

First-forbidden transitions in the reactor antineutrino anomaly

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Anomaly Introduction

What's it about in 3 steps:

Where is the anomaly?

Antineutrino's from β^- decay of reactor fission fragments

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When new physics lurks, look out for quirks!

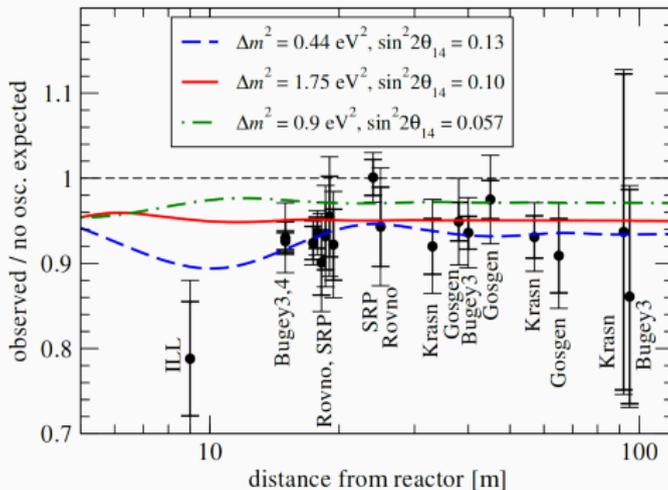
Deficiency and particle physics proposal

Current deficiency in neutrino count rate at 94% (2-3 σ)

$$P_{SBL}(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\alpha) \simeq 1$$

$$- \sin^2 2\theta_{\alpha 4} \sin^2 \left(\frac{\Delta m_{41}^2 L}{4E} \right)$$

Very exciting,
but... it is real?



Deficiency and particle physics proposal

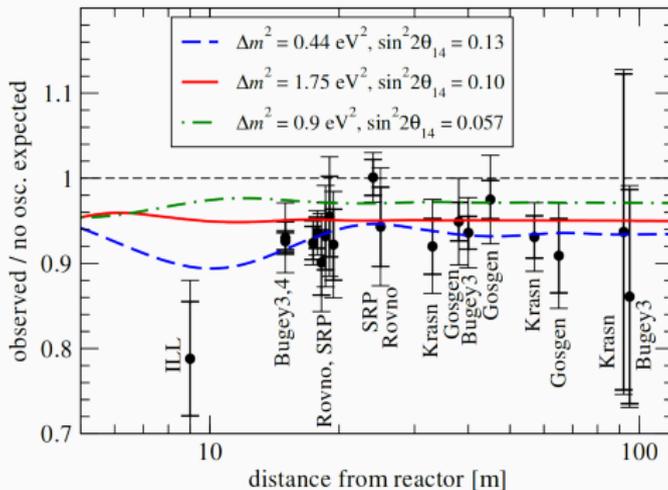
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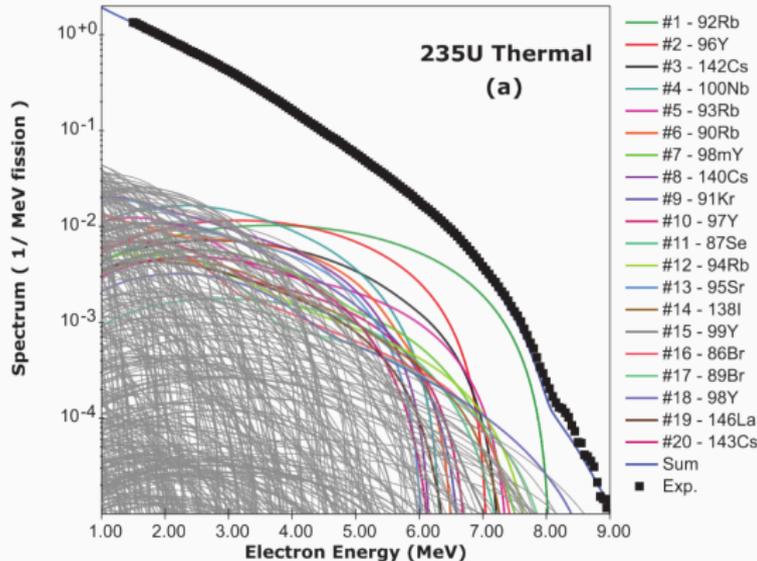
Understanding of
all corrections & nuclear
structure is **crucial!**



An *et al.* (Daya Bay Collab.), PRL 118 (2017) 251801 & J. Kopp *et al.*, JHEP 05 (2013) 050

Antineutrino origin

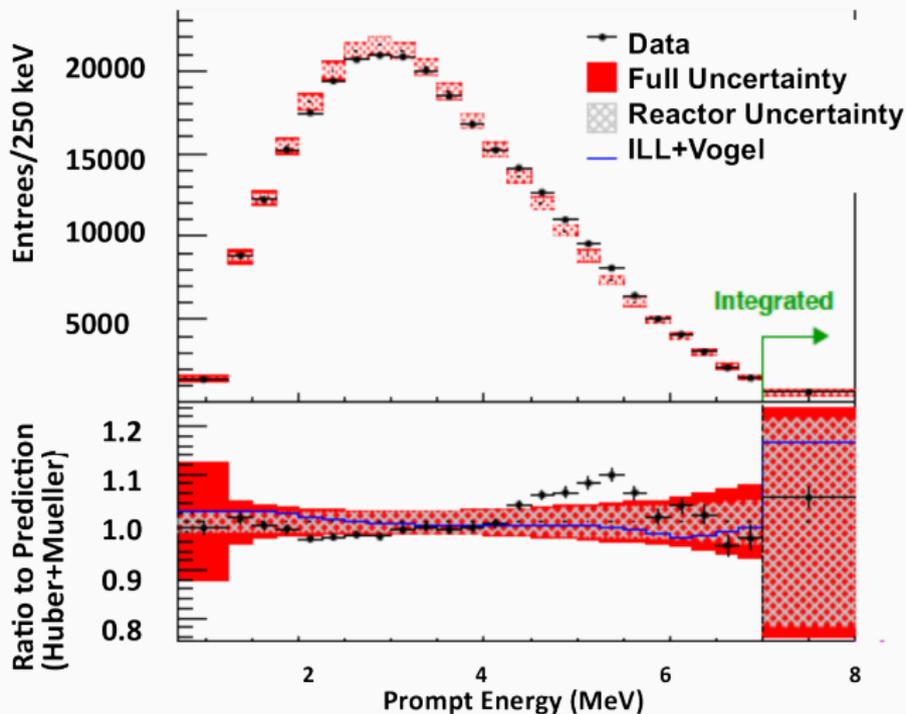
Fission fragments from ^{235}U , ^{238}U , ^{239}Pu and ^{241}Pu have many β^- branches, but can only measure **cumulative** spectrum.



