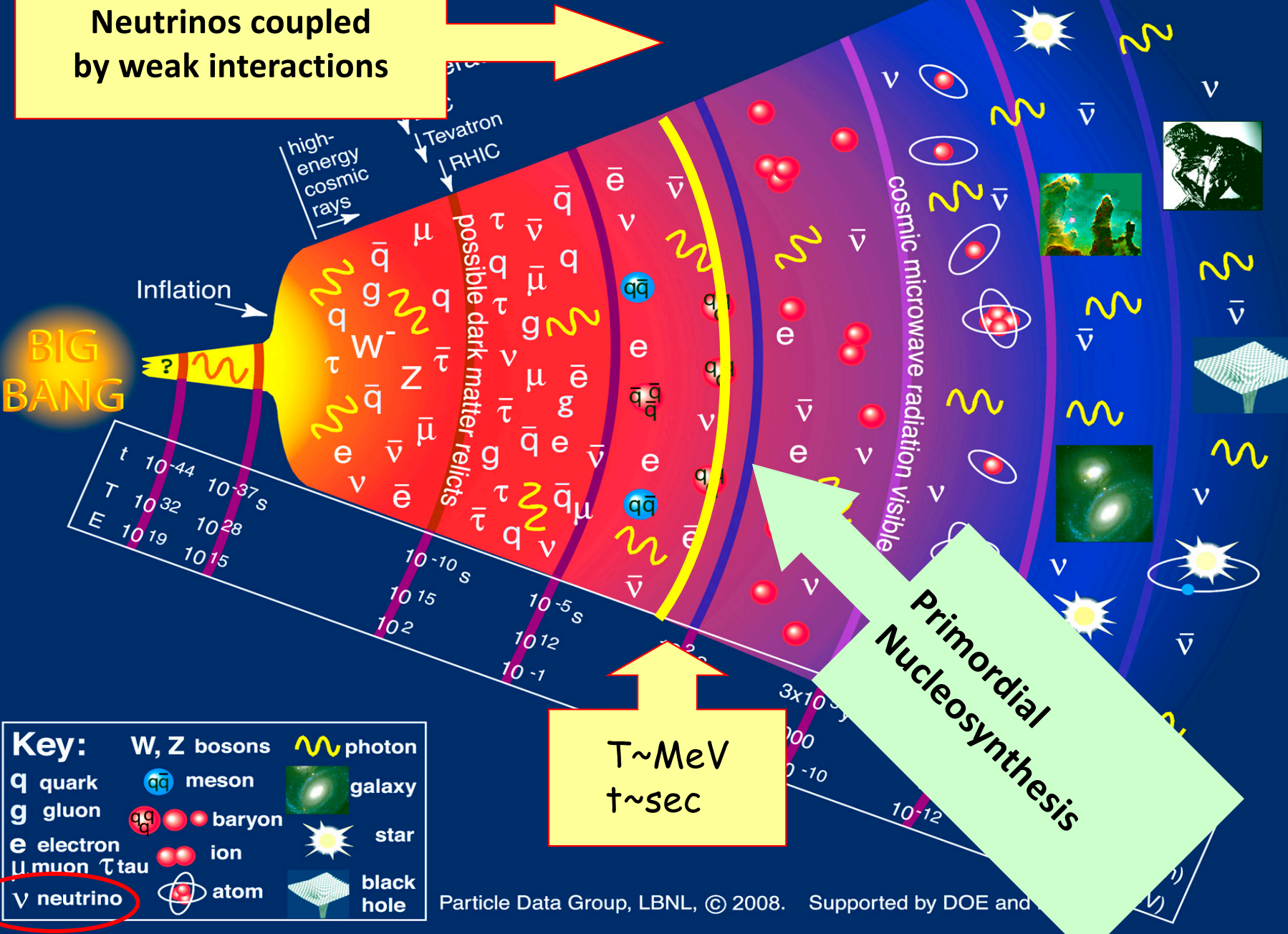
quark	neutrino	ion	star
gluon	bosons	atom	galaxy
electron	meson	photon	black hole
muon	baryon		
tau			

Role of neutrinos?

The concept for the above figure originated in a 1986 paper by Michael Turner.

# History of the Universe

Neutrinos coupled by weak interactions



high-energy cosmic rays  
Tevatron  
RHIC

Inflation

BIG BANG

t	10 <sup>-44</sup>	10 <sup>-37</sup> s
T	10 <sup>32</sup>	10 <sup>28</sup>
E	10 <sup>19</sup>	10 <sup>15</sup>

	10 <sup>-10</sup> s	
	10 <sup>-5</sup> s	
	10 <sup>-1</sup>	

T ~ MeV  
t ~ sec

Primordial Nucleosynthesis

**Key:**

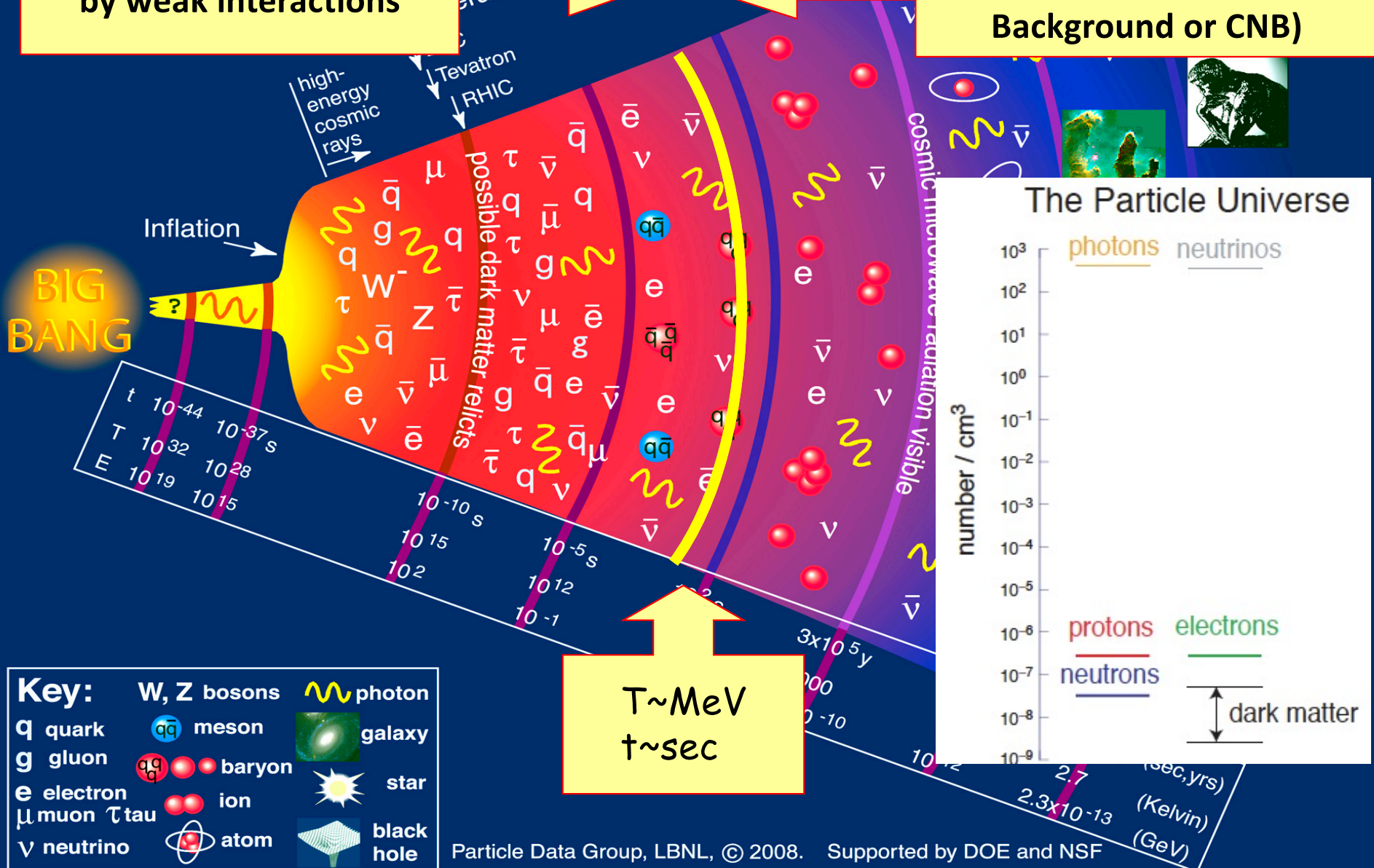
W, Z bosons	photon
q quark	meson
g gluon	baryon
e electron	ion
μ muon	τ tau
ν neutrino	atom
	galaxy
	star
	black hole



