

Anomalies and Sterile Neutrinos – Implications of New Theoretical Results

MEDEX'19

Joel Kostensalo

University of Jyväskylä

May 30, 2019

1 Motivation

2 Reactor anomaly

- Spectrum shape
- Results

3 Gallium anomaly

- Theoretical results
- Charge-exchange reaction results
- Angular distributions
- Tensor contributions

4 Summary

Why should we revisit the reactor antineutrino and gallium anomalies?

- ▶ They have been long unexplained.
- ▶ It has been suggested that new physics such as the existence of one or more eV-scale *sterile neutrinos* could be behind these discrepancies.
- ▶ Disagreement between experiment and theory has been reported at the $2\text{--}3\sigma$ level
- ▶ The previous theoretical estimates use **very crude approximations** but are often treated as reliable.

Why should we revisit the reactor antineutrino and gallium anomalies?

- ▶ They have been long unexplained.
- ▶ It has been suggested that new physics such as the existence of one or more eV-scale *sterile neutrinos* could be behind these discrepancies.
- ▶ Disagreement between experiment and theory has been reported at the $2-3\sigma$ level
- ▶ The previous theoretical estimates use **very crude approximations** but are often treated as reliable.

Why should we revisit the reactor antineutrino and gallium anomalies?

- ▶ They have been long unexplained.
- ▶ It has been suggested that new physics such as the existence of one or more eV-scale *sterile neutrinos* could be behind these discrepancies.
- ▶ Disagreement between experiment and theory has been reported at the $2\text{--}3\sigma$ level
- ▶ The previous theoretical estimates use **very crude approximations** but are often treated as reliable.

Why should we revisit the reactor antineutrino and gallium anomalies?

- ▶ They have been long unexplained.
- ▶ It has been suggested that new physics such as the existence of one or more eV-scale *sterile neutrinos* could be behind these discrepancies.
- ▶ Disagreement between experiment and theory has been reported at the $2-3\sigma$ level
- ▶ The previous theoretical estimates use **very crude approximations** but are often treated as reliable.

Why should we revisit the reactor antineutrino and gallium anomalies?

- ▶ They have been long unexplained.
- ▶ It has been suggested that new physics such as the existence of one or more eV-scale *sterile neutrinos* could be behind these discrepancies.
- ▶ Disagreement between experiment and theory has been reported at the $2\text{--}3\sigma$ level
- ▶ The previous theoretical estimates use **very crude approximations** but are often treated as reliable.

Short-baseline reactor neutrino experiments have two problems when compared to theory:

- 1) Total number of detected antineutrinos is 6 % lower
- 2) Detected energy spectrum has a bump

Problem

Many of the contributing decays are forbidden but often treated as allowed or unique to simplify the calculations.

Solution

Calculate the shape factors without these approximations using the nuclear shell model.

Short-baseline reactor neutrino experiments have two problems when compared to theory:

- 1) Total number of detected antineutrinos is 6 % lower
- 2) Detected energy spectrum has a bump

Problem

Many of the contributing decays are forbidden but often treated as allowed or unique to simplify the calculations.

Solution

Calculate the shape factors without these approximations using the nuclear shell model.

Short-baseline reactor neutrino experiments have two problems when compared to theory:

- 1) Total number of detected antineutrinos is 6 % lower
- 2) Detected energy spectrum has a bump

Problem

Many of the contributing decays are forbidden but often treated as allowed or unique to simplify the calculations.

Solution

Calculate the shape factors without these approximations using the nuclear shell model.

Short-baseline reactor neutrino experiments have two problems when compared to theory:

- 1) Total number of detected antineutrinos is 6 % lower
- 2) Detected energy spectrum has a bump

Problem

Many of the contributing decays are forbidden but often treated as allowed or unique to simplify the calculations.

Solution

Calculate the shape factors without these approximations using the nuclear shell model.

Short-baseline reactor neutrino experiments have two problems when compared to theory:

- 1) Total number of detected antineutrinos is 6 % lower
- 2) Detected energy spectrum has a bump

Problem

Many of the contributing decays are forbidden but often treated as allowed or unique to simplify the calculations.

Solution

Calculate the shape factors without these approximations using the nuclear shell model.

The β spectrum shape is given by

$$\frac{dN}{dW} = pW(W - W_0)^2 F(Z, W) C(Z, W) K(Z, W), \quad (1)$$

where

- ▶ $pW(W - W_0)^2$ Kinematics
- ▶ $F(Z, W)$ Fermi-function (interaction of beta particle with the nucleus)
- ▶ $C(Z, W)$ Shape factor \Leftarrow Nuclear physics!
- ▶ $K(Z, W)$ Higher-order corrections

The β spectrum shape is given by

$$\frac{dN}{dW} = pW(W - W_0)^2 F(Z, W) C(Z, W) K(Z, W), \quad (1)$$

where

- ▶ $pW(W - W_0)^2$ Kinematics
- ▶ $F(Z, W)$ Fermi-function (interaction of beta particle with the nucleus)
- ▶ $C(Z, W)$ Shape factor \Leftarrow Nuclear physics!
- ▶ $K(Z, W)$ Higher-order corrections