Conversion of all β branches is **tremendous** challenge

A. A. Sonzogni *et al.*, PRC **91** (2015) 011301(R)

Reactor bump



Something not understood, most likely **nuclear physics** problem

Experimental status

Very short baseline experiments

Since 2011, ~ 10 experiments started setting up

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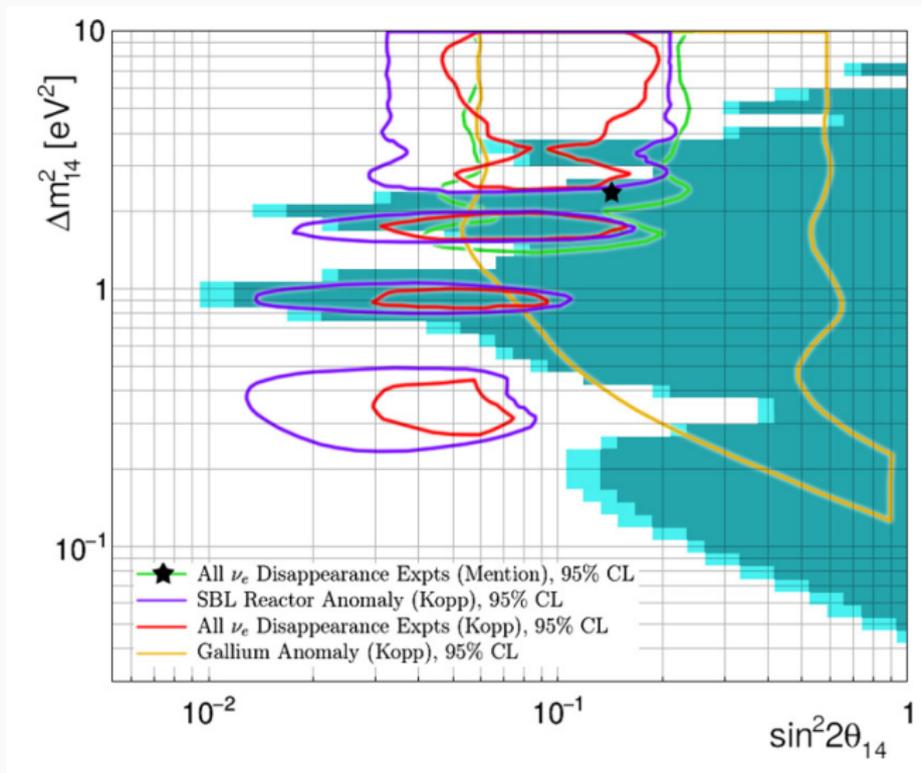
Several experiments came online late 2017/2018! Published data from

- DANSS (Russia) 1804.04046
- STEREO (France) 1806.02096
- PROSPECT (USA) 1806.02784
- NEOS (Korea) 1610.05134

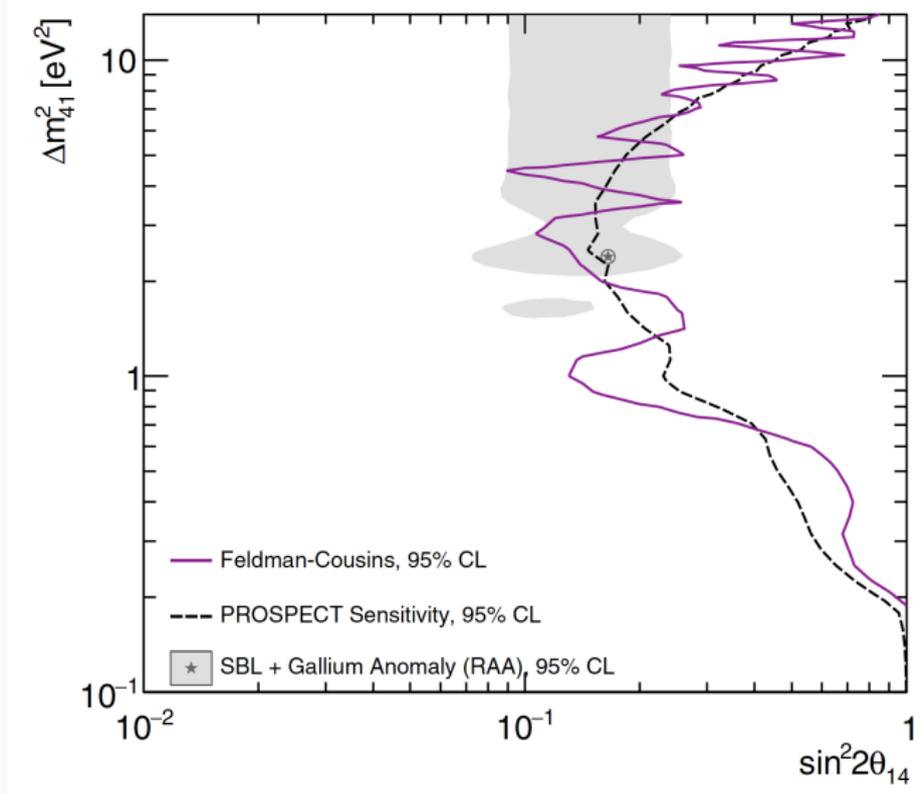
Very exciting & more coming soon!



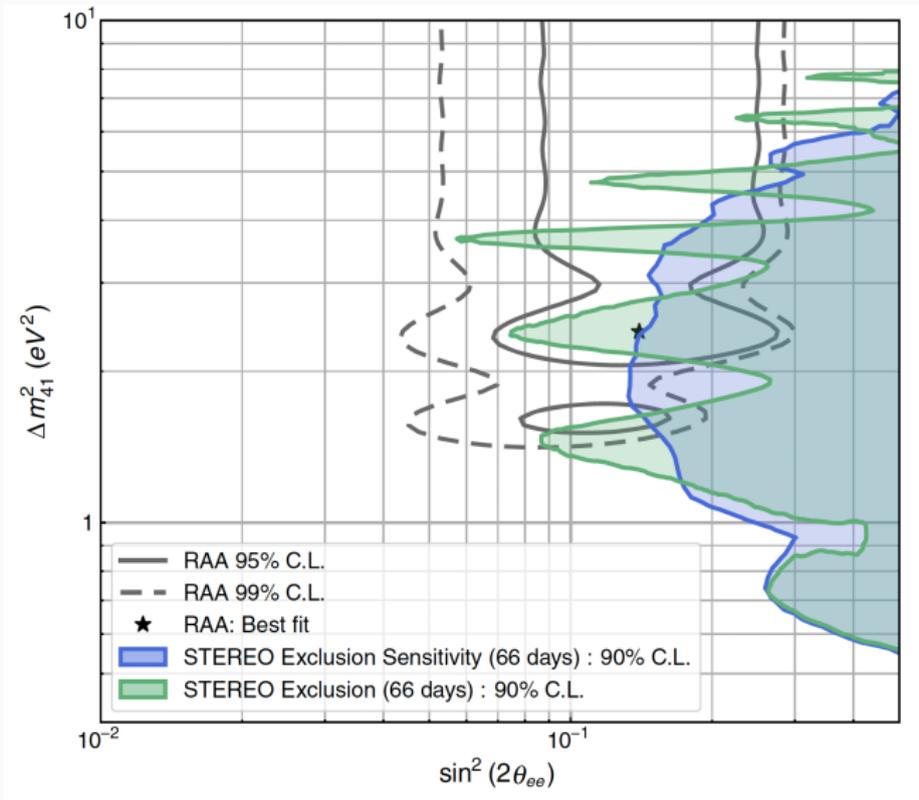
VSBL Results: DANSS



VSBL Results: PROSPECT



VSBL Results: STEREO



Faced with some interesting developments:

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Faced with some interesting developments:

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Things point to deficiencies in databases & **theoretical modeling**

Theory status

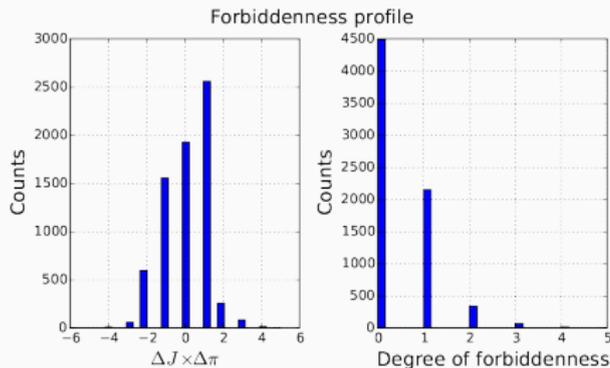
Theory: β participant sketch

Experiment sees nothing, what happens to theory?

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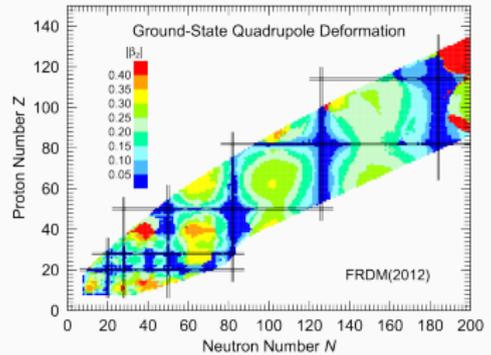
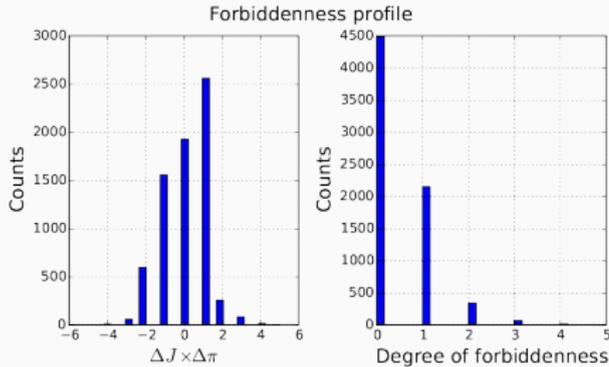
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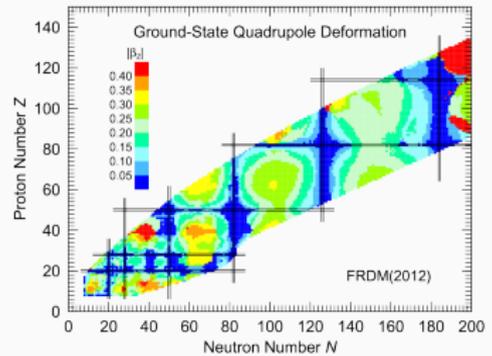
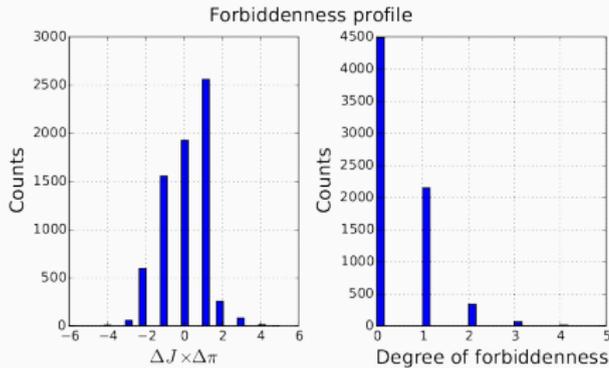


Both greatly influence the spectrum shape!

Theory: β participant sketch

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Additional lower order effects: Atomic, electrostatic, kinematic. . .

Analysis procedure

Experimental benchmark are ILL (Schreckenbach) cumulative electron spectra

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Approaches split up in 2:

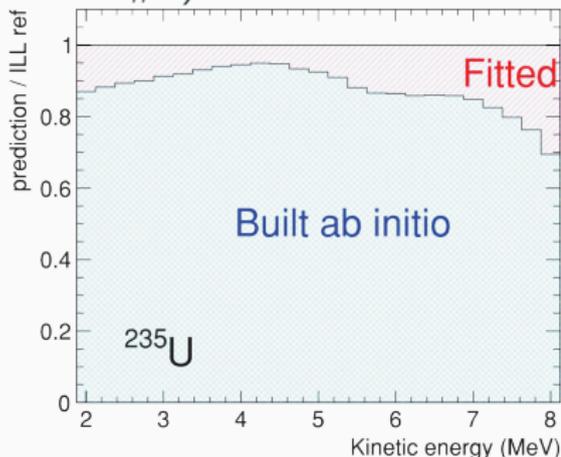
1. **Conversion** method: virtual β branch fits

Analysis procedure

Experimental benchmark are ILL (Schreckenbach) cumulative electron spectra

Approaches split up in 2:

1. **Conversion** method: virtual β branch fits
2. **Summation** method: Build from databases (& extrapolate a la #1)



Much of *summation* is based on same spectral assumptions Huber, PRC **84** (2011) 024617; Mueller *et al.*, PRC **83** (2011) 054615

Thoughts on state of the art

2 elements which require pause

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1. Central problem when comparing to ILL data

*Everything below 1.8 MeV in electron spectrum is unconstrained,
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*Everything below 1.8 MeV in electron spectrum is unconstrained,
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Everything that changes the shape below 1.8 MeV changes the anomaly → essential to get this right

Thoughts on state of the art

2 elements which require pause

2. Depending on method, questionable approximations

- Incorrectly estimates $(\alpha Z)^{n>1}$ effects, $\text{RAA}(\langle Z \rangle^{n>1}) \neq \langle \text{RAA}(Z^{N>1}) \rangle!$
- Estimated average b/Ac from spherical mirrors, but highly transition and deformation dependent
- All transitions assumed allowed/unique
- No Coulomb corrections to unique shape factors
- ...