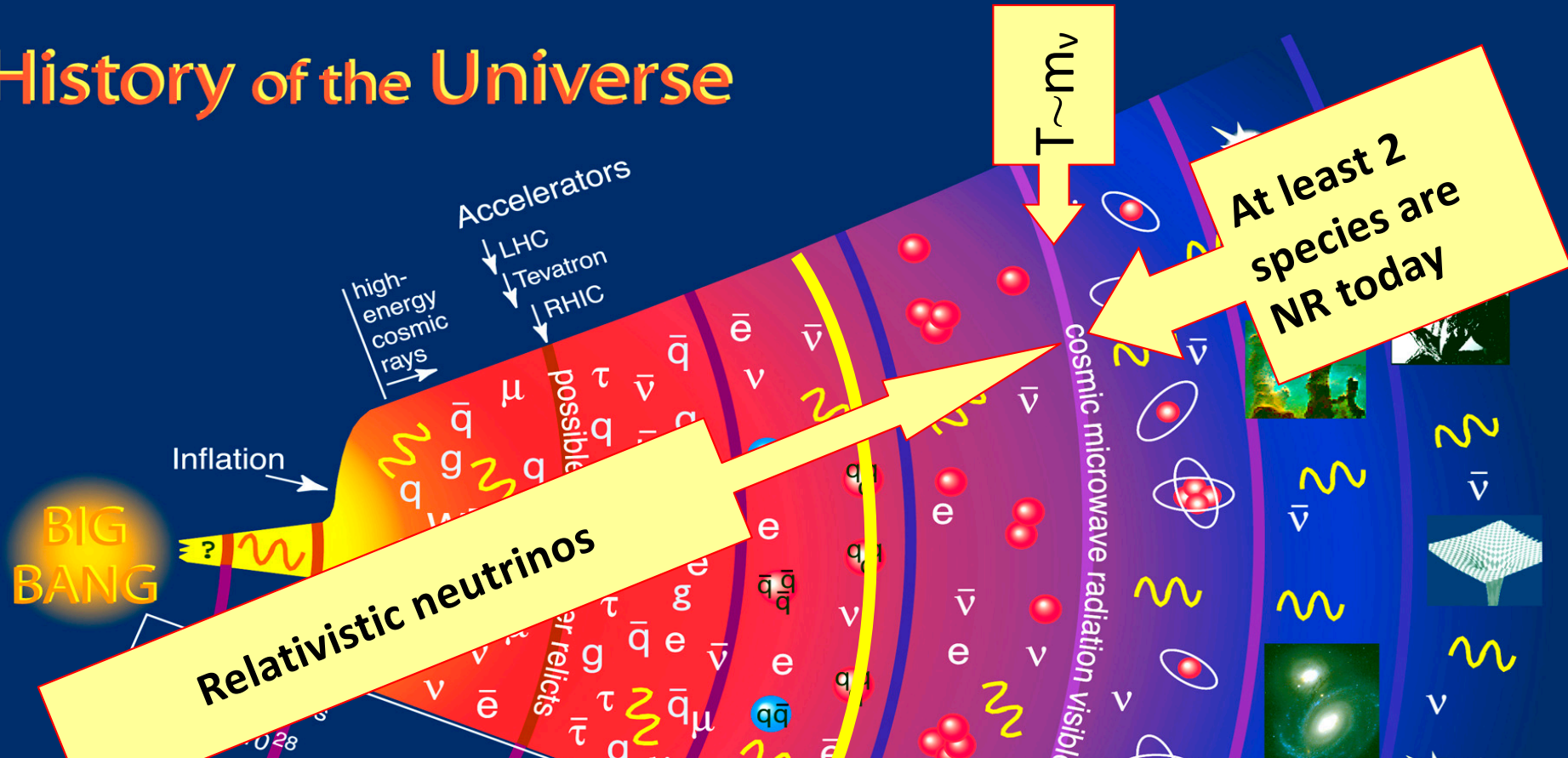
# History of the Universe

Neutrinos coupled by weak interactions

Decoupled neutrinos (Cosmic Neutrino Background or CNB)



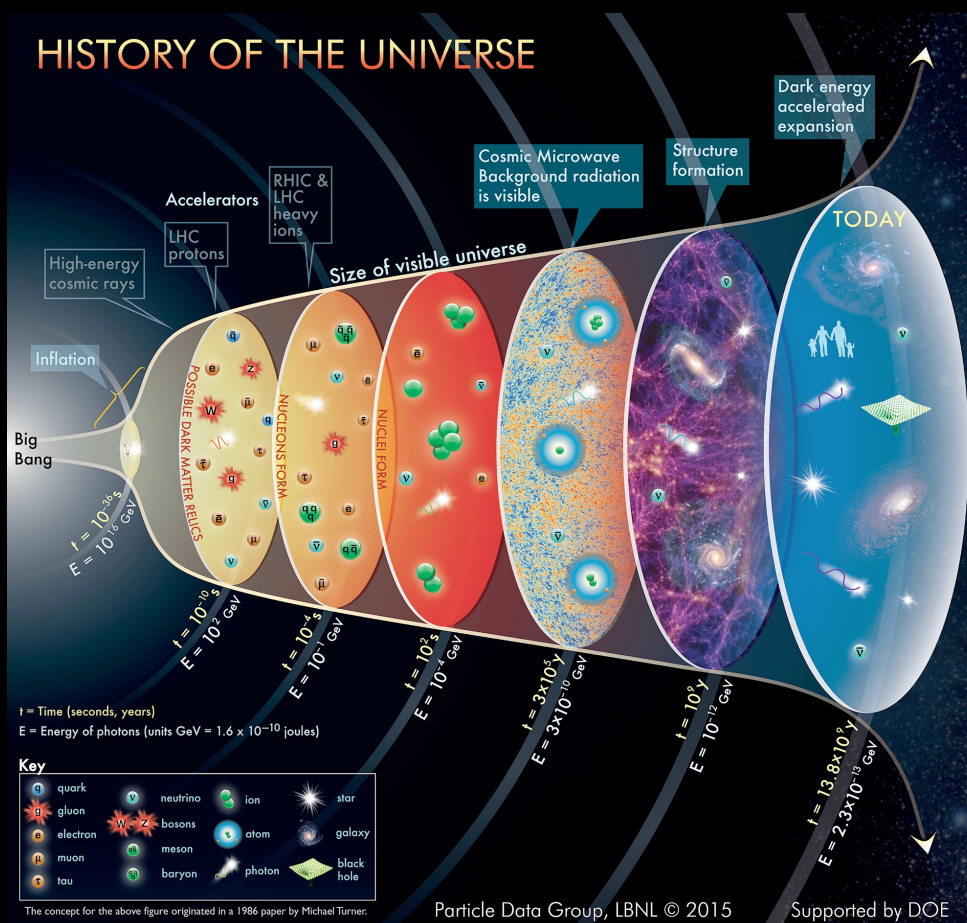
# History of the Universe



Neutrino cosmology is interesting because Relic neutrinos are very abundant:

- The CNB contributes to radiation at early times and to matter at late times (info on the number of neutrinos and their masses)
- Cosmological observables can be used to test standard or non-standard neutrino properties

$\mu$  muon  $\tau$  tau  
 $\nu$  neutrino atom black hole



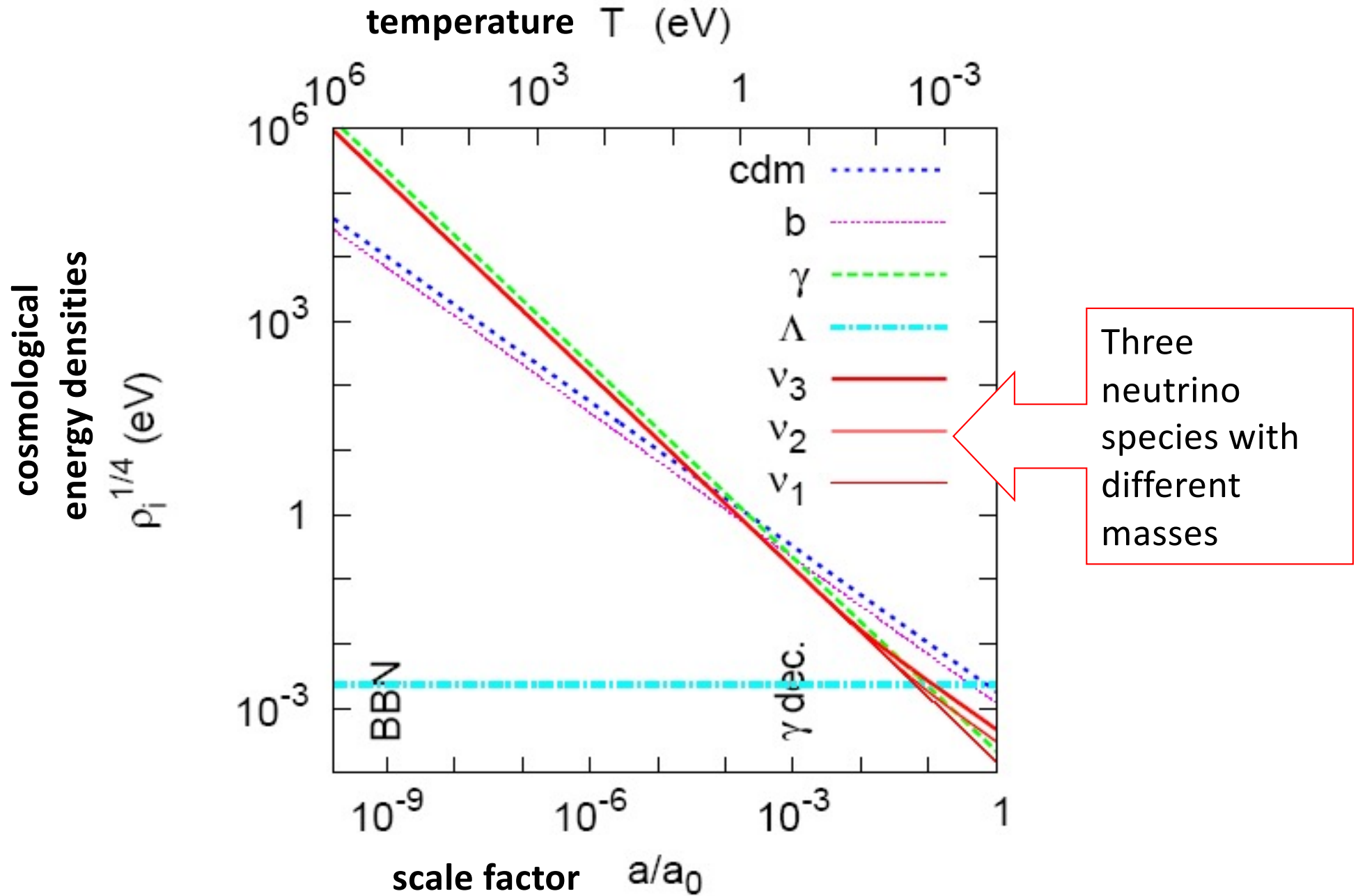
energy density:  $\rho(a) = a^{-3(1+w)}$