The β spectrum shape is given by

$$\frac{dN}{dW} = pW(W - W_0)^2 F(Z, W) C(Z, W) K(Z, W), \quad (1)$$

where

- ▶ $pW(W - W_0)^2$ Kinematics
- ▶ $F(Z, W)$ Fermi-function (interaction of beta particle with the nucleus)
- ▶ $C(Z, W)$ Shape factor \Leftarrow Nuclear physics!
- ▶ $K(Z, W)$ Higher-order corrections

The β spectrum shape is given by

$$\frac{dN}{dW} = pW(W - W_0)^2 F(Z, W) C(Z, W) K(Z, W), \quad (1)$$

where

- ▶ $pW(W - W_0)^2$ Kinematics
- ▶ $F(Z, W)$ Fermi-function (interaction of beta particle with the nucleus)
- ▶ $C(Z, W)$ Shape factor \Leftarrow Nuclear physics!
- ▶ $K(Z, W)$ Higher-order corrections

The β spectrum shape is given by

$$\frac{dN}{dW} = pW(W - W_0)^2 F(Z, W) C(Z, W) K(Z, W), \quad (1)$$

where

- ▶ $pW(W - W_0)^2$ Kinematics
- ▶ $F(Z, W)$ Fermi-function (interaction of beta particle with the nucleus)
- ▶ $C(Z, W)$ Shape factor \Leftarrow Nuclear physics!
- ▶ $K(Z, W)$ Higher-order corrections

In the first-order the shape factor is simple for allowed and unique decays, complicated for non-unique decays. For non-unique decays the shape factor depends on

- ▶ The nuclear matrix elements
- ▶ The effective value of g_A
- ▶ Kinematic factors

Uncertainty in the spectral shape can be estimated by varying the ratios of the matrix elements. Based on earlier research we consider $g_A = 0.7-1.27$ and enhance the axial-charge matrix element by a 40%–100%.

In the first-order the shape factor is simple for allowed and unique decays, complicated for non-unique decays. For non-unique decays the shape factor depends on

- ▶ The nuclear matrix elements
- ▶ The effective value of g_A
- ▶ Kinematic factors

Uncertainty in the spectral shape can be estimated by varying the ratios of the matrix elements. Based on earlier research we consider $g_A = 0.7-1.27$ and enhance the axial-charge matrix element by a 40%–100%.

In the first-order the shape factor is simple for allowed and unique decays, complicated for non-unique decays. For non-unique decays the shape factor depends on

- ▶ The nuclear matrix elements
- ▶ The effective value of g_A
- ▶ Kinematic factors

Uncertainty in the spectral shape can be estimated by varying the ratios of the matrix elements. Based on earlier research we consider $g_A = 0.7-1.27$ and enhance the axial-charge matrix element by a 40%–100%.

In the first-order the shape factor is simple for allowed and unique decays, complicated for non-unique decays. For non-unique decays the shape factor depends on

- ▶ The nuclear matrix elements
- ▶ The effective value of g_A
- ▶ Kinematic factors

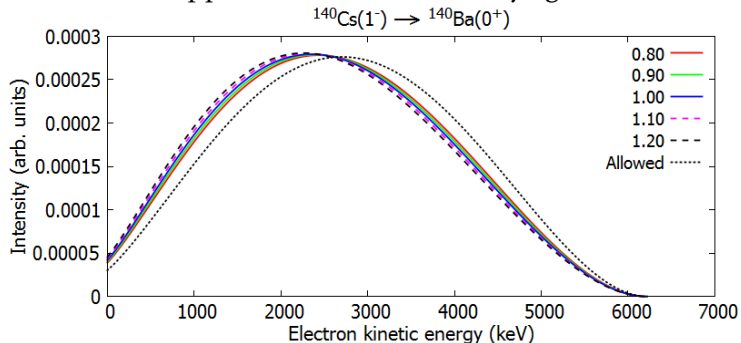
Uncertainty in the spectral shape can be estimated by varying the ratios of the matrix elements. Based on earlier research we consider $g_A = 0.7-1.27$ and enhance the axial-charge matrix element by a 40%–100%.

In the first-order the shape factor is simple for allowed and unique decays, complicated for non-unique decays. For non-unique decays the shape factor depends on

- ▶ The nuclear matrix elements
- ▶ The effective value of g_A
- ▶ Kinematic factors

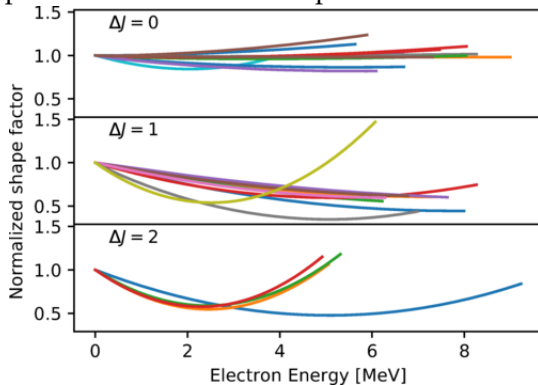
Uncertainty in the spectral shape can be estimated by varying the ratios of the matrix elements. Based on earlier research we consider $g_A = 0.7-1.27$ and enhance the axial-charge matrix element by a 40%–100%.