An *et al.* (Daya Bay Collab.), PRL 118 (2017) 251801 & Hayes *et al.*,
arXiv:1707.07728

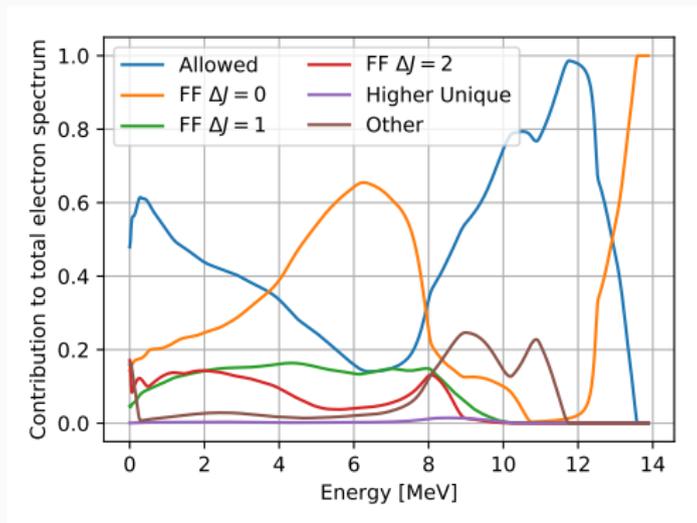
Forbidden decays

Forbidden shape factors

Roughly $\sim 30\%$ of 8000 transitions are forbidden, usually assumed of negligible importance for anomaly

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Experimental region of interest (2-8 MeV) is **dominated** by forbidden decays LH, J. Kostensalo, N. Severijns, J. Suhonen, PRC 99 (2019) 031301(R)

β spectrum shape

Central element in analysis is knowledge of β spectrum shape

$$\frac{dN}{dW} \propto pW(W_0 - W)^2 F(Z, W) C(Z, W) \dots$$

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- Treat as allowed
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Approximations in state-of-the-art for non-unique forbidden transitions

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are . . . not great

Shape factor

General shape factor

$$C(Z, W) = \sum_{k_e, k_\nu, K} \lambda_{k_e} \left\{ M_K^2(k_e, k_\nu) + m_K^2(k_e, k_\nu) - \frac{2\mu_{k_e} \gamma_{k_e}}{k_e W} M_K(k_e, k_\nu) m_K(k_e, k_\nu) \right\},$$

where

$$\lambda_{k_e} = \frac{\alpha_{-k_e}^2 + \alpha_{+k_e}^2}{\alpha_{-1}^2 + \alpha_{+1}^2},$$
$$\mu_{k_e} = \frac{\alpha_{-k_e}^2 - \alpha_{+k_e}^2}{\alpha_{-k_e}^2 + \alpha_{+k_e}^2} \frac{k_e W}{\gamma_{k_e}},$$

are Coulomb functions of $\mathcal{O}(1)$

Behrens, Bühring, Electron radial wave functions, 1982

First-forbidden transitions

Depending on spin-parity change, C can be simple ($R \sim 0.01$)

$$C_{0-} \propto 1 + \frac{2R}{3W} b + \mathcal{O}(\alpha ZR, W_0 R^2)$$

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$$C_{1-} \propto 1 + aW + \mu_1\gamma_1 \frac{b}{W} + cW^2$$

or rather simple, again

$$C_U \propto \sum_{k=1}^L \lambda_k \frac{p^{2(k-1)} q^{2(L-k)}}{(2k-1)! [2(L-k)+1]!}$$

First-forbidden transitions

There are several **complicating factors**, however

- Coulomb corrections at all levels: Fermi function, higher κ_e corrections, modified radial behaviour

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Challenging, but attempt to **establish uncertainty**

First-forbidden transitions

Cause for despair, but there's a helping hand:

First-forbidden transitions

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Higher in E you go, fewer branches contribute

First-forbidden transitions

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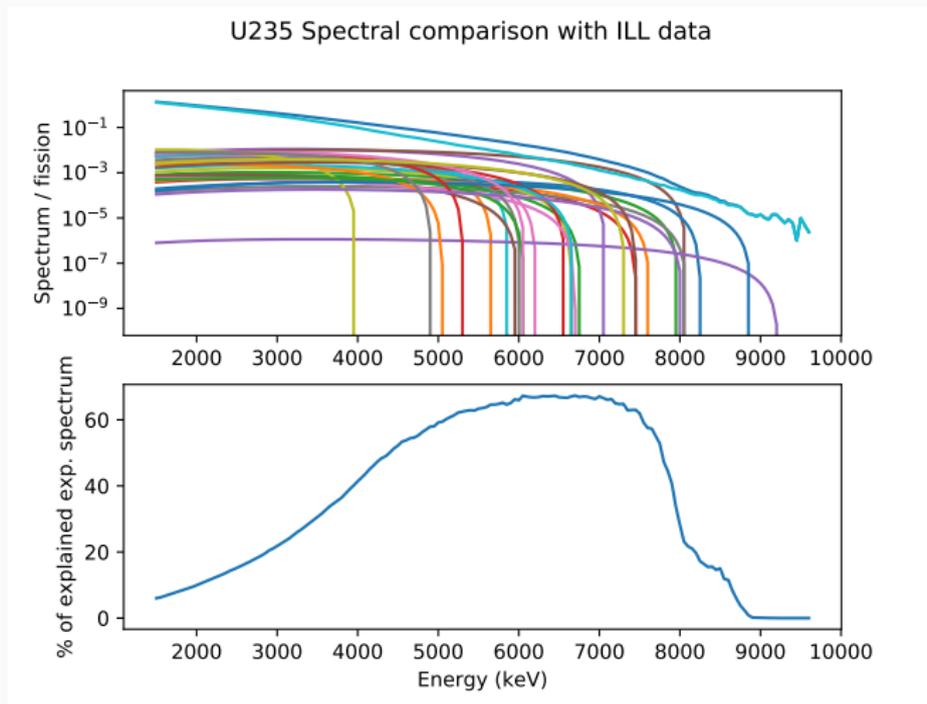
Higher in E you go, fewer branches contribute

From 5 MeV onwards: \gtrsim 90% of flux with less than 50 branches

Sonzogni *et al.*, 91 (2015) 011301

Forbidden shape factors

Picked 29 dominant forbidden transitions

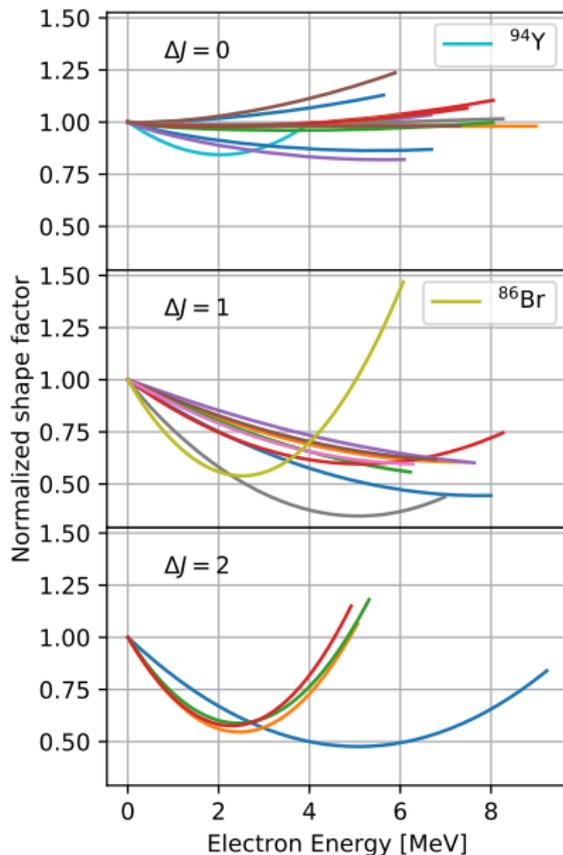


> 50% in region of interest (4-7 MeV)