$$\rho_R \sim a^{-4} \quad , \quad w = 1/3 \quad \text{(Radiation)}$$

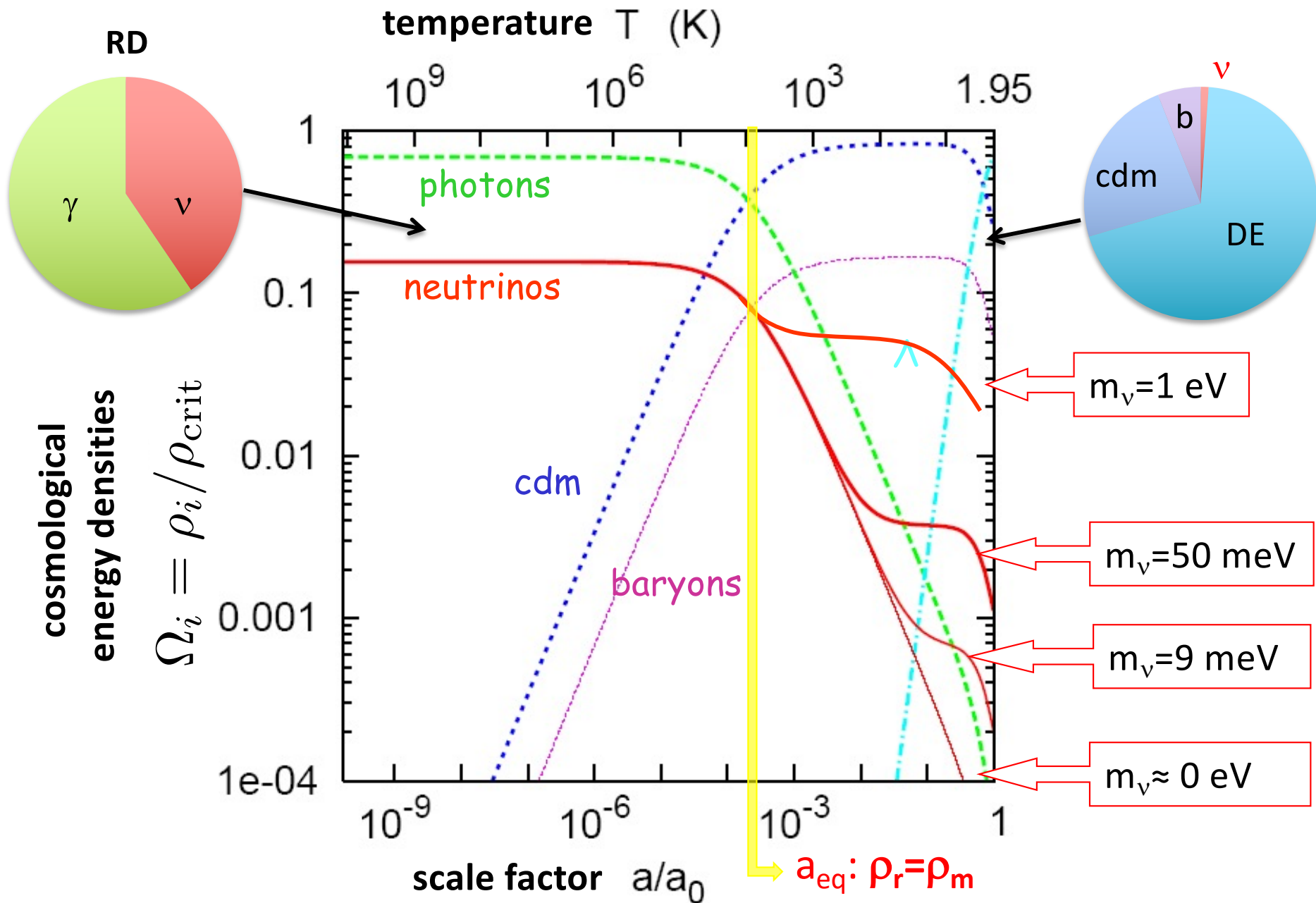
$$\rho_M \sim a^{-3} \quad , \quad w = 0 \quad \text{(Matter)}$$

$$\rho_\Lambda \sim \text{const.} \quad , \quad w = -1 \quad \text{(Cosmological constant)}$$

# Evolution of the background densities: 1 MeV $\rightarrow$ now



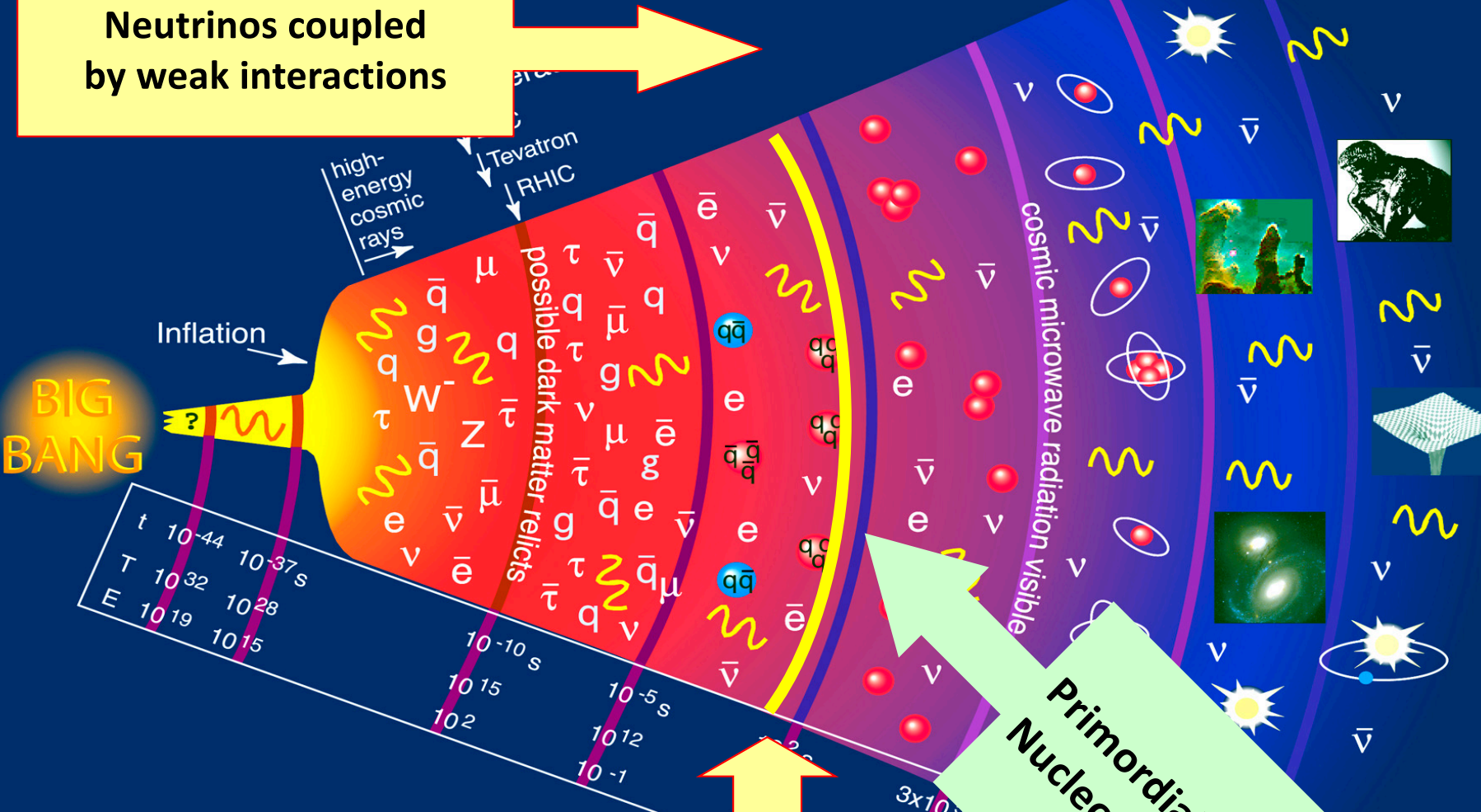
# Evolution of the background densities: 1 MeV $\rightarrow$ now



# **Production and decoupling of relic neutrinos**

# History of the Universe

Neutrinos coupled by weak interactions



t	10 <sup>-44</sup>	10 <sup>-37</sup> s
T	10 <sup>32</sup>	10 <sup>28</sup>
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	10 <sup>-10</sup> s	
	10 <sup>-5</sup> s	
	10 <sup>-1</sup>	

T ~ MeV  
t ~ sec

**Key:**

W, Z bosons	photon
q quark	meson
g gluon	baryon
e electron	ion
μ muon	τ tau
ν neutrino	atom
	galaxy
	star
	black hole

# Equilibrium thermodynamics

Distribution function of particle momenta in equilibrium

$$f_i^{eq}(p, T) = \left[ \exp \left( \frac{E_i - \mu_i}{T} \right) \mp 1 \right]^{-1}$$

Thermodynamical variables

VARIABLE	RELATIVISTIC		NON REL.
	BOSE	FERMI	
$n$	$\frac{\zeta(3)}{\pi^2} g T^3$	$\frac{3 \zeta(3)}{4 \pi^2} g T^3$	$g \left( \frac{mT}{2\pi} \right)^{3/2} e^{-m/T}$
$\rho$	$\frac{\pi^2}{30} g T^4$	$\frac{7 \pi^2}{830} g T^4$	$mn$
$p$	$\frac{\rho}{3}$		$nT \ll \rho$
$\langle E \rangle$	$2,701T$	$3,151T$	$m + \frac{3}{2}T$

Particles in equilibrium when T are high and interactions effective

$T \sim 1/a(t)$

$$n = g_i \int \frac{d^2 \vec{p}}{(2\pi)^3} f_i(p, T) \quad \rho = g_i \int \frac{d^2 \vec{p}}{(2\pi)^3} E_i f_i(p, T)$$

$$p = g_i \int \frac{d^2 \vec{p}}{(2\pi)^3} \frac{p^2}{3E_i} f_i(p, T) \quad \langle E \rangle = \rho/n$$