The allowed approximation is not always good:



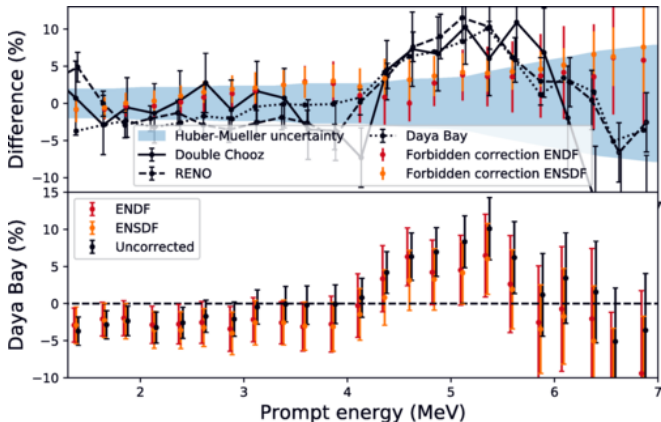
J. K. and J. Suhonen, *Int. J. Mod. Phys. A* **33**, 1843008 (2018).

The shape factors of 29 most important forbidden decays:



L. Hayen, J. K., N. Severijns, J. Suhonen, Phys Rev. C **99**, 031301(R) (2019).

The inclusion of the forbidden spectra mitigates the spectral shoulder and increases the uncertainties related to the antineutrino flux.



L. Hayen, J. K., N. Severijns, J. Suhonen, Phys Rev. C **99**, 031301(R) (2019).

The inclusion of the forbidden spectra mitigates the spectral shoulder and increases the uncertainties related to the antineutrino flux.

⇒ Decreases the statistical significance of the reactor neutrino anomaly significantly.

⇒ The forbidden distributions must be taken into account without using the allowed or unique approximations.

⇒ Precise measurement of these spectra is needed to verify/explain the reactor antineutrino anomaly

The inclusion of the forbidden spectra mitigates the spectral shoulder and increases the uncertainties related to the antineutrino flux.

⇒ Decreases the statistical significance of the reactor neutrino anomaly significantly.

⇒ The forbidden distributions must be taken into account without using the allowed or unique approximations.

⇒ Precise measurement of these spectra is needed to verify/explain the reactor antineutrino anomaly

The inclusion of the forbidden spectra mitigates the spectral shoulder and increases the uncertainties related to the antineutrino flux.

⇒ Decreases the statistical significance of the reactor neutrino anomaly significantly.

⇒ The forbidden transitions must be taken into account without using the allowed or unique approximations.

⇒ Precise measurement of these spectra is needed to verify/explain the reactor antineutrino anomaly

The inclusion of the forbidden spectra mitigates the spectral shoulder and increases the uncertainties related to the antineutrino flux.

⇒ Decreases the statistical significance of the reactor neutrino anomaly significantly.

⇒ The forbidden transitions must be taken into account without using the allowed or unique approximations.

⇒ Precise measurement of these spectra is needed to verify/explain the reactor antineutrino anomaly

The gallium anomaly refers to the missing electron-neutrino flux from ^{37}Ar and ^{51}Cr electron-capture decays as measured by the GALLEX and SAGE solar-neutrino detectors

Statistical analysis of Giunti and Laveder (2011) found a statistically significant difference between the experiments and the theoretical prediction of Bahcall at the 3.0σ level

Problem

The theoretical analysis assumes (p, n) -reaction BGTs and upper limits for BGTs are reliable estimates for weak BGTs

Solution?

Large-scale shell model calculation for the cross section.
Tensor contributions in charge-exchange reactions.

The gallium anomaly refers to the missing electron-neutrino flux from ^{37}Ar and ^{51}Cr electron-capture decays as measured by the GALLEX and SAGE solar-neutrino detectors

Statistical analysis of Giunti and Laveder (2011) found a statistically significant difference between the experiments and the theoretical prediction of Bahcall at the 3.0σ level

Problem

The theoretical analysis assumes (p, n) -reaction BGTs and upper limits for BGTs are reliable estimates for weak BGTs

Solution?

Large-scale shell model calculation for the cross section.
Tensor contributions in charge-exchange reactions.