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$$\frac{dN}{dE} \propto pE(E_0 - E)^2 F(Z, E)$$

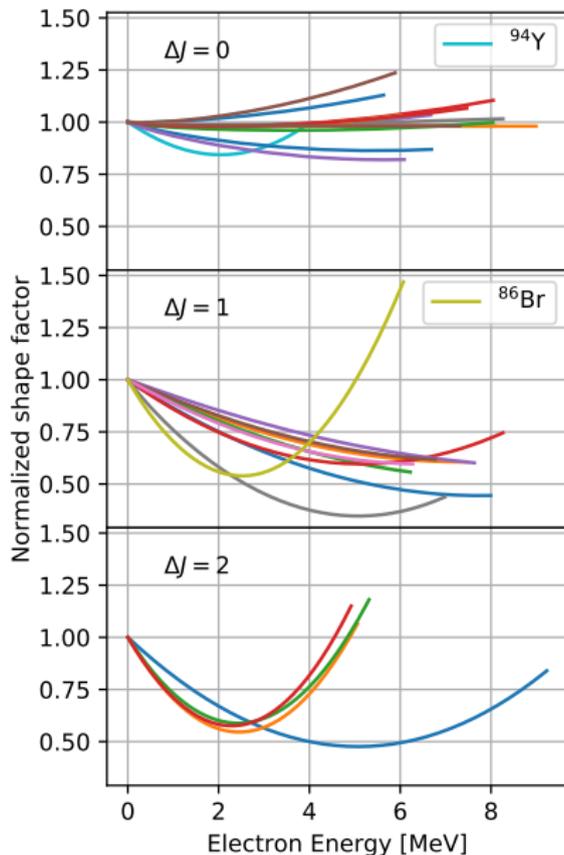
$$C(Z, E)$$

Allowed: $C \approx 1$

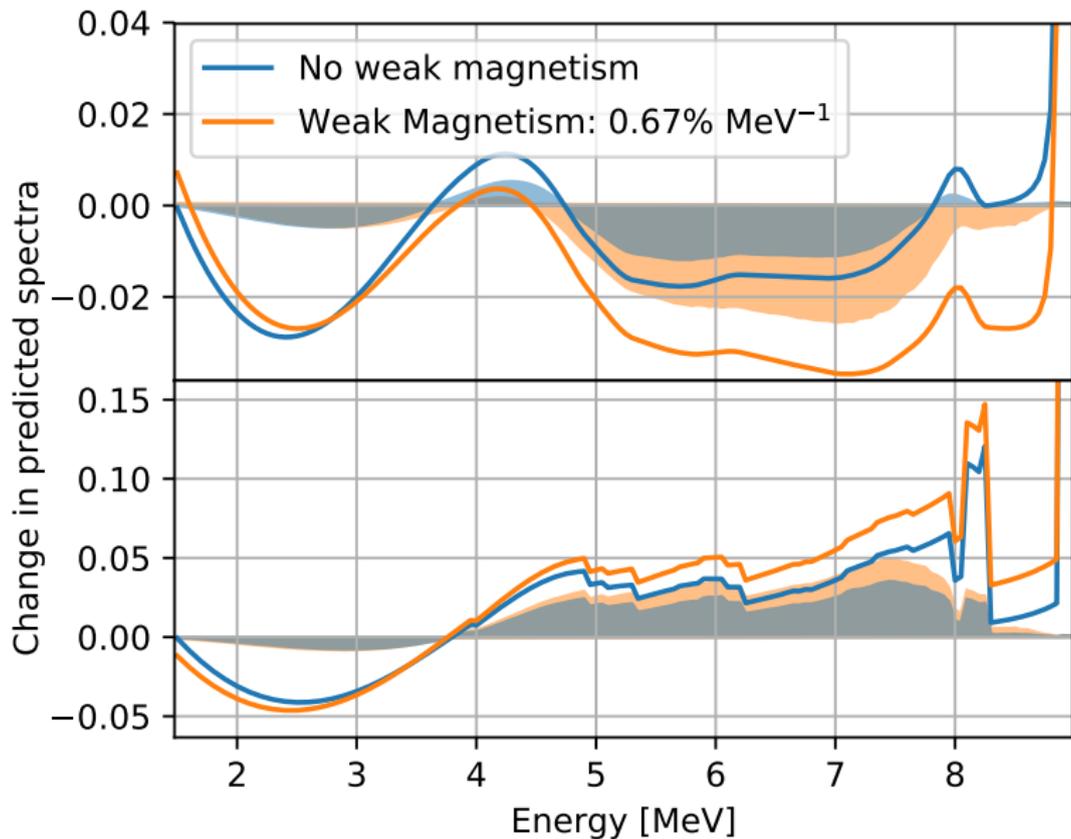
As expected,

large spectral changes

LH *et al.*, PRC 99 (2019) 013301(R)



Spectral changes



Can we use knowledge of these transitions to say something about the other 3000?

Parametrization

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Uniform behaviour for each ΔJ separately invites a parametrization

Parametrization

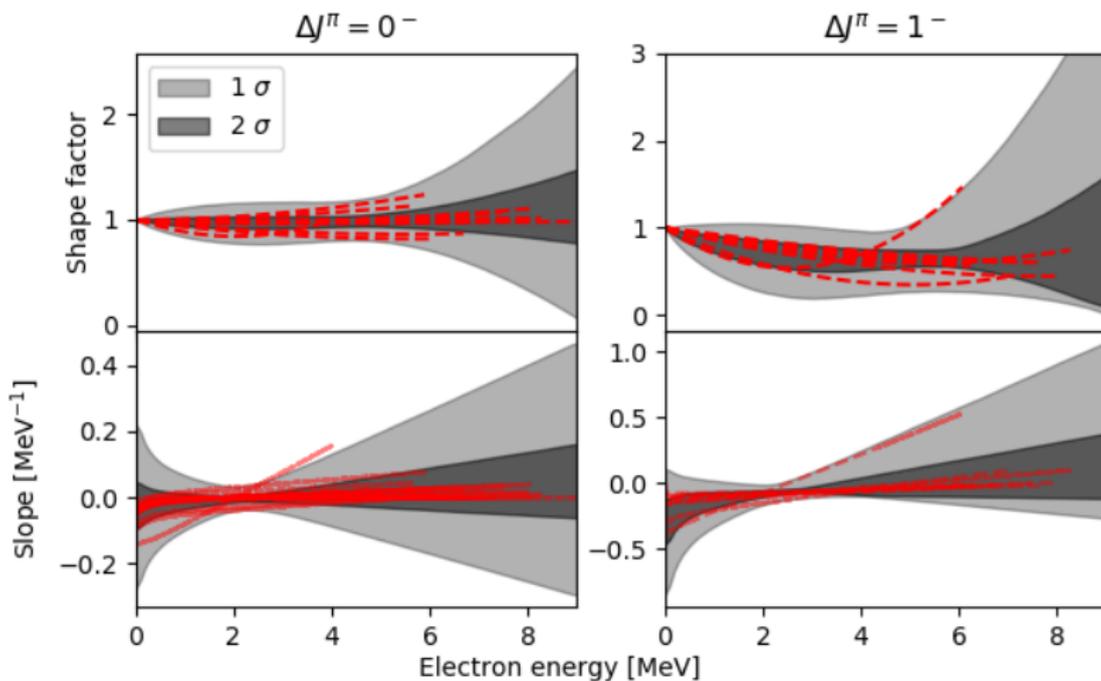
Can we use knowledge of these transitions to say something about the other 3000?