# Neutrinos in Equilibrium

$$1 \text{ MeV} \lesssim T \lesssim m_\mu$$

$$T_\nu = T_{e^\pm} = T_\gamma$$

$$\nu_\alpha \nu_\beta \leftrightarrow \nu_\alpha \nu_\beta$$

$$\nu_\alpha \bar{\nu}_\beta \leftrightarrow \nu_\alpha \bar{\nu}_\beta$$

$$\nu_\alpha e^\pm \leftrightarrow \nu_\alpha e^\pm$$

$$\nu_\alpha \bar{\nu}_\alpha \leftrightarrow e^+ e^-$$

$$\mathcal{L}_{\text{SM}} = -2\sqrt{2}G_F \left\{ (\bar{\nu}_e \gamma^\mu L \nu_e) (\bar{e} \gamma_\mu L e) + \sum_{P,\alpha} g_P (\bar{\nu}_\alpha \gamma^\mu L \nu_\alpha) (\bar{e} \gamma_\mu P e) \right\}$$

$$P = L, R = (1 \mp \gamma_5)/2$$

$$g_L = -\frac{1}{2} + \sin^2 \theta_W \text{ and } g_R = \sin^2 \theta_W$$

# Cosmological energy densities: radiation

Energy density of **relativistic particles** with  $f_i(p)$

$$\rho_i = g \int \frac{d^3 p}{(2\pi)^3} \frac{p}{e^{p/T_i} \pm 1}$$

$$\rho_i = 3P_i = \begin{cases} \frac{\pi^2}{30} g T_i^4, & \text{boson} \\ \frac{7}{8} \frac{\pi^2}{30} g T_i^4, & \text{fermion} \end{cases}$$
$$\rho_\gamma = \frac{\pi^2}{15} T_\gamma^4$$
$$\rho_\nu = 3 \times \frac{7}{8} \times \frac{\pi^2}{15} T_\nu^4$$

# Neutrino decoupling

As the Universe expands, particle densities are diluted and temperatures fall. Weak interactions become **ineffective** to keep neutrinos in good thermal contact with the e.m. plasma

Rough, but quite accurate estimate of the decoupling temperature

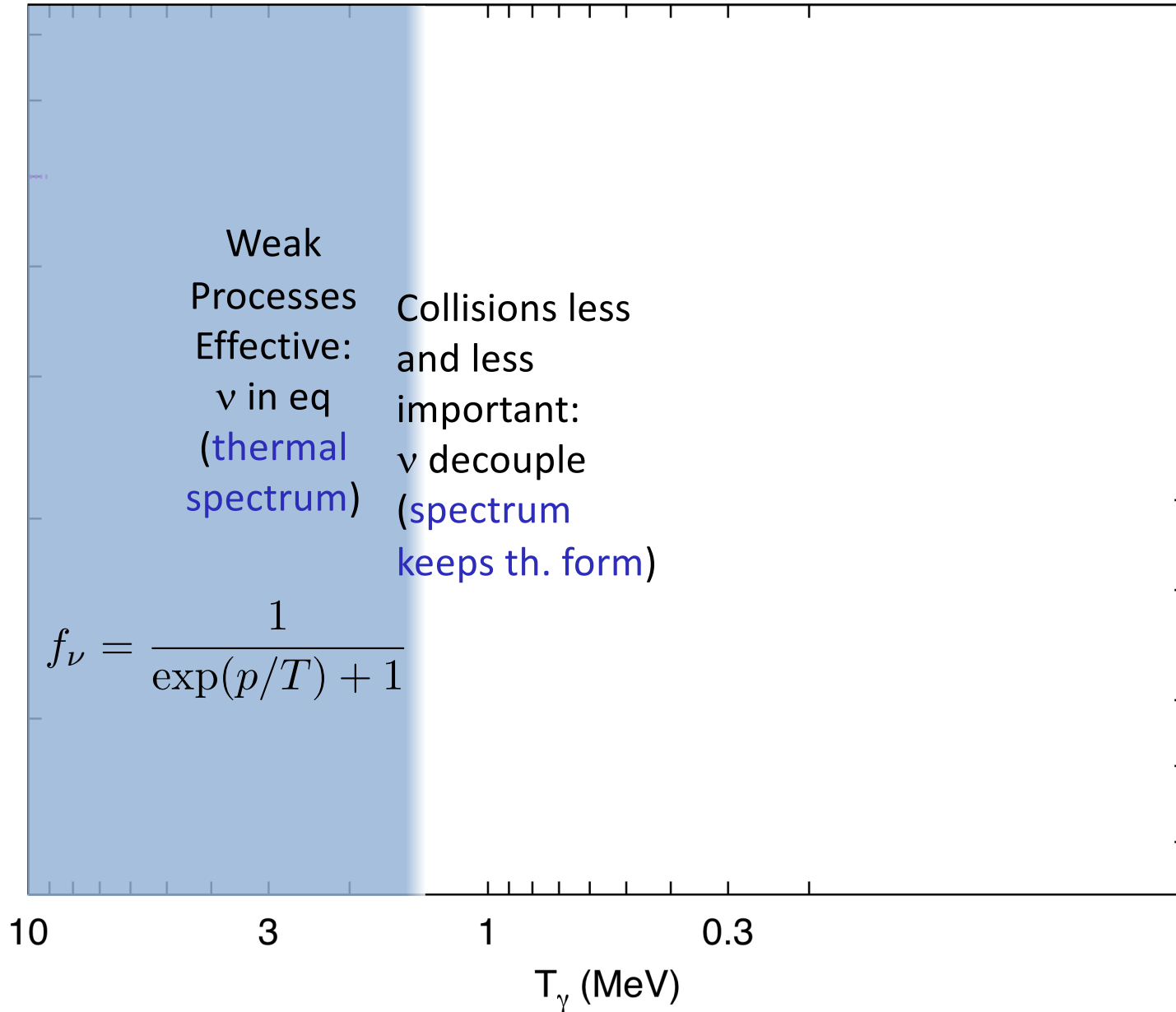
Rate of weak processes  $\sim$  Hubble expansion rate

$$\Gamma_W \approx \sigma_W |v| n, \quad H^2 = \frac{8\pi\rho_{\text{rad}}}{3M_P^2} \rightarrow G_F^2 T^5 \approx \sqrt{\frac{8\pi\rho_{\text{rad}}}{3M_P^2}} \rightarrow T_{\text{dec}}(\nu) \approx 1 \text{ MeV}$$

Since  $\nu_e$  have both CC and NC interactions with  $e^\pm$

$$T_{\text{dec}}(\nu_e) \simeq 2 \text{ MeV} \quad T_{\text{dec}}(\nu_{\mu,\tau}) \simeq 3 \text{ MeV}$$

# Neutrino decoupling



Expansion of the Universe

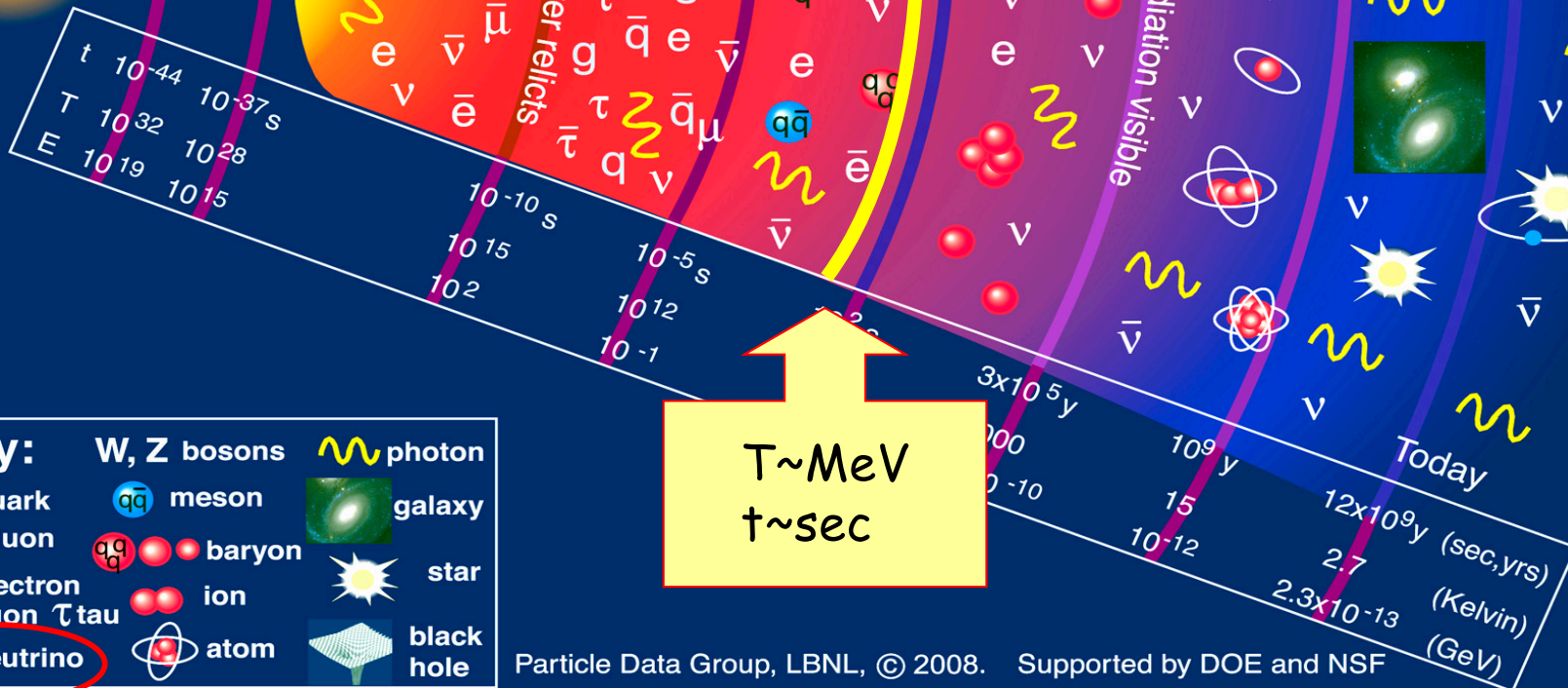


# History of the Universe

Neutrinos coupled  
by weak interactions

$$f_\nu(p, T) = \frac{1}{\exp(p/T) + 1}$$

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**Key:**

W, Z bosons		photon	
q quark		meson	
g gluon		baryon	
e electron		ion	
$\mu$ muon		atom	
$\tau$ tau		star	
$\nu$ neutrino		black hole	
		galaxy	