The gallium anomaly refers to the missing electron-neutrino flux from ^{37}Ar and ^{51}Cr electron-capture decays as measured by the GALLEX and SAGE solar-neutrino detectors

Statistical analysis of Giunti and Laveder (2011) found a statistically significant difference between the experiments and the theoretical prediction of Bahcall at the 3.0σ level

Problem

The theoretical analysis assumes (p,n) -reaction BGTs and upper limits for BGTs are reliable estimates for weak BGTs

Solution?

Large-scale shell model calculation for the cross section.
Tensor contributions in charge-exchange reactions.

The gallium anomaly refers to the missing electron-neutrino flux from ^{37}Ar and ^{51}Cr electron-capture decays as measured by the GALLEX and SAGE solar-neutrino detectors

Statistical analysis of Giunti and Laveder (2011) found a statistically significant difference between the experiments and the theoretical prediction of Bahcall at the 3.0σ level

Problem

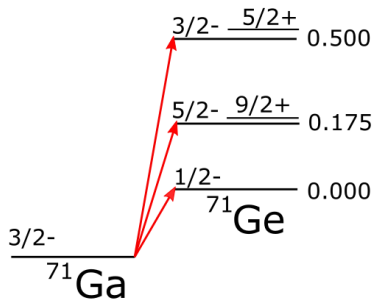
The theoretical analysis assumes (p, n) -reaction BGTs and upper limits for BGTs are reliable estimates for weak BGTs

Solution?

Large-scale shell model calculation for the cross section.
Tensor contributions in charge-exchange reactions.

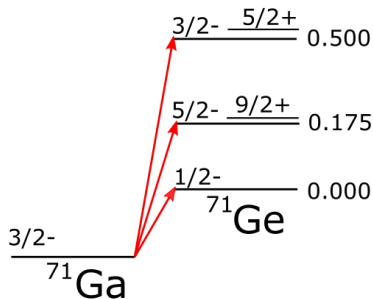
Evaluating the cross section:

- ▶ Gs-to-gs cross section can be deduced from beta decay of ^{71}Ge
- ▶ For the excited states other methods must be used (calculations, CERs)
- ▶ Bahcall used (p, n) -BGTs (more specifically half of the old upper limit <0.056 for $\text{BGT}_{5/2-}/\text{BGT}_{\text{g.s.}}$)



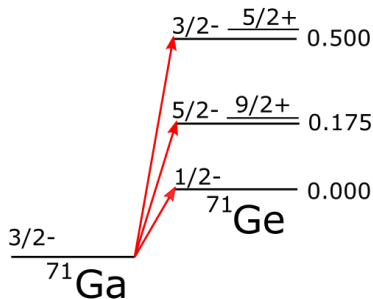
Evaluating the cross section:

- ▶ Gs-to-gs cross section can be deduced from beta decay of ^{71}Ge
- ▶ For the excited states other methods must be used (calculations, CERs)
- ▶ Bahcall used (p, n) -BGTs (more specifically half of the old upper limit <0.056 for $\text{BGT}_{5/2-}/\text{BGT}_{\text{g.s.}}$)



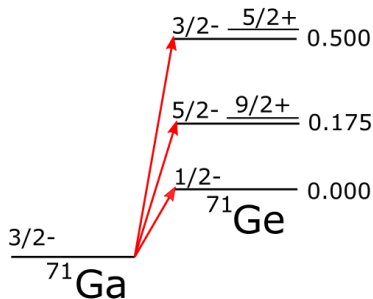
Evaluating the cross section:

- ▶ Gs-to-gs cross section can be deduced from beta decay of ^{71}Ge
- ▶ For the excited states other methods must be used (calculations, CERs)
- ▶ Bahcall used (p, n) -BGTs (more specifically half of the old upper limit <0.056 for $\text{BGT}_{5/2-}/\text{BGT}_{\text{g.s.}}$)



Evaluating the cross section:

- ▶ Gs-to-gs cross section can be deduced from beta decay of ^{71}Ge
- ▶ For the excited states other methods must be used (calculations, CERs)
- ▶ Bahcall used (p, n) -BGTs (more specifically half of the old upper limit <0.056 for $\text{BGT}_{5/2-}/\text{BGT}_{\text{g.s.}}$)



We ran new calculations using the nuclear shell model in the whole $0f_{5/2} - 1p - 0g_{9/2}$ model space using several effective Hamiltonians of which the best turned out to be JUN45

- ▶ Reproduces the excitation spectrum relatively well
- ▶ Reproduces the ^{71}Ge half-life with $g_A = 0.955$
- ▶ Agreement with experimental dipole and quadrupole moments

We ran new calculations using the nuclear shell model in the whole $0f_{5/2} - 1p - 0g_{9/2}$ model space using several effective Hamiltonians of which the best turned out to be JUN45

- ▶ Reproduces the excitation spectrum relatively well
- ▶ Reproduces the ^{71}Ge half-life with $g_A = 0.955$
- ▶ Agreement with experimental dipole and quadrupole moments

We ran new calculations using the nuclear shell model in the whole $0f_{5/2} - 1p - 0g_{9/2}$ model space using several effective Hamiltonians of which the best turned out to be JUN45

- ▶ Reproduces the excitation spectrum relatively well
- ▶ Reproduces the ^{71}Ge half-life with $g_A = 0.955$
- ▶ Agreement with experimental dipole and quadrupole moments

We ran new calculations using the nuclear shell model in the whole $0f_{5/2} - 1p - 0g_{9/2}$ model space using several effective Hamiltonians of which the best turned out to be JUN45

- ▶ Reproduces the excitation spectrum relatively well
- ▶ Reproduces the ^{71}Ge half-life with $g_A = 0.955$
- ▶ Agreement with experimental dipole and quadrupole moments

Table: Cross-section results for the ^{51}Cr neutrinos.