Uniform behaviour for each ΔJ separately invites a parametrization

Fit each calculated shape factor using simple polynomial, obtain distributions of correlated fit parameters for each ΔJ

Parametrization

Construct conservative shape factor distributions for each ΔJ



Directly applicable to conversion method!

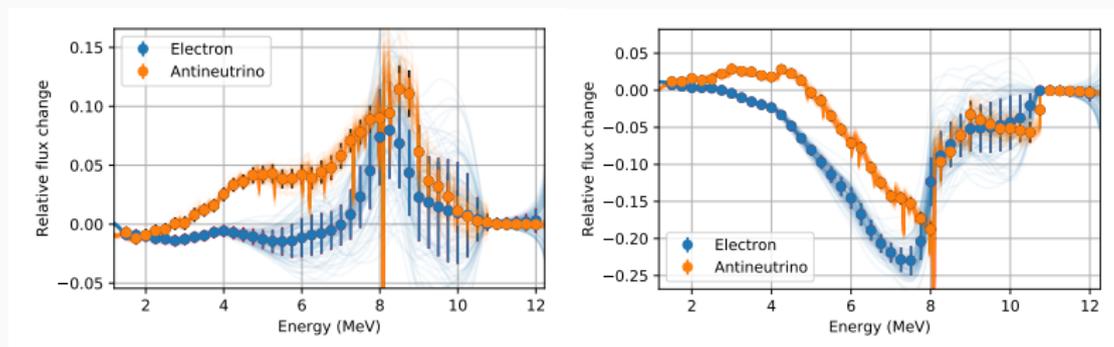
Forbidden spectral changes

Perform Monte Carlo sampling over **all** forbidden branches to propagate uncertainty into final calculation

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Look at difference in cumulative spectrum shapes



Allowed

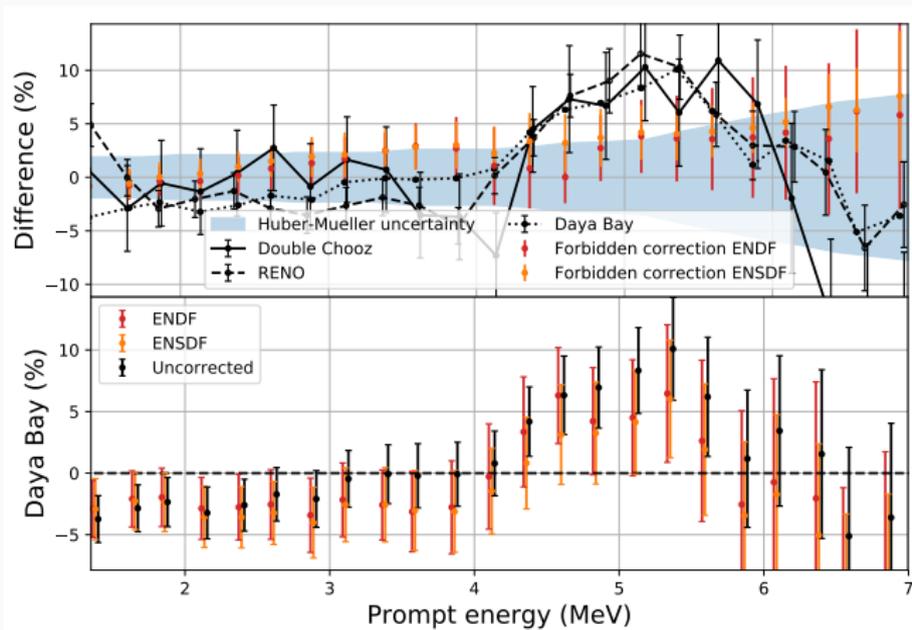
Unique

Large spectral changes for **all** actinides

Monte Carlo allows for uncertainty estimation

Forbidden transitions & the bump

Use spectrum changes with Schreckenbach correspondence



Bump **significantly mitigated**, still further research

LH, J. Kostensalo, N. Severijns, J. Suhonen, PRC 99 (2019) 031301(R)

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Estimate uncertainty using Monte Carlo methods

Reactor bump is **significantly mitigated**, increased uncertainty weakens anomaly