$T \sim \text{MeV}$   
 $t \sim \text{sec}$

# History of the Universe

Neutrinos coupled by weak interactions

Free-streaming neutrinos (decoupled): Cosmic Neutrino Background

$$f_\nu(p, T) = \frac{1}{\exp(p/T) + 1}$$

Neutrinos keep the energy spectrum of a relativistic fermion with eq form

$T \sim \text{MeV}$   
 $t \sim \text{sec}$

BIG BANG

Inflation

t	10 <sup>-44</sup>	10 <sup>-37</sup> s
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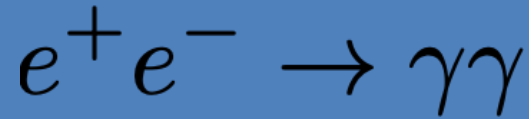
possible dark matter relics

**Key:**

W, Z bosons	photon
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ν neutrino	atom
	black hole

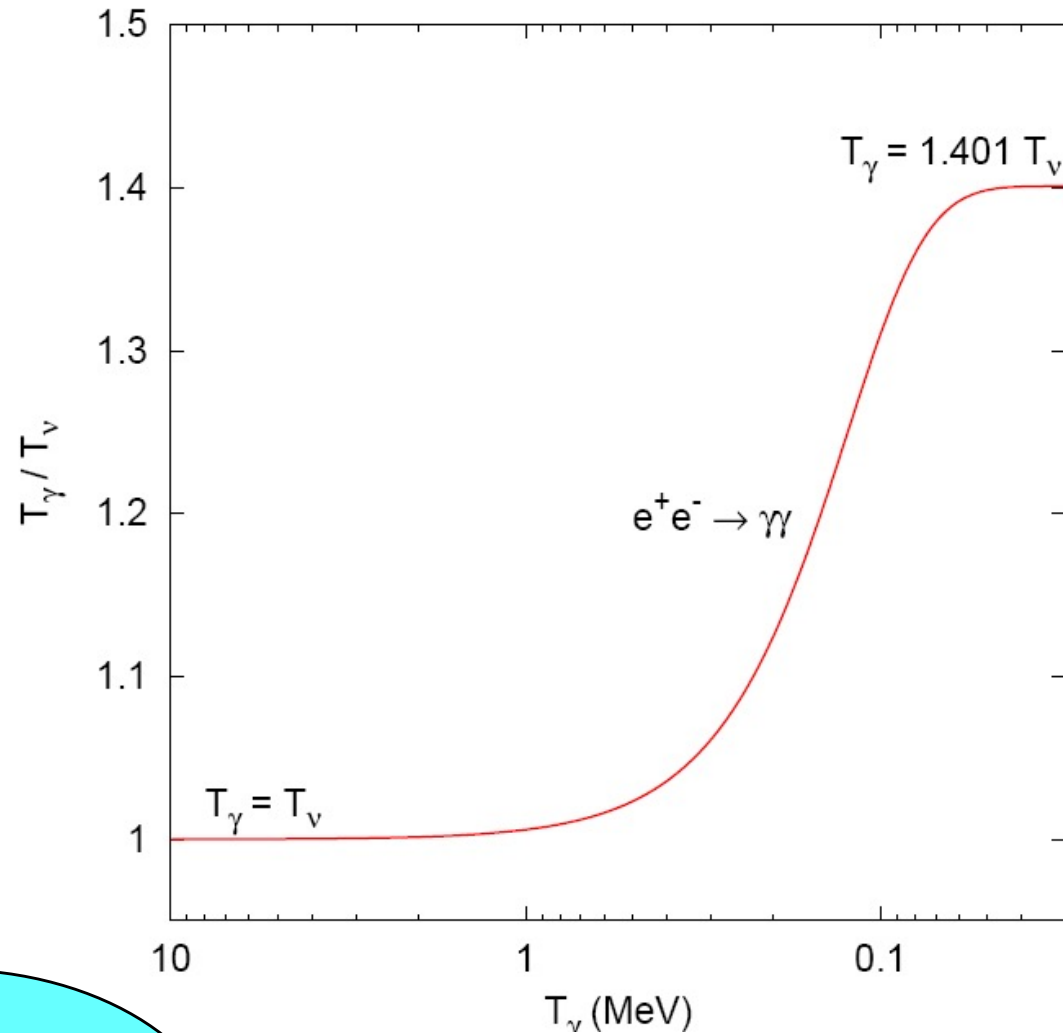
# Neutrino and photon (CMB) temperatures

At  $T \sim m_e$ ,  
electron-  
positron pairs  
annihilate



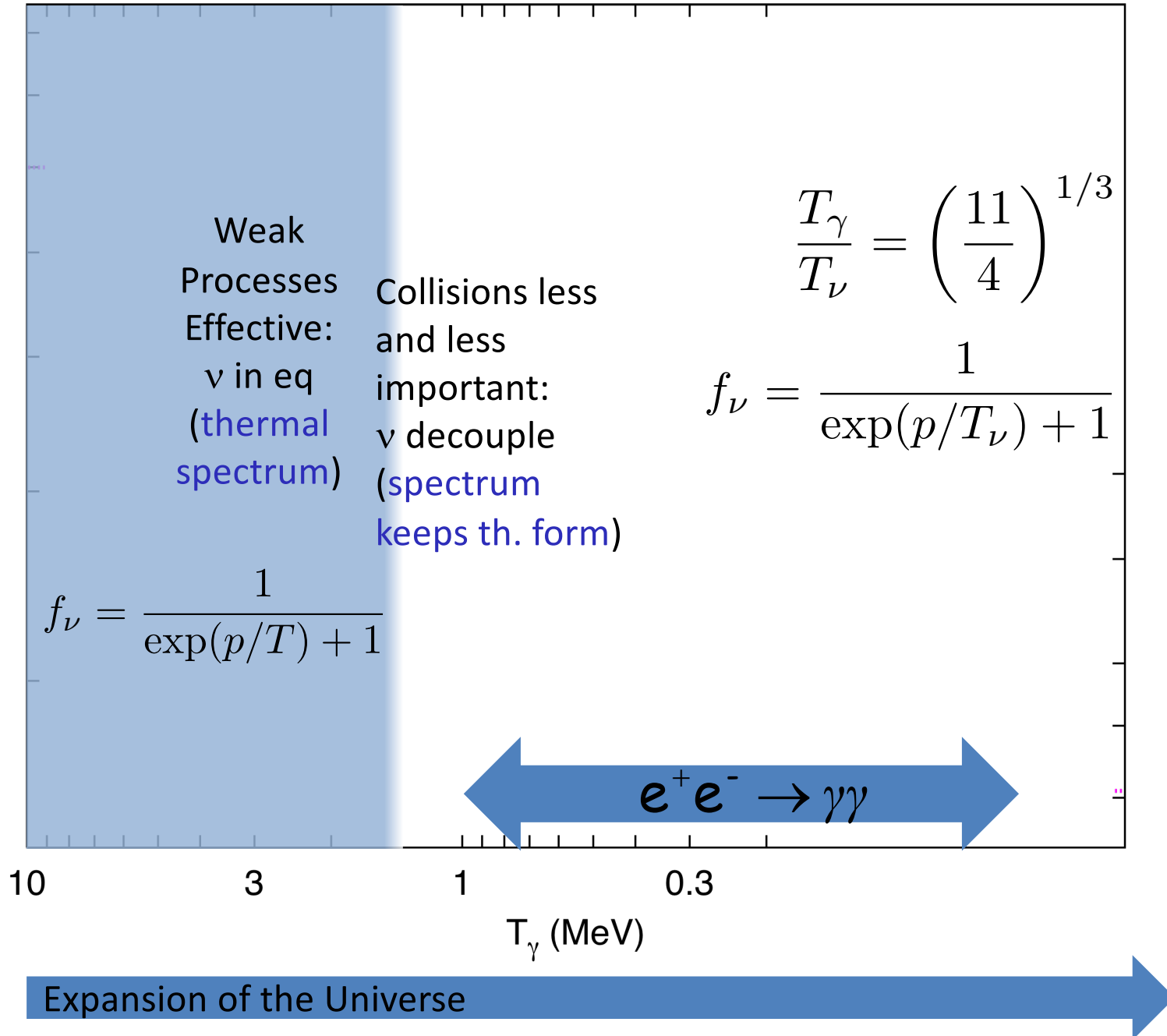
heating photons  
but not the  
decoupled  
neutrinos

$$\frac{T_\gamma}{T_\nu} = \left( \frac{11}{4} \right)^{1/3}$$



$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

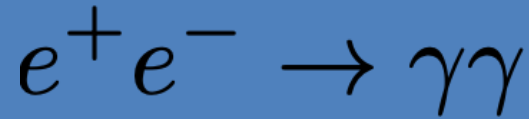
# Neutrino decoupling and $e^\pm$ annihilations