$1/2_{\text{g.s.}}^-$	$5.53 \pm 0.07 \times 10^{-45}$
$5/2_1^-$	$1.21 \pm 0.61 \times 10^{-46}$
$9/2_1^+$	$\leq 10^{-56}$
$3/2_1^-$	$1.94 \pm 0.97 \times 10^{-47}$
total	$5.67 \pm 0.10 \times 10^{-45}$

Table: Cross-section results for the ^{37}Ar neutrinos.

$1/2_{\text{g.s.}}^-$	$6.62 \pm 0.09 \times 10^{-45}$
$5/2_1^-$	$1.51 \pm 0.76 \times 10^{-46}$
$9/2_1^+$	$\leq 10^{-56}$
$3/2_1^-$	$2.79 \pm 1.40 \times 10^{-47}$
$5/2_1^+$	$5.91 \pm 2.96 \times 10^{-51}$
total	$6.80 \pm 0.12 \times 10^{-45}$

Shell model cross sections:

$$5.67 \pm 0.10 \times 10^{-45} \text{ cm}^2 \quad ({}^{51}\text{Cr})$$

$$6.80 \pm 0.12 \times 10^{-45} \text{ cm}^2 \quad ({}^{37}\text{Ar})$$

Bahcall cross sections:

$$5.81^{+0.21}_{-0.16} \times 10^{-45} \text{ cm}^2 \quad ({}^{51}\text{Cr})$$

$$7.00^{+0.49}_{-0.21} \times 10^{-45} \text{ cm}^2 \quad ({}^{37}\text{Ar})$$

Statistical significance of the gallium anomaly using the shell model cross sections with one-way t -test:

$$P = 0.09,$$

i.e. not statistically significant.

Problem

Charge-exchange reactions predict higher cross sections for the excited states. Why?

Statistical significance of the gallium anomaly using the shell model cross sections with one-way t -test:

$$P = 0.09,$$

i.e. not statistically significant.

Problem

Charge-exchange reactions predict higher cross sections for the excited states. Why?

Cross section can be expressed as

$$\sigma = \sigma_{\text{gs}} \left(1 + \xi_{5/2^-} \frac{\text{BGT}_{5/2^-}}{\text{BGT}_{\text{gs}}} + \xi_{3/2^-} \frac{\text{BGT}_{3/2^-}}{\text{BGT}_{\text{gs}}} \right)$$

Study	Method	$\frac{\text{BGT}_{5/2^-}}{\text{BGT}_{\text{gs}}}$	$\frac{\text{BGT}_{3/2^-}}{\text{BGT}_{\text{gs}}}$
Krofcheck et al.	(p, n)	<0.057	0.126 ± 0.023
Bahcall		0.028	0.146
Frekers et al.	$(^3\text{He}, t)$	0.039 ± 0.030	0.202 ± 0.016
Present	ISM	0.033 ± 0.017	0.016 ± 0.008

Possible problems in extracting the BGT value:

- ▶ Extraction of the $[J_{pro} J_{tar} J_{rel}] = [110]$ component at 0° . Is the nuclear structure input valid and what are the uncertainties related to this?
- ▶ Relating the $[J_{pro} J_{tar} J_{rel}] = [110]$ component at 0° to the GT strength: possible significant contributions from $L = 2$ matrix element.

Possible problems in extracting the BGT value:

- ▶ Extraction of the $[J_{pro} J_{tar} J_{rel}] = [110]$ component at 0° .
Is the nuclear structure input valid and what are the uncertainties related to this?
- ▶ Relating the $[J_{pro} J_{tar} J_{rel}] = [110]$ component at 0° to the GT strength: possible significant contributions from $L = 2$ matrix element.

Possible problems in extracting the BGT value:

- ▶ Extraction of the $[J_{pro} J_{tar} J_{rel}] = [110]$ component at 0° . Is the nuclear structure input valid and what are the uncertainties related to this?
- ▶ Relating the $[J_{pro} J_{tar} J_{rel}] = [110]$ component at 0° to the GT strength: possible significant contributions from $L = 2$ matrix element.