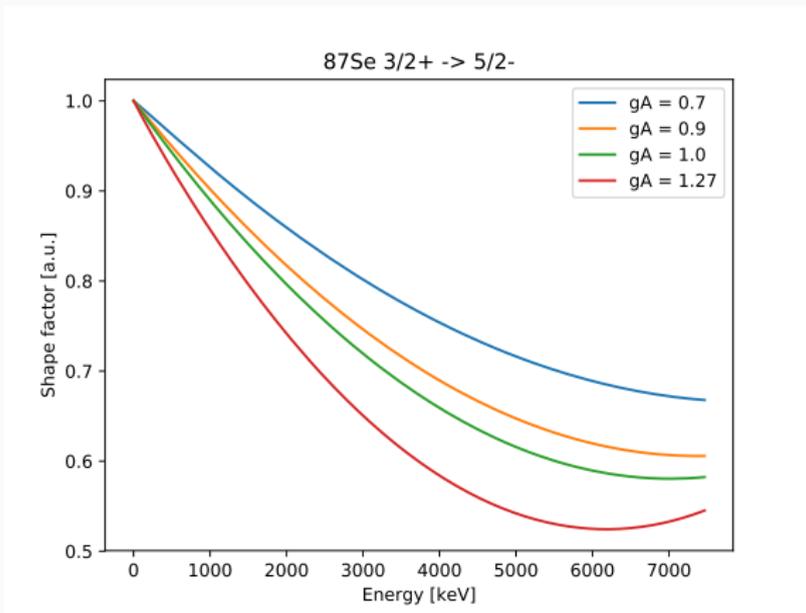
Uncertainty estimation

Care only about shape, not absolute magnitude

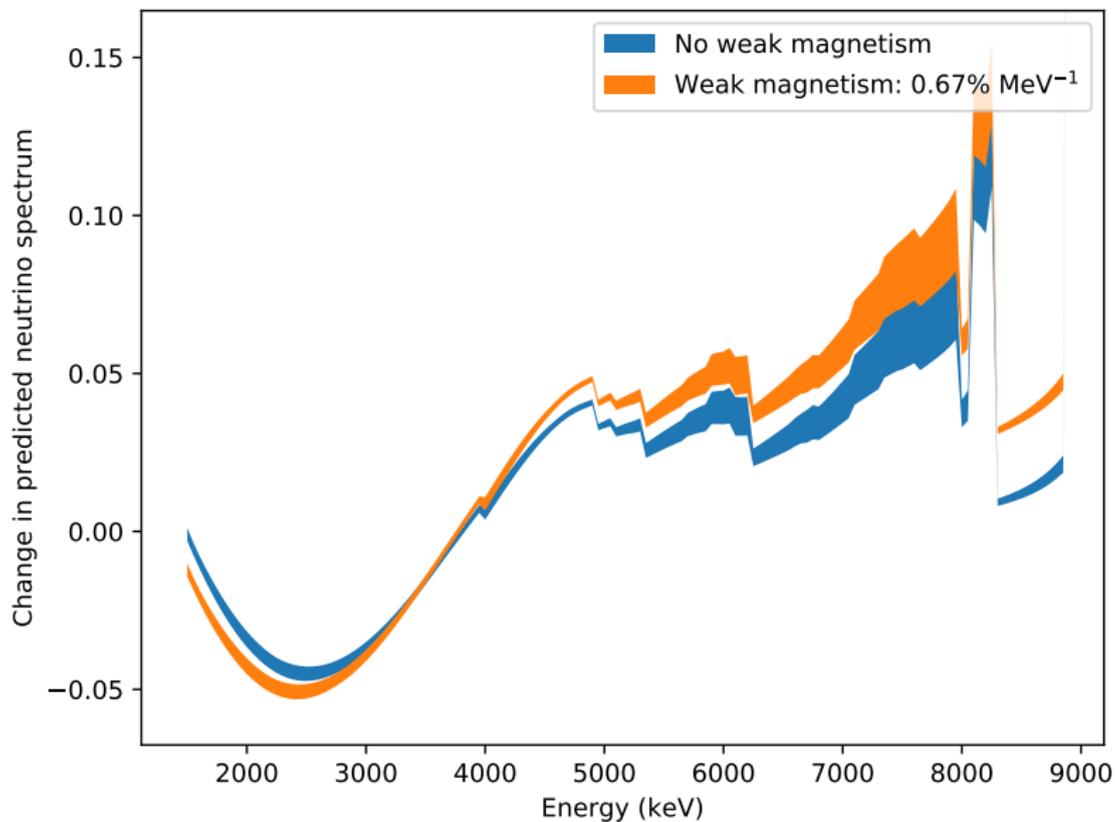
Uncertainty estimation

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Shape factor can have (significant) dependence on what value for g_A^{eff} is used



Uncertainty Estimation



Matrix elements

Up to first order, deal with 6 matrix elements

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ΔJ	Matrix elements	Forbidden
0	${}^A\mathcal{M}_{000}$	-
1	${}^A\mathcal{M}_{111}, {}^A\mathcal{M}_{121}, {}^V\mathcal{M}_{101}, {}^A\mathcal{M}_{110}$	$0 \rightarrow 0$
2	${}^A\mathcal{M}_{211}$	$0 \rightarrow 0, \frac{1}{2} \rightarrow \frac{1}{2}, 1 \rightarrow 0$

Behrens-Bühning notation: ${}^{V/A}\mathcal{M}_{KLs}$

- s : timelike (0, scalar) or spacelike (1, vector)
- L : Angular momentum from multipole decomposition
- K : total J of operator ($|L - s| \leq K \leq L + s$)

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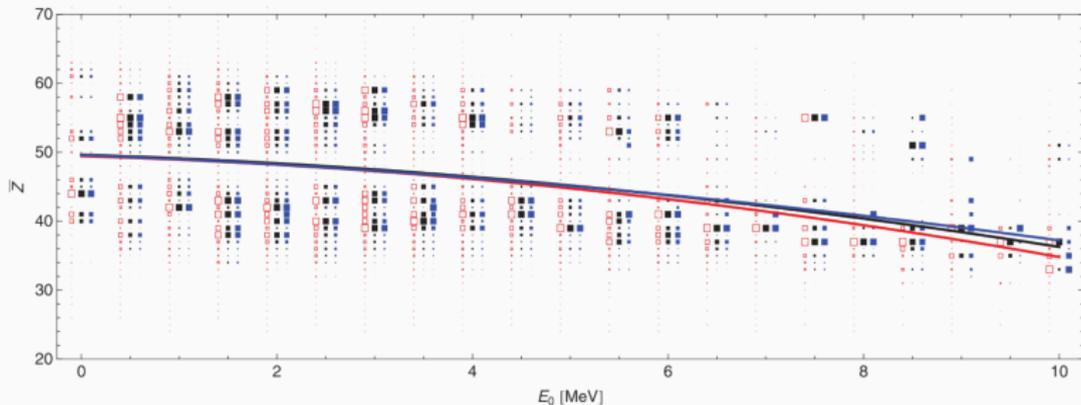
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General: 1 dominant matrix element \rightarrow easy, $> 1 \rightarrow$ harder

Modern conversion analysis

Extrapolation & Virtual branches

How to construct these fictitious β branches?



Parametrised $\bar{Z}(E_0)$ fit with simple polynomial

P. Huber, PRC **84** (2011) 024617

Extrapolation & Virtual branches

Typical procedure

1. Make grid for E_0 in $[2, 12]$ MeV
2. Every gridpoint $E_{0,i}$, choose $Z(E_{0,i})$
3. Assume allowed shape, extrapolate average nuclear matrix elements
4. Fit VB intensities to cumulative exp. spectrum

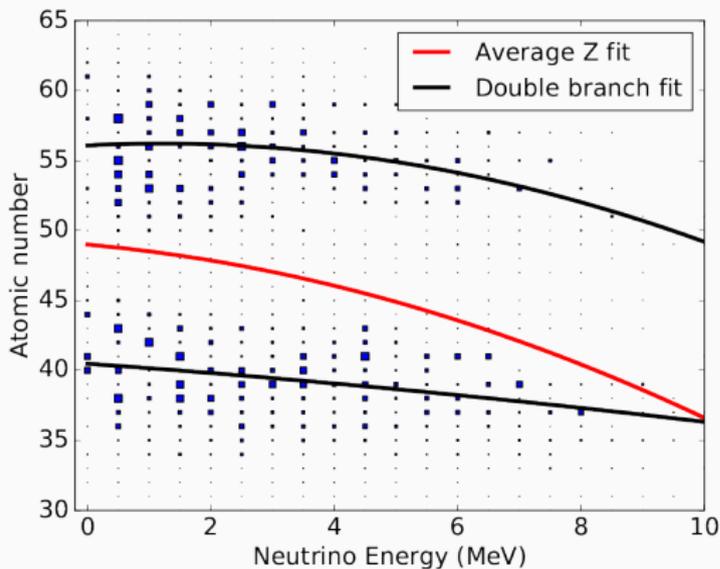
$$S(E_e) = \sum_i c_i S(E_e, \bar{Z}(E_{0,i}), E_{0,i})$$