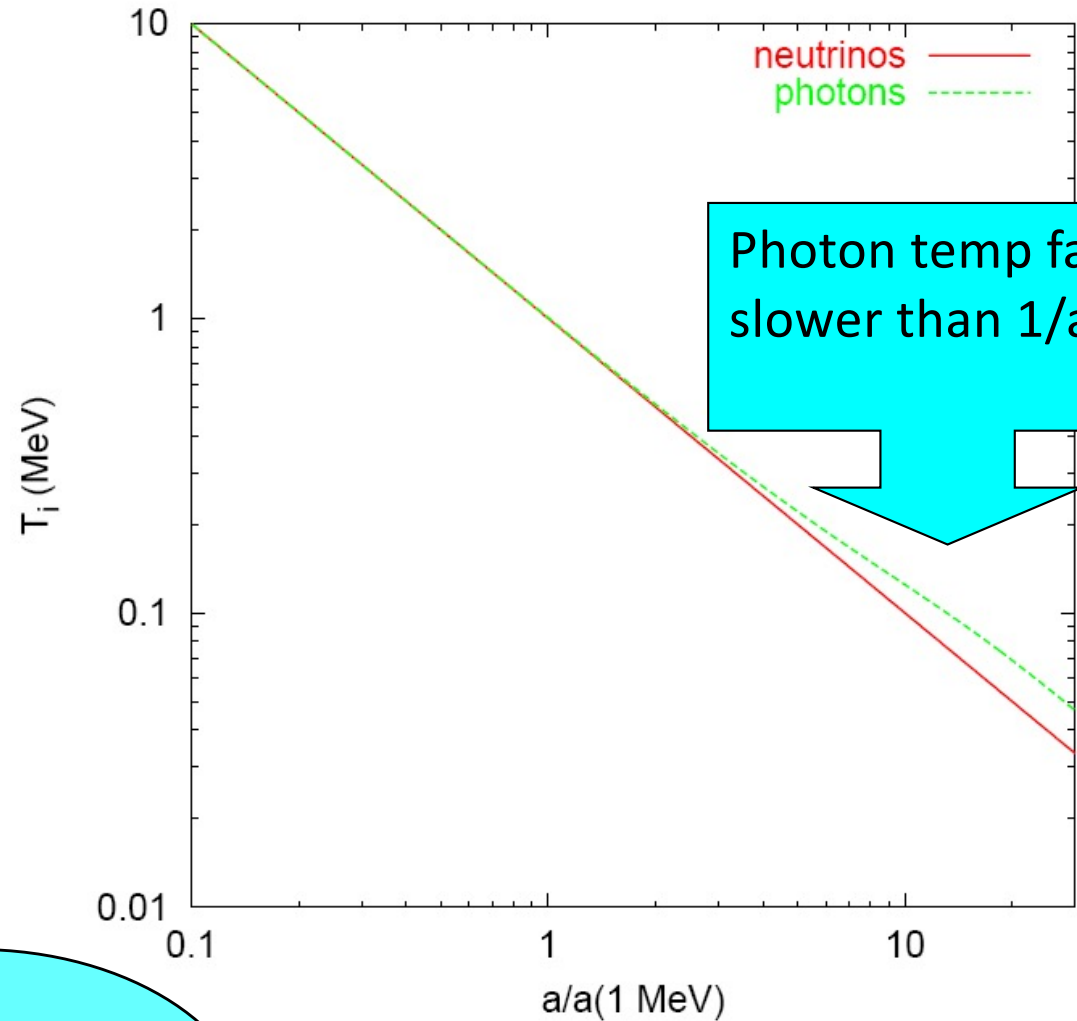
# Neutrino and Photon (CMB) temperatures

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heating photons  
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$$\frac{T_\gamma}{T_\nu} = \left( \frac{11}{4} \right)^{1/3}$$



$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

# The Cosmic Neutrino Background

Neutrinos decoupled at  $T \sim \text{MeV}$ , keeping a spectrum as that of a relativistic species

$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

- Number density

$$n_\nu = \int \frac{d^3p}{(2\pi)^3} f_\nu(p, T_\nu) = \frac{3}{11} n_\gamma = \frac{6\zeta(3)}{11\pi^2} T_{\text{CMB}}^3$$

- Energy density

$$\rho_{\nu_i} = \int \sqrt{p^2 + m_{\nu_i}^2} \frac{d^3p}{(2\pi)^3} f_\nu(p, T_\nu) \rightarrow \begin{cases} \frac{7\pi^2}{120} \left(\frac{4}{11}\right)^{4/3} T_{\text{CMB}}^4 & \text{Massless} \\ m_{\nu_i} n_\nu & \text{Massive } m_\nu \gg T \end{cases}$$

# The Cosmic Neutrino Background

Neutrinos decoupled at  $T \sim \text{MeV}$ , keeping a spectrum as that of a relativistic species

$$f_\nu(p, T) = \frac{1}{\exp(p/T_\nu) + 1}$$

- Number density

At present  $112 (\nu + \bar{\nu}) \text{ cm}^{-3}$  per flavour

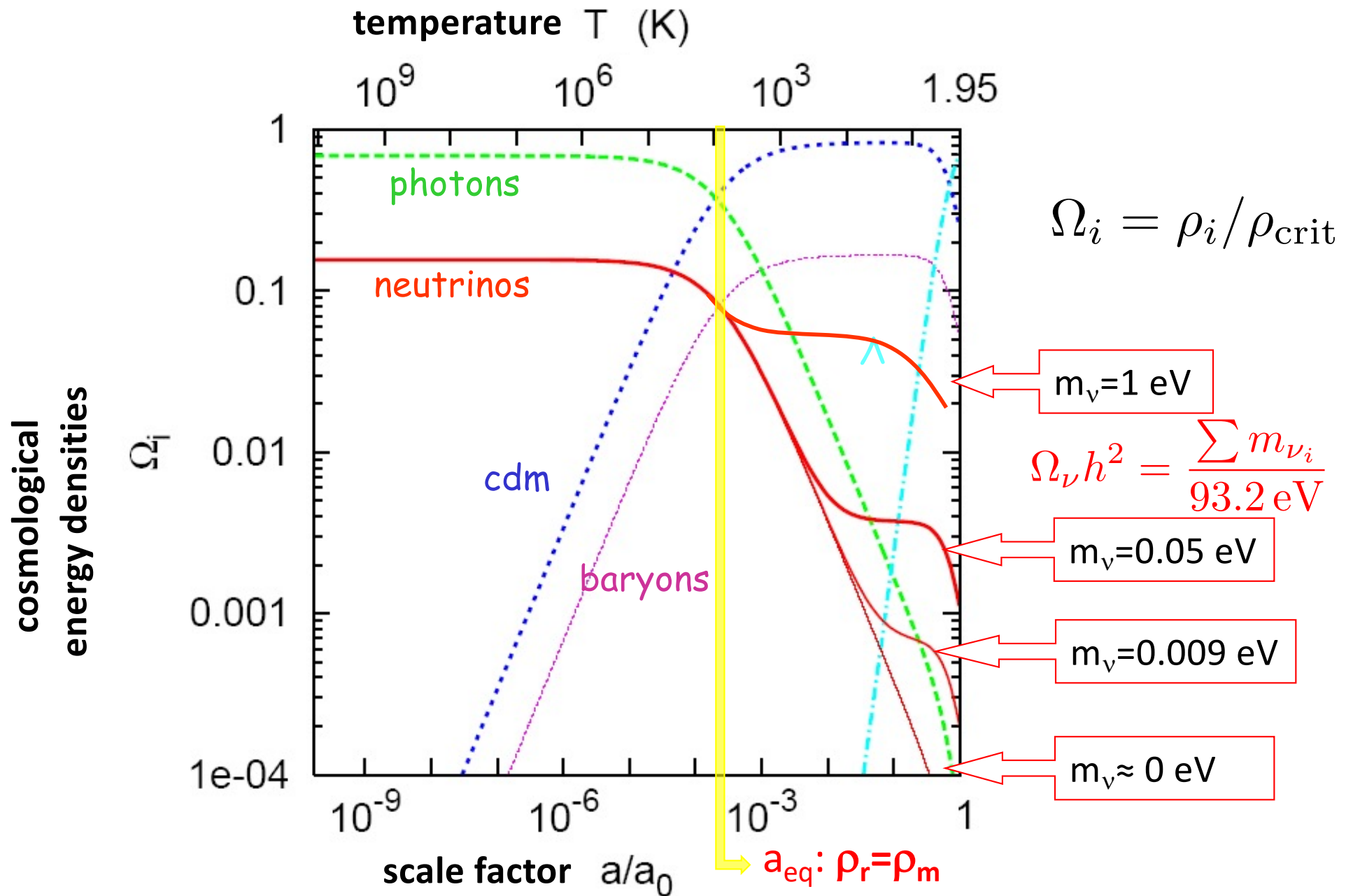
- Energy density

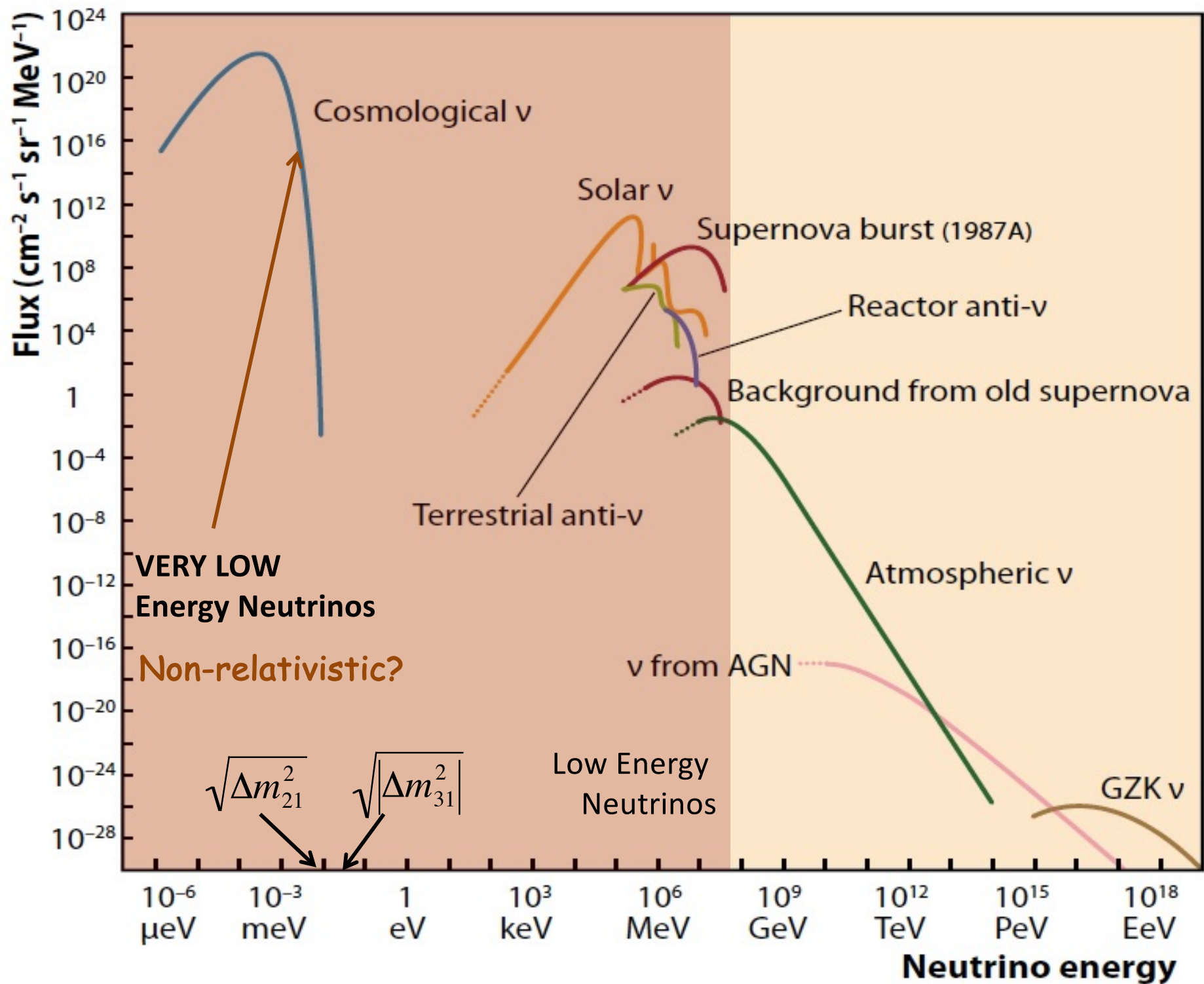
Contribution to the energy density of the Universe