Angular distribution analysis

"One-body transition densities (OBTDs) were calculated in the shell-model code NuShellX using the GXPF1a interaction in the full fp -model space"

Frekers et al. (2011)

Excitation spectrum of ^{71}Ge using this Hamiltonian:

$5/2^-$ 0.000 MeV

$1/2^-$ 0.388 MeV

$3/2^-$ 1.496 MeV

Angular distribution analysis

"One-body transition densities (OBTDs) were calculated in the shell-model code NuShellX using the GXPF1a interaction in the full fp -model space"

Frekers et al. (2011)

Excitation spectrum of ^{71}Ge using this Hamiltonian:

$5/2^-$ 0.000 MeV

$1/2^-$ 0.388 MeV

$3/2^-$ 1.496 MeV

This interaction's one-body transition densities (OBTDs) give the BGT values

$$\text{BGT}_{1/2^-} = 0.390$$

$$\text{BGT}_{5/2^-} = 0.001$$

$$\text{BGT}_{3/2^-} = 0.271$$

Requires $g_A \approx 0.6$ to reproduce the experimental half life of ${}^{71}\text{Ge}$.

Possible problem

With these the ground-state transition is 92%, transition to $5/2^-$ state is 40%, and the transition to $3/2^-$ state is 87 % [110]. How accurate are these?

This interaction's one-body transition densities (OBTDs) give the BGT values

$$\text{BGT}_{1/2^-} = 0.390$$

$$\text{BGT}_{5/2^-} = 0.001$$

$$\text{BGT}_{3/2^-} = 0.271$$

Requires $g_A \approx 0.6$ to reproduce the experimental half life of ^{71}Ge .

Possible problem

With these the ground-state transition is 92%, transition to $5/2^-$ state is 40%, and the transition to $3/2^-$ state is 87 % [110]. How accurate are these?

This interaction's one-body transition densities (OBTDs) give the BGT values

$$\text{BGT}_{1/2^-} = 0.390$$

$$\text{BGT}_{5/2^-} = 0.001$$

$$\text{BGT}_{3/2^-} = 0.271$$

Requires $g_A \approx 0.6$ to reproduce the experimental half life of ^{71}Ge .

Possible problem

With these the ground-state transition is 92%, transition to $5/2^-$ state is 40%, and the transition to $3/2^-$ state is 87 % [110]. How accurate are these?

In (p, n) -reactions the interference between the Gamow-Teller (GT) and tensor (T) NMEs is described by the effective linear combination

$$\langle f \| O_{(p,n)} \| i \rangle = \langle f \| O_{GT} \| i \rangle + \delta \langle f \| O_{L=2} \| i \rangle, \quad (2)$$

where i (f) is the initial (final) nuclear state and $\delta \approx 0.1$ is the mixing parameter.

The interference can be constructive or destructive.

In (p, n) -reactions the interference between the Gamow-Teller (GT) and tensor (T) NMEs is described by the effective linear combination

$$\langle f \| O_{(p,n)} \| i \rangle = \langle f \| O_{GT} \| i \rangle + \delta \langle f \| O_{L=2} \| i \rangle, \quad (2)$$

where i (f) is the initial (final) nuclear state and $\delta \approx 0.1$ is the mixing parameter.

The interference can be constructive or destructive.

In (p, n) -reactions the interference between the Gamow-Teller (GT) and tensor (T) NMEs is described by the effective linear combination

$$\langle f \| O_{(p,n)} \| i \rangle = \langle f \| O_{GT} \| i \rangle + \delta \langle f \| O_{L=2} \| i \rangle, \quad (2)$$

where i (f) is the initial (final) nuclear state and $\delta \approx 0.1$ is the mixing parameter.

The interference can be constructive or destructive.

Table: Results for ^{71}Ga with $\delta = 0.097$.

State	$\langle f \ O_{\text{GT}} \ i \rangle$	$\langle f \ O_{L=2} \ i \rangle$	$\text{BGT}_{\beta}^{\text{SM}}$	$\text{BGT}_{(p,n)}^{\text{SM}}$
$1/2_{\text{g.s.}}^-$	-0.795	0.465	0.158	0.141
$5/2_1^-$	0.144	-1.902	0.0052	0.0004
$3/2_1^-$	0.100	0.0482	0.0025	0.0027

- ▶ There is a known large destructive interference for the $5/2^-$ state (Haxton 1998)
- ▶ New calculations show that there is a smaller destructive interference for the ground state
- ▶ There is a constructive interference for the $3/2^-$ state
- ▶ The ratio $\text{BGT}_{3/2^-} / \text{BGT}_{\text{gs}}$ is over estimated in CER

Table: Results for ^{71}Ga with $\delta = 0.097$.