5. Invert spectra using $E_\nu = E_0 - E_e$

Database extrapolation

Database contains much more information to use

Trivial extension
to improve
 $(\alpha Z)^2$ behaviour,
fixed weights

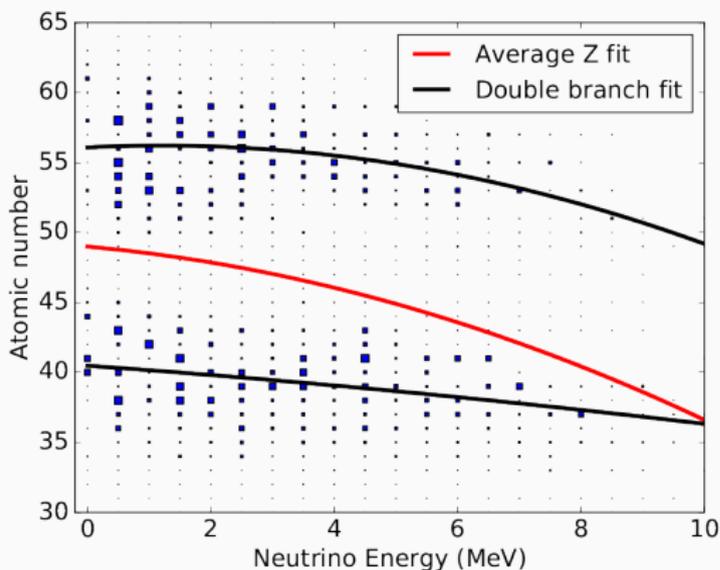


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Employ
Machine Learning
clustering
algorithms to find
better patterns



Nuclear β decays live in high-dimensional vector spaces

- Z, A
- Log ft values
- Branching Ratio, E_0 , daughter excitation
- $\Delta J^{\Delta\pi}$ (forbiddenness, unique)
- Initial and final deformation
- ...

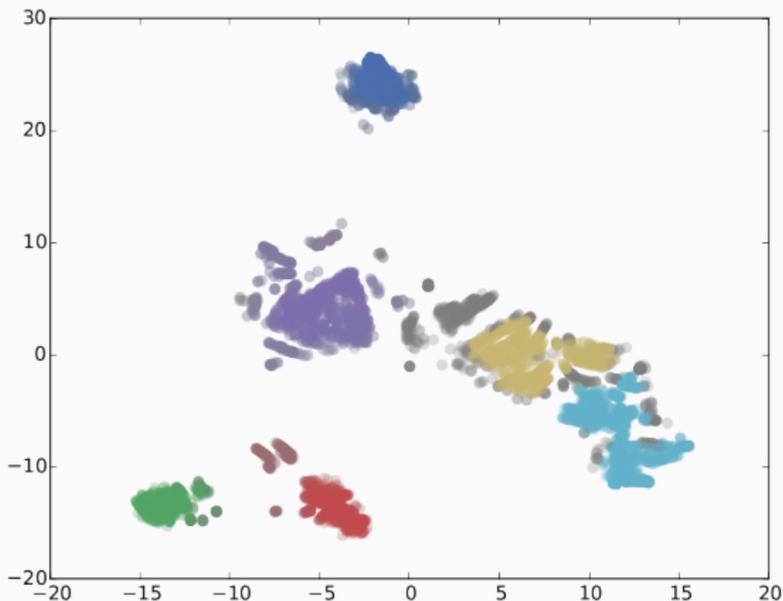
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Clusters in high dimensions are smeared in 2D projections

Clustering visualisation

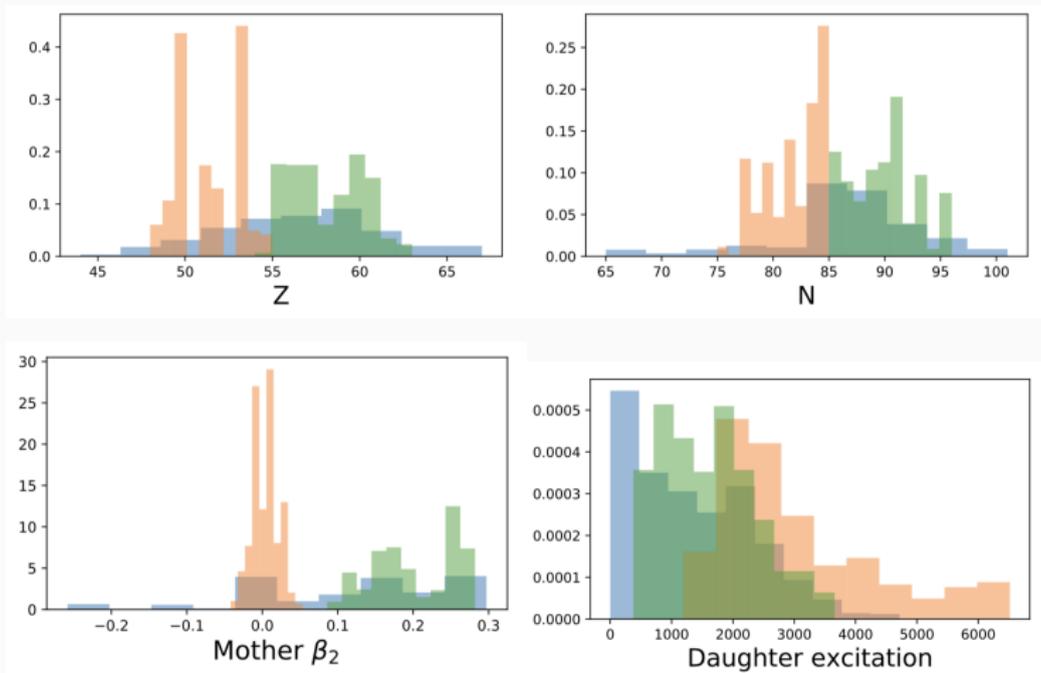
Use dimensional reduction (t-SNE) to visualise results



Clear clusters, intercluster distance irrelevant here

Intercluster comparison

Example comparison for 3 clusters



Large differences visible for simple histograms!

How to combine these results?

Instead of a single $Z(E_0)$ fit, use

Multidimensional Cluster Markov Chain Monte Carlo (MC³)

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Build a **distribution** of anomaly → better uncertainty estimate

Virtual β branch creation

Procedure:

For each E_0 bin, for each cluster, build sampling distribution

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$$P(\theta|d) \propto P(\theta)P(d|\theta)$$

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fission yield \times BR

Likelihood ($P(d|\theta)$): probability for point to belong to cluster

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Modification of prior allows for compensation/study of
pandemonium

MC³ moving forward

Clusters contain nuclear structure information, can stochastically deduce matrix element corrections

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Done correctly, *realistic uncertainty & anomaly* including correlations