$$\Omega_\nu h^2 \simeq 1.7 \times 10^{-5} \quad \text{Massless}$$

$$\Omega_\nu h^2 = \frac{\sum_i m_{\nu_i}}{94.1 \text{ eV}} \quad \text{Massive} \\ m_\nu \gg T$$

# Evolution of the background densities: 1 MeV $\rightarrow$ now





# The radiation content of the Universe ( $N_{\text{eff}}$ )

# Relativistic particles in the universe

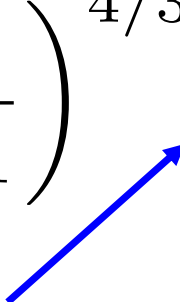
At  $T < m_e$ , the radiation content of the Universe is

$$\rho_{\text{rad}} = \rho_{\gamma} + \rho_{\nu} = \rho_{\gamma} \left[ 1 + \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} \times 3 \right]$$

Valid for standard neutrinos in the  
instantaneous decoupling approximation

# Relativistic particles in the universe

At  $T < m_e$ , the radiation content of the Universe is

$$\rho_{\text{rad}} = \rho_{\gamma} + \rho_{\nu} + \rho_x = \rho_{\gamma} \left[ 1 + \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} N_{\text{eff}} \right]$$


effective number of relativistic neutrino species  
(effective number of neutrinos)

$N_{\text{eff}}$  is a way to measure the ratio  $\frac{\rho_{\nu} + \rho_x}{\rho_{\gamma}}$

1960s-1970s :  $N_{\text{eff}} = N_{\nu}$ , **extra neutrinos** would enhance the cosmological expansion

>1980s:  $N_{\text{eff}} =$  **additional relativistic particles**

Number of light neutrino types (LEP data)  $N_{\nu} = 2.984 \pm 0.008$



# Relativistic particles in the universe

At  $T < m_e$ , the radiation content of the Universe is

$$\rho_{\text{rad}} = \rho_{\gamma} + \rho_{\nu} + \rho_x = \rho_{\gamma} \left[ 1 + \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} N_{\text{eff}} \right]$$

**effective number of neutrinos**

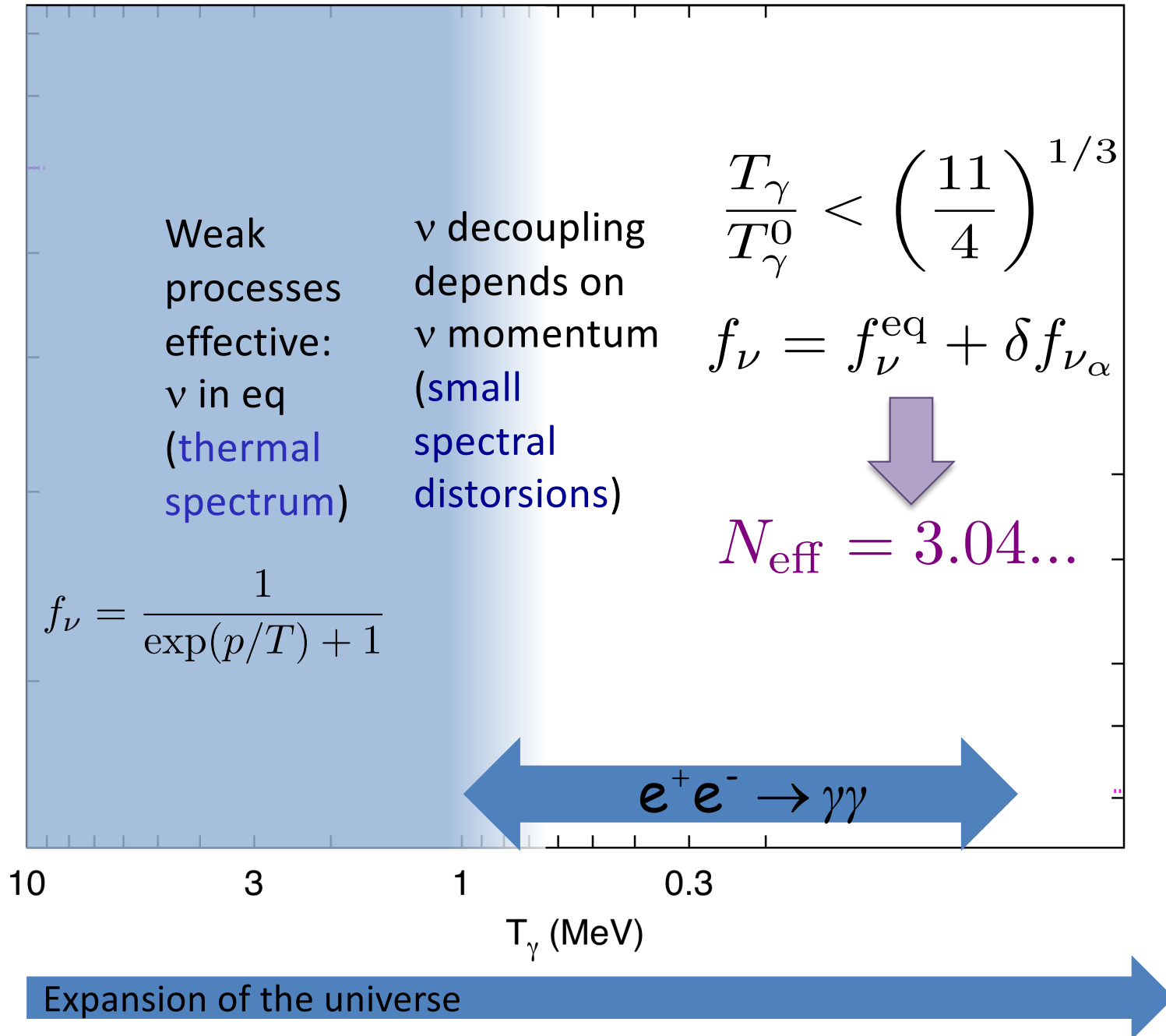
$N_{\text{eff}} \neq 3$

**additional relativistic particles** (scalars, pseudoscalars, decay products of heavy particles,...)

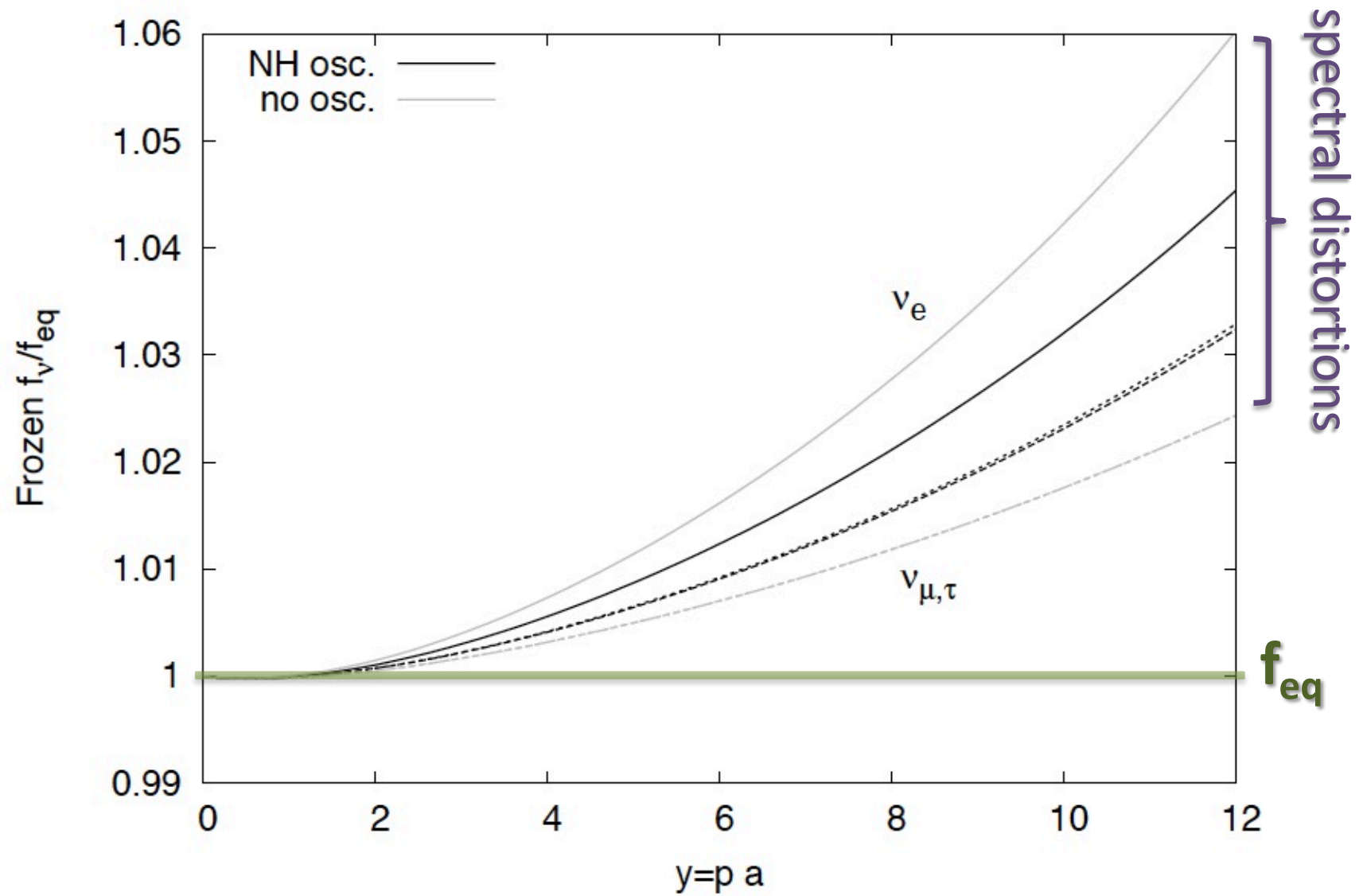
**non-standard neutrino physics** (primordial neutrino asymmetries, totally or partially thermalized light sterile neutrinos, non-standard interactions with electrons,...)