State	$\langle f \ O_{\text{GT}} \ i \rangle$	$\langle f \ O_{L=2} \ i \rangle$	$\text{BGT}_{\beta}^{\text{SM}}$	$\text{BGT}_{(p,n)}^{\text{SM}}$
$1/2_{\text{g.s.}}^{-}$	-0.795	0.465	0.158	0.141
$5/2_1^{-}$	0.144	-1.902	0.0052	0.0004
$3/2_1^{-}$	0.100	0.0482	0.0025	0.0027

- ▶ There is a known large destructive interference for the $5/2^{-}$ state (Haxton 1998)
- ▶ New calculations show that there is a smaller destructive interference for the ground state
- ▶ There is a constructive interference for the $3/2^{-}$ state
- ▶ The ratio $\text{BGT}_{3/2^{-}}/\text{BGT}_{\text{gs}}$ is over estimated in CER

Table: Results for ^{71}Ga with $\delta = 0.097$.

State	$\langle f \ O_{\text{GT}} \ i \rangle$	$\langle f \ O_{L=2} \ i \rangle$	$\text{BGT}_{\beta}^{\text{SM}}$	$\text{BGT}_{(p,n)}^{\text{SM}}$
$1/2_{\text{g.s.}}^{-}$	-0.795	0.465	0.158	0.141
$5/2_1^{-}$	0.144	-1.902	0.0052	0.0004
$3/2_1^{-}$	0.100	0.0482	0.0025	0.0027

- ▶ There is a known large destructive interference for the $5/2^{-}$ state (Haxton 1998)
- ▶ New calculations show that there is a smaller destructive interference for the ground state
- ▶ There is a constructive interference for the $3/2^{-}$ state
- ▶ The ratio $\text{BGT}_{3/2^{-}}/\text{BGT}_{\text{gs}}$ is over estimated in CER

Table: Results for ^{71}Ga with $\delta = 0.097$.

State	$\langle f \ O_{\text{GT}} \ i \rangle$	$\langle f \ O_{L=2} \ i \rangle$	$\text{BGT}_{\beta}^{\text{SM}}$	$\text{BGT}_{(p,n)}^{\text{SM}}$
$1/2_{\text{g.s.}}^{-}$	-0.795	0.465	0.158	0.141
$5/2_1^{-}$	0.144	-1.902	0.0052	0.0004
$3/2_1^{-}$	0.100	0.0482	0.0025	0.0027

- ▶ There is a known large destructive interference for the $5/2^{-}$ state (Haxton 1998)
- ▶ New calculations show that there is a smaller destructive interference for the ground state
- ▶ There is a constructive interference for the $3/2^{-}$ state
- ▶ The ratio $\text{BGT}_{3/2^{-}}/\text{BGT}_{\text{gs}}$ is over estimated in CER

Table: Results for ^{71}Ga with $\delta = 0.097$.

State	$\langle f \ O_{\text{GT}} \ i \rangle$	$\langle f \ O_{L=2} \ i \rangle$	$\text{BGT}_{\beta}^{\text{SM}}$	$\text{BGT}_{(p,n)}^{\text{SM}}$
$1/2_{\text{g.s.}}^-$	-0.795	0.465	0.158	0.141
$5/2_1^-$	0.144	-1.902	0.0052	0.0004
$3/2_1^-$	0.100	0.0482	0.0025	0.0027

- ▶ There is a known large destructive interference for the $5/2^-$ state (Haxton 1998)
- ▶ New calculations show that there is a smaller destructive interference for the ground state
- ▶ There is a constructive interference for the $3/2^-$ state
- ▶ The ratio $\text{BGT}_{3/2^-}/\text{BGT}_{\text{gs}}$ is over estimated in CER

$L = 2$ matrix element

Suggests that the ratio $BGT_{3/2^-}/BGT_{gs}$ is over estimated by at least 30 % in CERs.

Problem

There is still a factor 8 difference between the shell model results and the CER. Is the problem in theory or experiment?

Solution 1

Uncertainties in the mixing parameter: $L = 2$ contribution is actually larger. Also the $[1\ 1\ 0]$ component at 0° might be smaller/larger for one of these transitions. No new particles.

Solution 1

The shell model wave functions are not accurate. New physics.

$L = 2$ matrix element

Suggests that the ratio $BGT_{3/2^-}/BGT_{gs}$ is over estimated by at least 30 % in CERs.

Problem

There is still a factor 8 difference between the shell model results and the CER. Is the problem in theory or experiment?

Solution 1

Uncertainties in the mixing parameter: $L = 2$ contribution is actually larger. Also the $[1\ 1\ 0]$ component at 0° might be smaller/larger for one of these transitions. No new particles.

Solution 1

The shell model wave functions are not accurate. New physics.

$L = 2$ matrix element

Suggests that the ratio $BGT_{3/2^-}/BGT_{gs}$ is over estimated by at least 30 % in CERs.

Problem

There is still a factor 8 difference between the shell model results and the CER. Is the problem in theory or experiment?

Solution 1

Uncertainties in the mixing parameter: $L = 2$ contribution is actually larger. Also the $[1\ 1\ 0]$ component at 0° might be smaller/larger for one of these transitions. No new particles.

Solution 1

The shell model wave functions are not accurate. New physics.

$L = 2$ matrix element

Suggests that the ratio $BGT_{3/2^-}/BGT_{gs}$ is over estimated by at least 30 % in CERs.

Problem

There is still a factor 8 difference between the shell model results and the CER. Is the problem in theory or experiment?

Solution 1

Uncertainties in the mixing parameter: $L = 2$ contribution is actually larger. Also the $[1\ 1\ 0]$ component at 0° might be smaller/larger for one of these transitions. No new particles.

Solution 1

The shell model wave functions are not accurate. New physics.

- ▶ **New calculations for the reactor neutrino anomaly:**
 - 1) Decrease the difference between experiment and theory
 - 2) Mitigate the spectral shoulder

- ▶ **New calculations for the gallium anomaly:**
 - 1) Decrease the difference between experimental and theoretical neutrino-nucleus scattering cross-sections
 - 2) Explain partially the difference between the charge-exchange BGTs and the GALLEX/SAGE results.

- ▶ **Conclusion:** forbidden spectral shapes and tensor contributions in CERs must be taken into account in order to make strong claims regarding the reactor and gallium anomalies.

- ▶ **New calculations for the reactor neutrino anomaly:**
 - 1) Decrease the difference between experiment and theory
 - 2) Mitigate the spectral shoulder

- ▶ **New calculations for the gallium anomaly:**
 - 1) Decrease the difference between experimental and theoretical neutrino-nucleus scattering cross-sections
 - 2) Explain partially the difference between the charge-exchange BGTs and the GALLEX/SAGE results.

- ▶ **Conclusion:** forbidden spectral shapes and tensor contributions in CERs must be taken into account in order to make strong claims regarding the reactor and gallium anomalies.

- ▶ New calculations for the reactor neutrino anomaly:
 - 1) Decrease the difference between experiment and theory
 - 2) Mitigate the spectral shoulder

- ▶ New calculations for the gallium anomaly:
 - 1) Decrease the difference between experimental and theoretical neutrino-nucleus scattering cross-sections
 - 2) Explain partially the difference between the charge-exchange BGTs and the GALLEX/SAGE results.

- ▶ Conclusion: forbidden spectral shapes and tensor contributions in CERs must be taken into account in order to make strong claims regarding the reactor and gallium anomalies.

- ▶ New calculations for the reactor neutrino anomaly:
 - 1) Decrease the difference between experiment and theory
 - 2) Mitigate the spectral shoulder

- ▶ New calculations for the gallium anomaly:
 - 1) Decrease the difference between experimental and theoretical neutrino-nucleus scattering cross-sections
 - 2) Explain partially the difference between the charge-exchange BGTs and the GALLEX/SAGE results.

- ▶ Conclusion: forbidden spectral shapes and tensor contributions in CERs must be taken into account in order to make strong claims regarding the reactor and gallium anomalies.

- ▶ New calculations for the reactor neutrino anomaly:
 - 1) Decrease the difference between experiment and theory
 - 2) Mitigate the spectral shoulder

- ▶ New calculations for the gallium anomaly:
 - 1) Decrease the difference between experimental and theoretical neutrino-nucleus scattering cross-sections
 - 2) Explain partially the difference between the charge-exchange BGTs and the GALLEX/SAGE results.

- ▶ Conclusion: forbidden spectral shapes and tensor contributions in CERs must be taken into account in order to make strong claims regarding the reactor and gallium anomalies.

- ▶ New calculations for the reactor neutrino anomaly:
 - 1) Decrease the difference between experiment and theory
 - 2) Mitigate the spectral shoulder

- ▶ New calculations for the gallium anomaly:
 - 1) Decrease the difference between experimental and theoretical neutrino-nucleus scattering cross-sections
 - 2) Explain partially the difference between the charge-exchange BGTs and the GALLEX/SAGE results.

- ▶ Conclusion: forbidden spectral shapes and tensor contributions in CERs must be taken into account in order to make strong claims regarding the reactor and gallium anomalies.

- ▶ New calculations for the reactor neutrino anomaly:
 - 1) Decrease the difference between experiment and theory
 - 2) Mitigate the spectral shoulder

- ▶ New calculations for the gallium anomaly:
 - 1) Decrease the difference between experimental and theoretical neutrino-nucleus scattering cross-sections
 - 2) Explain partially the difference between the charge-exchange BGTs and the GALLEX/SAGE results.

- ▶ Conclusion: forbidden spectral shapes and tensor contributions in CERs must be taken into account in order to make strong claims regarding the reactor and gallium anomalies.

 L. Hayen, J. K., N. Severijns, J. Suhonen, Phys. Rev. C **99**, 031301(R) (2019).

 Carlo Giunti and Marco Laveder Phys. Rev. C **83**, 065504 (2011).

 John N. Bahcall Phys. Rev. C **56**, 3391 (1997).

 W. Haxton, PLB 431 (1998).

 Frekers et al. PLB 706 134–138 (2011).

Thank you!