**$N_{\text{eff}} \neq 3$  in the standard case**

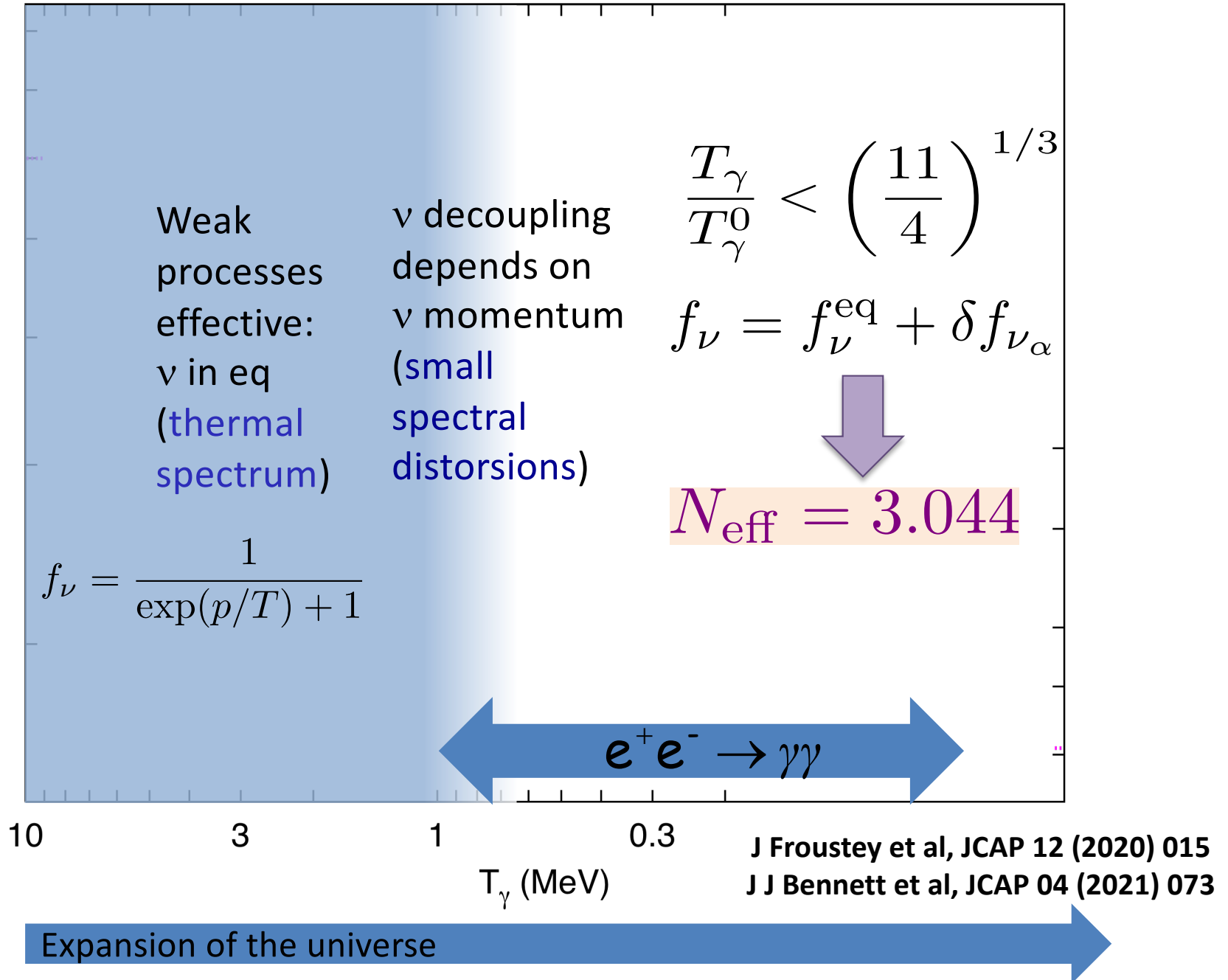
# $N_{\text{eff}} > 3$ : small neutrino heating



# Standard scenario: final distortions



# $N_{\text{eff}} > 3$ : small neutrino heating

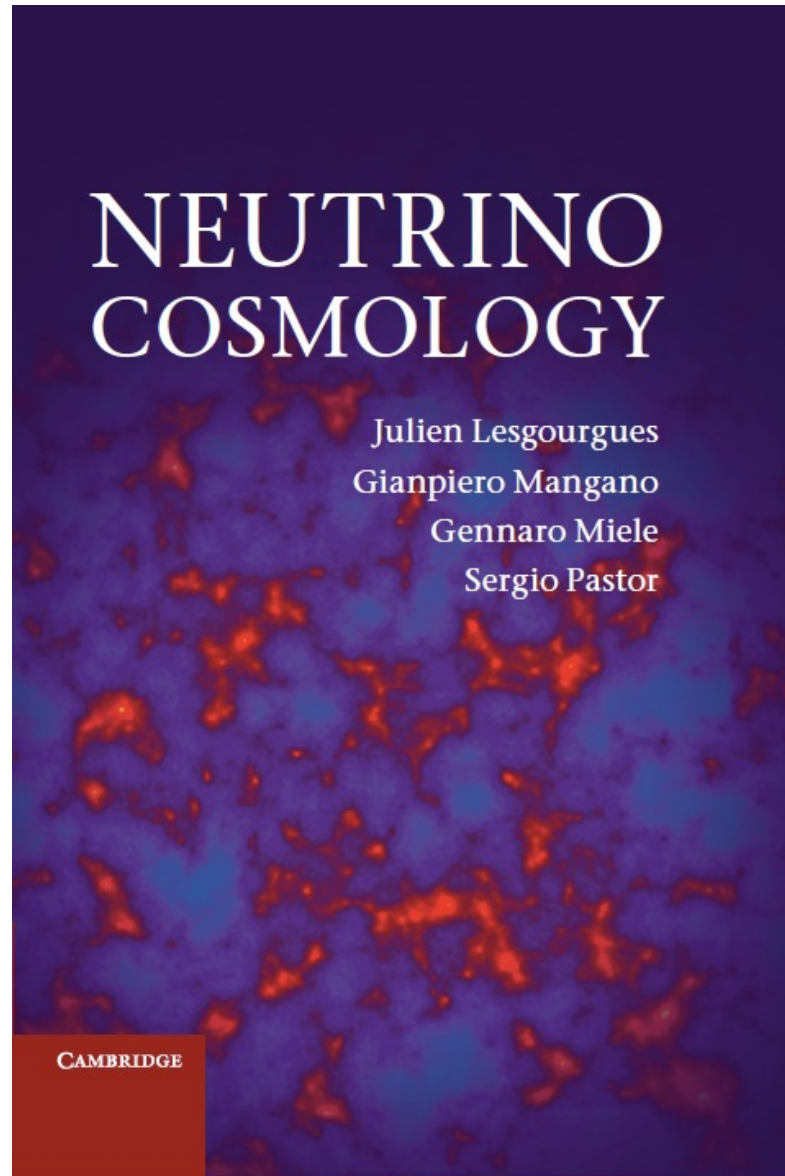


**End**

## Exercises: try to calculate...

- The present number density of massive/massless neutrinos  $n_\nu^0$  in  $\text{cm}^{-3}$
  - The present energy density of massive/massless neutrinos  $\Omega_\nu^0$  and find the limits on the total neutrino mass from  $\Omega_\nu^0 < 1$  and  $\Omega_\nu^0 < \Omega_m^0$
  - The final ratio  $T_\gamma/T_\nu$  using the conservation of entropy density before/after  $e^\pm$  annihilations
  - The decoupling temperature of relic neutrinos using  $\Gamma_w \approx H$
- 
- The photon temperature / redshift of the matter radiation equality for  $m_\nu = 1 \text{ eV}$

**For more details...**



Ed. Cambridge Univ. Press, 2013



## Suggested References

### Books

*Modern Cosmology*, S. Dodelson (Academic Press, 2003)

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*Neutrino Cosmology*, Mangano, Miele, Lesgourgues & SP (Cambridge U., 2013)

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*Ann. Rev. Nucl. Part. Sci.* 61 (2011) 69-98 [[arXiv:1111.1436](#)]

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